

MA 4448  
Assignment 1  
Due 2 February 2011

Id: 4448-1011-1.m4,v 1.3 2011/02/14 23:49:59 john Exp john

1. Figure out how the ordinary derivatives  $A_{j_1 \dots j_s, k}^{i_1 \dots i_r}$  of an arbitrary  $(r, s)$  tensor  $A_{j_1 \dots j_s}^{i_1 \dots i_r}$  transform under coordinate changes.

If you get lost in the indices then can do this for a  $(2, 2)$ , although the general case is actually easier.

*Solution:* In the coordinates  $\bar{x}^0, \bar{x}^1, \bar{x}^2, \bar{x}^3$ ,

$$\begin{aligned} \bar{A}_{j_1 \dots j_s, k}^{i_1 \dots i_r} &= \frac{\partial}{\partial \bar{x}^k} \bar{A}_{j_1 \dots j_s}^{i_1 \dots i_r} \\ &= \frac{\partial x^n}{\partial \bar{x}^k} \frac{\partial}{\partial x^n} \left( \prod_{1 \leq p \leq r} \frac{\partial \bar{x}^{i_p}}{\partial x^{l_p}} \prod_{1 \leq q \leq s} \frac{\partial x^{m_q}}{\partial \bar{x}^{j_q}} A_{m_1 \dots m_s}^{l_1 \dots l_r} \right) \\ &= \prod_{1 \leq p \leq r} \frac{\partial \bar{x}^{i_p}}{\partial x^{l_p}} \prod_{1 \leq q \leq s} \frac{\partial x^{m_q}}{\partial \bar{x}^{j_q}} \frac{\partial x^n}{\partial \bar{x}^k} \frac{\partial}{\partial x^n} A_{m_1 \dots m_s}^{l_1 \dots l_r} \\ &\quad + \sum_{1 \leq u \leq r} \prod_{\substack{1 \leq p \leq r \\ p \neq u}} \frac{\partial \bar{x}^{i_p}}{\partial x^{l_p}} \prod_{1 \leq q \leq s} \frac{\partial x^{m_q}}{\partial \bar{x}^{j_q}} \frac{\partial x^n}{\partial \bar{x}^k} \frac{\partial^2 \bar{x}^{i_u}}{\partial x^n \partial x^{l_u}} A_{m_1 \dots m_s}^{l_1 \dots l_r} \\ &\quad + \sum_{1 \leq v \leq s} \prod_{1 \leq p \leq r} \frac{\partial \bar{x}^{i_p}}{\partial x^{l_p}} \prod_{\substack{1 \leq q \leq s \\ q \neq v}} \frac{\partial x^{m_q}}{\partial \bar{x}^{j_q}} \frac{\partial x^n}{\partial \bar{x}^k} \frac{\partial^2 x^{m_v}}{\partial x^n \partial \bar{x}^{j_v}} A_{m_1 \dots m_s}^{l_1 \dots l_r}. \end{aligned}$$

2. Compute the Christoffel symbols of Minkowski space in

- (a) Cartesian coordinates:  $(x^0, x^1, x^2, x^3) = (t, x, y, z)$ . For your sanity and mine, write indices as  $t, x, y$  and  $z$  rather than 0, 1, 2 and 3. So the metric is  $g_{tt} = -c^2, g_{xx} = g_{yy} = g_{zz} = 1$  and all other components are zero.

*Solution:* Every derivative of every metric component is zero, so all the connection coefficients are zero.

- (b) Spherical Coordinates:  $(x^0, x^1, x^2, x^3)(t, r, \theta, \varphi)$ . The metric in these coordinates is  $g_{tt} = -c^2$ ,  $g_{rr} = 1$ ,  $g_{\theta\theta} = r^2$ ,  $g_{\varphi\varphi} = r^2 \sin^2 \theta$  and all other components are zero.

*Solution:*

The inverse metric is easily computed because the matrix representation of the metric is diagonal in this basis.

$$g^{tt} = -c^{-2} \quad g^{rr} = 1 \quad g^{\theta\theta} = r^{-2} \quad g^{\varphi\varphi} = r^{-2} \csc^2 \theta$$

The non-zero derivatives of the metric components are

$$g_{\theta\theta,r} = 2r \quad g_{\varphi\varphi,r} = 2r \sin^2 \theta \quad g_{\varphi\varphi,\theta} = 2r^2 \sin \theta \cos \theta.$$

Using the formula

$$\Gamma_{jk}^i = \frac{1}{2} g^{il} (g_{kl,j} + g_{jl,k} - g_{jk,l})$$

we see that, of the 64 entries of the Christoffel symbol, the only non-zero ones are

$$\Gamma_{\theta\theta}^r = -r \quad \Gamma_{r\theta}^\theta = \Gamma_{\theta r}^\theta = r^{-1} = \Gamma_{r\varphi}^\varphi = \Gamma_{\varphi r}^\varphi$$

$$\Gamma_{\varphi\varphi}^r = -r \sin^2 \theta \quad \Gamma_{\varphi\varphi}^\theta = -\sin \theta \cos \theta \quad \Gamma_{\varphi\theta}^\varphi = \Gamma_{\theta\varphi}^\varphi = \cot \theta$$

3. Read about raising and lowering of indices in Section 3.1.1 of George Ellis' notes. It's fairly simple. Then,

- (a) Show that  $g_{;k}^{ij}$  is zero.

*Solution:*

There is no need to deal with Christoffel symbols. We already know that the covariant derivative of  $g$  with two lower indices is zero. That was part of the definition of the covariant derivative. We also know that the covariant derivative  $g$  with an upper and lower index is zero. That was proved in lecture. So,

$$0 = g_{l;k}^i = (g^{im} g_{lm})_{;k} = g_{;k}^{im} g_{lm} + g^{im} g_{lm;k} = g_{;k}^{im} g_{lm}.$$

Contract with  $g^{lj}$ ,

$$0 = g_{;k}^{im} g_{lm} g^{jl} = g_{;k}^{im} g_m^j = g_{;k}^{ij}.$$

- (b) What do raising and lowering correspond to in the coordinate-free description of tensors?

*Solution:* The metric tensor is a non degenerate quadratic form on  $T_xM$ . It gives us an isomorphism between  $T_xM$  and  $T_x^*M$  as follows. Take a vector in  $T_xM$ . Taking the quadratic form applied to this vector and some other vector, we get something linear in the second vector. In other words, we get an element of  $T_x^*M$ . This element depends linearly on the first vector. So we have a linear map from  $T_xM$  to  $T_x^*M$ . The fact that this has zero nullspace follows immediately from non-degeneracy. The vector spaces have the same dimension, so an injective linear function is an isomorphism. Lowering an index corresponds to applying this isomorphism to one of the  $T_xM$  factors in the tensor product. Raising an index corresponds to applying its inverse to one of the  $T_x^*M$  factors.

- (c) Prove that we get the same result from differentiating and then raising/lowering as we get from raising/lowering and then differentiating. Differentiating here means covariant differentiation.

*Solution:* There is no need to mention Christoffel symbols at all. We raise an index by contracting with  $g^{ij}$  and lower by contracting with  $g_{ij}$ . Applying the Leibnitz rule, we get two terms. The first is  $g$  contracted with the covariant derivative of our tensor. That's what we want. The second term is zero, because all covariant derivatives of  $g$  are zero, whether it has upper or lower indices.