

PYU44T20 Quantum Optics and Information

Problem Set 1 due 14/02/2023

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SS Theoretical Physics

1 Linear Vector Spaces & Operators

1.1)

For a given Hermitian operator $\hat{A}$ we have

$$\hat{A} = \begin{pmatrix} a_{00} & a_{01} \\ a_{10} & a_{11} \end{pmatrix} = \begin{pmatrix} \alpha + \beta_z & \beta_x - i\beta_y \\ \beta_x + i\beta_y & \alpha - \beta_z \end{pmatrix}.$$

This is equivalent to

$$a_{00} = \alpha + \beta_z, \quad (1)$$

$$a_{11} = \alpha - \beta_z, \quad (2)$$

$$a_{01} = \beta_x - i\beta_y, \quad (3)$$

$$a_{10} = \beta_x + i\beta_y. \quad (4)$$

These can be rearranged to obtain expressions for α and $\vec{\beta}$ as

$$(1) + (2) \implies \alpha = \frac{1}{2} (a_{00} + a_{11}), \quad (5)$$

$$(1) - (2) \implies \beta_z = \frac{1}{2} (a_{00} - a_{11}), \quad (6)$$

$$(3) + (4) \implies \beta_x = \frac{1}{2} (a_{01} + a_{10}), \quad (7)$$

$$(3) - (4) \implies \beta_y = \frac{i}{2} (a_{01} - a_{10}). \quad (8)$$

If $\hat{A} = \hat{0}$, i.e. $a_{00} = a_{01} = a_{10} = a_{11} = 0$, then we also have $\alpha = \beta_x = \beta_y = \beta_z = 0$, and so the elements of $\{\hat{1}, \hat{\sigma}_i\}$ are linearly independent.

For general $\hat{A}$, from $\hat{A}^\dagger = \hat{A}$ we have

$$a_{00}^* = a_{00} \implies a_{00} \in \mathbb{R}, \quad (9)$$

$$a_{11}^* = a_{11} \implies a_{11} \in \mathbb{R}, \quad (10)$$

$$a_{01} = a_{10}^*. \quad (11)$$

Combining these constraints with the expressions for $\alpha, \vec{\beta}$ gives

$$(5), (9), (10) \implies \alpha \in \mathbb{R}$$

$$(6), (9), (10) \implies \beta_z \in \mathbb{R}$$

$$(7), (11) \implies \beta_x^* = \beta_x \implies \beta_x \in \mathbb{R}$$

$$(8), (11) \implies \beta_y^* = \beta_y \implies \beta_y \in \mathbb{R}$$

Thus $\{\hat{1}, \hat{\sigma}_i\}$ spans the space of Hermitian operators, and so any $\hat{A}$ can be expressed in the desired form $\hat{A} = \alpha \hat{1} + \vec{\beta} \cdot \vec{\sigma}$.

1.2)

$$\begin{aligned}
\left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right)^2 &= (\beta_x \hat{\sigma}_x + \beta_y \hat{\sigma}_y + \beta_z \hat{\sigma}_z)^2 \\
&= (\beta_x^2 \hat{\sigma}_x \hat{\sigma}_x + \beta_y^2 \hat{\sigma}_y \hat{\sigma}_y + \beta_z^2 \hat{\sigma}_z \hat{\sigma}_z) \\
&\quad + \beta_x \beta_y (\hat{\sigma}_x \hat{\sigma}_y + \hat{\sigma}_y \hat{\sigma}_x) + \beta_x \beta_z (\hat{\sigma}_x \hat{\sigma}_z + \hat{\sigma}_z \hat{\sigma}_x) + \beta_y \beta_z (\hat{\sigma}_y \hat{\sigma}_z + \hat{\sigma}_z \hat{\sigma}_y) \\
&= \left(\frac{\beta_x^2}{2} \{\hat{\sigma}_x, \hat{\sigma}_x\} + \frac{\beta_y^2}{2} \{\hat{\sigma}_y, \hat{\sigma}_y\} + \frac{\beta_z^2}{2} \{\hat{\sigma}_z, \hat{\sigma}_z\} \right) \\
&\quad + \beta_x \beta_y \{\hat{\sigma}_x, \hat{\sigma}_y\} + \beta_x \beta_z \{\hat{\sigma}_x, \hat{\sigma}_z\} + \beta_y \beta_z \{\hat{\sigma}_y, \hat{\sigma}_z\} \\
&= (\beta_x^2 + \beta_y^2 + \beta_z^2) \hat{1} + 0 \quad \text{(from } \{\hat{\sigma}_i, \hat{\sigma}_j\} = 2\delta_{ij} \hat{1})
\end{aligned}$$

1.3)

$$\begin{aligned}
\left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right)^{2n} &= \left[\left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right)^2 \right]^n \\
&= \left[(\vec{\beta} \cdot \vec{\beta}) \hat{1} \right]^n \\
&= (\vec{\beta} \cdot \vec{\beta})^n \hat{1}
\end{aligned}$$

$$\begin{aligned}
\left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right)^{2n+1} &= \left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right)^{2n} \left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right) \\
&= (\vec{\beta} \cdot \vec{\beta})^n \hat{1} \left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right) \\
&= (\vec{\beta} \cdot \vec{\beta})^n \left(\vec{\beta} \cdot \hat{\vec{\sigma}}\right)
\end{aligned}$$

1.4)

$$\hat{H}_0 = \begin{pmatrix} E & K \\ K^* & E \end{pmatrix} = \alpha \hat{1} + \vec{\beta} \cdot \hat{\vec{\sigma}}$$

$$(5), (6), (7), (8) \implies \alpha = E, \quad \beta_x = \frac{K + K^*}{2} = \text{Re}(K), \quad \beta_y = \frac{i(K - K^*)}{2} = -\text{Im}(K), \quad \beta_z = 0$$

$$\begin{aligned} \hat{U}(t) &= \exp\left(-\frac{i\hat{H}_0 t}{\hbar}\right) \\ &= \exp\left(-\frac{it}{\hbar} (\alpha \hat{1} + \vec{\beta} \cdot \hat{\vec{\sigma}})\right) \\ &= \exp\left(-\frac{it\alpha}{\hbar} \hat{1}\right) \exp\left(-\frac{it}{\hbar} \vec{\beta} \cdot \hat{\vec{\sigma}}\right) \\ &= \exp\left(-\frac{it\alpha}{\hbar} \hat{1}\right) \sum_{k=0}^{\infty} \frac{\left(-\frac{it}{\hbar}\right)^k}{k!} (\vec{\beta} \cdot \hat{\vec{\sigma}})^k \\ &= \exp\left(-\frac{iEt}{\hbar} \hat{1}\right) \sum_{n=0}^{\infty} \left[\frac{\left(-\frac{it}{\hbar}\right)^{2n}}{(2n)!} (\vec{\beta} \cdot \hat{\vec{\sigma}})^{2n} + \frac{\left(-\frac{it}{\hbar}\right)^{2n+1}}{(2n+1)!} (\vec{\beta} \cdot \hat{\vec{\sigma}})^{2n+1} \right] \\ &= \exp\left(-\frac{iEt}{\hbar} \hat{1}\right) \sum_{n=0}^{\infty} \left[\frac{(-1)^n \left(\frac{t}{\hbar}\right)^{2n}}{(2n)!} (\beta_x^2 + \beta_y^2 + \beta_z^2)^n \hat{1} - i \frac{(-1)^n \left(\frac{t}{\hbar}\right)^{2n+1}}{(2n+1)!} (\beta_x^2 + \beta_y^2 + \beta_z^2)^n (\vec{\beta} \cdot \hat{\vec{\sigma}}) \right] \\ &= \exp\left(-\frac{iEt}{\hbar} \hat{1}\right) \sum_{n=0}^{\infty} \left[\frac{(-1)^n \left(\frac{t}{\hbar} \sqrt{\text{Re}(K)^2 + (-\text{Im}(K))^2}\right)^{2n}}{(2n)!} \hat{1} \right. \\ &\quad \left. - \frac{i(-1)^n \left(\frac{t}{\hbar} \sqrt{\text{Re}(K)^2 + (-\text{Im}(K))^2}\right)^{2n+1}}{(2n+1)! \sqrt{\text{Re}(K)^2 + (-\text{Im}(K))^2}} (\beta_x \hat{\sigma}_x + \beta_y \hat{\sigma}_y) \right] \\ &= \exp\left(-\frac{iEt}{\hbar} \hat{1}\right) \left[\cos\left(\frac{|K|t}{\hbar}\right) \hat{1} - \frac{i}{|K|} \sin\left(\frac{|K|t}{\hbar}\right) (\beta_x \hat{\sigma}_x + \beta_y \hat{\sigma}_y) \right] \\ &= \exp\left(-\frac{iEt}{\hbar}\right) \left[\cos\left(\frac{|K|t}{\hbar}\right) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \frac{i \sin\left(\frac{|K|t}{\hbar}\right)}{|K|} \left(\text{Re}(K) \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + i \text{Im}(K) \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \right) \right] \\ &= \exp\left(-\frac{iEt}{\hbar}\right) \begin{pmatrix} \cos\left(\frac{|K|t}{\hbar}\right) & -i \frac{K}{|K|} \sin\left(\frac{|K|t}{\hbar}\right) \\ -i \frac{K^*}{|K|} \sin\left(\frac{|K|t}{\hbar}\right) & \cos\left(\frac{|K|t}{\hbar}\right) \end{pmatrix} \end{aligned}$$

1.5)

$$\begin{aligned}
0 &= \det(\hat{H}_0 - \lambda \hat{1}) \\
&= \begin{vmatrix} E - \lambda & K \\ K^* & E - \lambda \end{vmatrix} \\
&= (E - \lambda)^2 - |K|^2 \\
\Rightarrow \lambda_{\pm} &= E \pm |K|
\end{aligned}$$

$$\hat{H}_0 |v_{\pm}\rangle = \lambda_{\pm} |v_{\pm}\rangle \Rightarrow \begin{pmatrix} E & K \\ K^* & E \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix} = \begin{pmatrix} aE + bK \\ aK^* + bE \end{pmatrix} = \begin{pmatrix} a(E \pm |K|) \\ b(E \pm |K|) \end{pmatrix}$$

$$\Rightarrow aE + bK = a(E \pm |K|), \quad aK^* + bE = b(E \pm |K|)$$

$$\Rightarrow b = \pm a \frac{|K|}{K} = \pm a \frac{K^*}{|K|} \Rightarrow |v_{\pm}\rangle = a \left(|0\rangle \pm \frac{|K|}{K} |1\rangle \right)$$

$$\begin{aligned}
1 &= \langle v_{\pm} | v_{\pm} \rangle \\
&= |a|^2 \left(1^2 + \frac{|K|^2}{|K|^2} \right) \\
\Rightarrow |a|^2 &= \frac{1}{2} \\
\Rightarrow |v_{\pm}\rangle &= \frac{e^{i\phi}}{\sqrt{2}} \left(|0\rangle \pm \frac{|K|}{K} |1\rangle \right), \phi \in \mathbb{R}
\end{aligned}$$

We can express $\hat{H}_0$ as a matrix in the standard basis or the eigenbasis, i.e.

$$\begin{aligned}
\hat{H}_0 &= E(|0\rangle\langle 0| + |1\rangle\langle 1|) + K|0\rangle\langle 1| + K^*|1\rangle\langle 0| & \hat{H}_0 &= \lambda_+ |v_+\rangle\langle v_+| + \lambda_- |v_-\rangle\langle v_-| \\
&= \begin{pmatrix} E & K \\ K^* & E \end{pmatrix}_{|0\rangle, |1\rangle} & &= \begin{pmatrix} E + |K| & 0 \\ 0 & E - |K| \end{pmatrix}_{|v_{\pm}\rangle}
\end{aligned}$$

The corresponding form of $\hat{U}(t)$ as a matrix in the eigenbasis is thus given by

$$\begin{aligned}
\hat{U}(t) &= \exp\left(-\frac{i\hat{H}_0 t}{\hbar}\right) \\
&= \exp\left(-\frac{it}{\hbar} \text{diag}(E + |K|, E - |K|)_{|v_{\pm}\rangle}\right) \\
\hat{U}(t) &= \text{diag}\left(\exp\left(-\frac{it}{\hbar}(E + |K|)\right), \exp\left(-\frac{it}{\hbar}(E - |K|)\right)\right)_{|v_{\pm}\rangle} \\
&= \begin{pmatrix} \exp\left(-\frac{it}{\hbar}(E + |K|)\right) & 0 \\ 0 & \exp\left(-\frac{it}{\hbar}(E - |K|)\right) \end{pmatrix}_{|v_{\pm}\rangle}
\end{aligned}$$

2 Projection Operators

2.1)

$$\begin{aligned}
 \hat{P}_i^2 &= (|i\rangle \langle i|)(|i\rangle \langle i|) & \hat{P}_i^\dagger &= (|i\rangle \langle i|)^\dagger \\
 &= |i\rangle \langle i| i \langle i| & &= (\langle i|)^\dagger (|i\rangle)^\dagger \\
 &= |i\rangle \hat{1} \langle i| & &= |i\rangle \langle i| \\
 &= \hat{P}_i & &= \hat{P}_i
 \end{aligned}$$

$$\begin{aligned}
 |i\rangle \text{ not normalised} &\implies \hat{P}_i^2 = \langle i|i\rangle \hat{P}_i \neq \hat{P}_i \\
 &\implies \text{first condition } \hat{P}_i^2 = \hat{P}_i \text{ is not satisfied}
 \end{aligned}$$

2.2)

$$\begin{aligned}
 (\hat{P}_i + \hat{P}_j)^2 &= (|i\rangle \langle i| + |j\rangle \langle j|)(|i\rangle \langle i| + |j\rangle \langle j|) \\
 &= |i\rangle \langle i|i\rangle \langle i| + |i\rangle \langle i|j\rangle \langle j| + |j\rangle \langle j|i\rangle \langle i| + |j\rangle \langle j|j\rangle \langle j| \\
 &= |i\rangle \langle i| + 0 + 0 + |j\rangle \langle j| \\
 &= \hat{P}_i + \hat{P}_j
 \end{aligned}$$

$$\begin{aligned}
 (\hat{P}_i + \hat{P}_j)^\dagger &= (|i\rangle \langle i| + |j\rangle \langle j|)^\dagger \\
 &= (|i\rangle \langle i|)^\dagger + (|j\rangle \langle j|)^\dagger \\
 &= \hat{P}_i^\dagger + \hat{P}_j^\dagger \\
 &= \hat{P}_i + \hat{P}_j
 \end{aligned}$$

2.3)

$$\begin{aligned}
 \sum_{i=1}^{\dim \mathcal{H}} \langle C_i | \hat{A} | C_i \rangle &= \sum_{i=1}^{\dim \mathcal{H}} \langle C_i | \hat{A} \left(\sum_{j=1}^{\dim \mathcal{H}} |B_j\rangle \langle B_j| C_i \rangle \right) \\
 &= \sum_{i,j=1}^{\dim \mathcal{H}} \langle C_i | \hat{A} | B_j \rangle \langle B_j | C_i \rangle \\
 &= \sum_{j=1}^{\dim \mathcal{H}} \left(\sum_{i=1}^{\dim \mathcal{H}} \langle B_j | C_i \rangle \langle C_i | \right) \hat{A} | B_j \rangle \\
 &= \sum_{j=1}^{\dim \mathcal{H}} \langle B_j | \hat{A} | B_j \rangle
 \end{aligned}$$

2.4)

$$\begin{aligned}
\text{Tr}[\hat{P}_i \hat{A}] &= \sum_{j=1}^{\dim \mathcal{H}} \langle B_j | \hat{P}_i \hat{A} | B_j \rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} \langle B_j | B_i \rangle \langle B_i | \hat{A} | B_j \rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} \delta_{ji} \langle B_i | \hat{A} | B_j \rangle \\
&= \langle B_i | \hat{A} | B_i \rangle
\end{aligned}$$

2.5)

$$\begin{aligned}
\text{Tr}[\hat{A}_1 \hat{A}_2] &= \sum_{i=1}^{\dim \mathcal{H}} \langle B_i | \hat{A}_1 \hat{A}_2 | B_i \rangle \\
&= \sum_{i=1}^{\dim \mathcal{H}} \langle B_i | \hat{A}_1 \hat{1} \hat{A}_2 | B_i \rangle \\
&= \sum_{i=1}^{\dim \mathcal{H}} \langle B_i | \hat{A}_1 \left(\sum_{j=1}^{\dim \mathcal{H}} |B_j\rangle \langle B_j| \right) \hat{A}_2 | B_i \rangle \\
&= \sum_{i,j=1}^{\dim \mathcal{H}} \langle B_i | \hat{A}_1 | B_j \rangle \langle B_j | \hat{A}_2 | B_i \rangle \\
&= \sum_{i,j=1}^{\dim \mathcal{H}} \langle B_j | \hat{A}_2 | B_i \rangle \langle B_i | \hat{A}_1 | B_j \rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} \langle B_j | \hat{A}_2 \left(\sum_{i=1}^{\dim \mathcal{H}} |B_i\rangle \langle B_i| \right) \hat{A}_1 | B_j \rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} \langle B_j | \hat{A}_2 \hat{1} \hat{A}_1 | B_j \rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} \langle B_j | \hat{A}_2 \hat{A}_1 | B_j \rangle \\
&= \text{Tr}[\hat{A}_2 \hat{A}_1]
\end{aligned}$$

2.6)

$$\begin{aligned}
|D_i\rangle &= \hat{1} |D_i\rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} |D_j\rangle \langle D_j | D_i \rangle \\
&= \sum_{j=1}^{\dim \mathcal{H}} \langle D_j | D_i \rangle |D_j\rangle \\
\text{also } |D_i\rangle &= \sum_{j=1}^{\dim \mathcal{H}} \delta_{ji} |D_j\rangle \implies \langle D_j | D_i \rangle = \delta_{ji} \implies |D_i\rangle \text{ are orthonormal}
\end{aligned}$$

3 Translation Operators

3.1)

$$\begin{aligned}
\left(\prod_{i=1}^N \hat{U}_i\right) \left(\prod_{j=1}^N \hat{U}_j\right)^\dagger &= \hat{U}_1 \cdots \hat{U}_N \hat{U}_N^\dagger \cdots \hat{U}_1^\dagger \\
&= \hat{U}_1 \cdots \hat{U}_{N-1} \hat{\mathbb{1}} \hat{U}_{N-1}^\dagger \cdots \hat{U}_1^\dagger \\
&= \hat{U}_1 \cdots \hat{U}_{N-1} \hat{U}_{N-1}^\dagger \cdots \hat{U}_1^\dagger \\
&= \dots \\
&= \hat{U}_1 \hat{U}_1^\dagger \\
&= \hat{\mathbb{1}} \\
\Rightarrow \left(\prod_{i=1}^N \hat{U}_i\right)^\dagger &= \left(\prod_{i=1}^N \hat{U}_i\right)^{-1} \Rightarrow \text{unitary}
\end{aligned}$$

3.2)

$$\begin{aligned}
(\hat{U}_1 + \hat{U}_2) (\hat{U}_1 + \hat{U}_2)^\dagger &= (\hat{U}_1 + \hat{U}_2) (\hat{U}_1^\dagger + \hat{U}_2^\dagger) \\
&= \hat{U}_1 \hat{U}_1^\dagger + \hat{U}_1 \hat{U}_2^\dagger + \hat{U}_2 \hat{U}_1^\dagger + \hat{U}_2 \hat{U}_2^\dagger \\
&= \hat{\mathbb{1}} + \hat{U}_1 \hat{U}_2^\dagger + \hat{U}_2 \hat{U}_1^\dagger + \hat{\mathbb{1}}
\end{aligned}$$

Thus the sum is only unitary if $\hat{U}_1 \hat{U}_2^\dagger + \hat{U}_2 \hat{U}_1^\dagger = -\hat{\mathbb{1}}$, which is not generally true for arbitrary unitary $\hat{U}_1, \hat{U}_2$, e.g.

$$\begin{aligned}
\hat{\sigma}_i \hat{\sigma}_j^\dagger + \hat{\sigma}_j \hat{\sigma}_i^\dagger &= \hat{\sigma}_i \hat{\sigma}_j + \hat{\sigma}_j \hat{\sigma}_i \\
&= 2\delta_{ij} \hat{\mathbb{1}} \neq -\hat{\mathbb{1}} \\
\hat{U} \hat{U}^\dagger + \hat{U} \hat{U}^\dagger &= 2\hat{\mathbb{1}} \neq -\hat{\mathbb{1}}
\end{aligned}$$

3.3)

$$\begin{aligned}
(\hat{T}_1)^2 &= \left(\sum_{i=-\infty}^{\infty} |x_{i-1}\rangle \langle x_i|\right) \left(\sum_{j=-\infty}^{\infty} |x_{j-1}\rangle \langle x_j|\right) \\
&= \sum_{i,j=-\infty}^{\infty} |x_{i-1}\rangle \langle x_i | x_{j-1}\rangle \langle x_j| \\
&= \sum_{i,j=-\infty}^{\infty} |x_{i-1}\rangle \delta_{i,j-1} \langle x_j| \\
&= \sum_{j=-\infty}^{\infty} |x_{j-2}\rangle \langle x_j| \\
&= \hat{T}_2
\end{aligned}$$

3.4)

We can prove this by induction, i.e. prove for $k = 1$, and then prove for $k = n + 1$ assuming true for $k = n$.

$$\left(\hat{T}_1\right)^1 = \hat{T}_1 \implies \text{true for } k = 1$$

Now assume this is true for $k = n$.

$$\begin{aligned} \left(\hat{T}_1\right)^{n+1} &= \left(\hat{T}_1\right)^n \hat{T}_1 \\ &= \hat{T}_n \hat{T}_1 \\ &= \left(\sum_{i=-\infty}^{\infty} |x_{i-n}\rangle \langle x_i|\right) \left(\sum_{j=-\infty}^{\infty} |x_{j-1}\rangle \langle x_j|\right) \\ &= \sum_{i,j=-\infty}^{\infty} |x_{i-n}\rangle \langle x_i | x_{j-1}\rangle \langle x_j| \\ &= \sum_{i,j=-\infty}^{\infty} |x_{i-n}\rangle \delta_{i,j-1} \langle x_j| \\ &= \sum_{j=-\infty}^{\infty} |x_{j-(n+1)}\rangle \langle x_j| \\ &= \hat{T}_{n+1} \implies \text{true for } k = n + 1 \end{aligned}$$

$$\begin{aligned} \text{True for } k = 1 &\implies \text{true for } k = 2 \\ &\implies \text{true for } k = 3 \\ &\implies \dots \\ &\implies \text{true for } k \in \mathbb{Z}^+ \end{aligned}$$

3.5)

$$\begin{aligned} \hat{T}_1 \hat{T}_1^\dagger &= \left(\sum_{i=-\infty}^{\infty} |x_{i-1}\rangle \langle x_i|\right) \left(\sum_{j=-\infty}^{\infty} |x_{j-1}\rangle \langle x_j|\right)^\dagger \\ &= \left(\sum_{i=-\infty}^{\infty} |x_{i-1}\rangle \langle x_i|\right) \left(\sum_{j=-\infty}^{\infty} |x_j\rangle \langle x_{j-1}|\right) \\ &= \sum_{i,j=-\infty}^{\infty} |x_{i-1}\rangle \langle x_i | x_j\rangle \langle x_{j-1}| \\ &= \sum_{i,j=-\infty}^{\infty} |x_{i-1}\rangle \delta_{ij} \langle x_{j-1}| \\ &= \sum_{j=-\infty}^{\infty} |x_{j-1}\rangle \langle x_{j-1}| \\ &= \sum_{k=-\infty}^{\infty} |x_k\rangle \langle x_k| \quad (k = j - 1) \\ &= \hat{\mathbb{1}} \\ &\implies \hat{T}_1^\dagger = \hat{T}_1^{-1} \implies \hat{T}_1 \text{ is unitary} \end{aligned}$$

From 3.4), $\hat{T}_k = \left(\hat{T}_1\right)^k$, and so is a product of k unitary operators. From 3.1), this product is also a unitary operator, and thus $\hat{T}_k$ must also be unitary.

3.6)

$$\begin{aligned}
\hat{T}_p \hat{X} \hat{T}_p^\dagger &= \left(\sum_{i=-\infty}^{\infty} |x_{i-p}\rangle \langle x_i| \right) \left(\sum_{j=-\infty}^{\infty} x_j |x_j\rangle \langle x_j| \right) \left(\sum_{k=-\infty}^{\infty} |x_{k-p}\rangle \langle x_k| \right)^\dagger \\
&= \sum_{i,j,k=-\infty}^{\infty} x_j |x_{i-p}\rangle \langle x_i | x_j\rangle \langle x_j | x_k\rangle \langle x_{k-p}| \\
&= \sum_{i,j,k=-\infty}^{\infty} j a |x_{i-p}\rangle \delta_{ij} \delta_{jk} \langle x_{k-p}| \\
&= \sum_{j=-\infty}^{\infty} j a |x_{j-p}\rangle \langle x_{j-p}| \\
&= \sum_{l=-\infty}^{\infty} (l+p) a |x_l\rangle \langle x_l| \quad (l = j - p) \\
&= \sum_{l=-\infty}^{\infty} l a |x_l\rangle \langle x_l| + \sum_{l=-\infty}^{\infty} p a |x_l\rangle \langle x_l| \\
&= \sum_{l=-\infty}^{\infty} x_l |x_l\rangle \langle x_l| + p a \sum_{l=-\infty}^{\infty} |x_l\rangle \langle x_l| \\
&= \hat{X} + p a \hat{\mathbb{1}}
\end{aligned}$$

3.7)

$$\begin{aligned}
\hat{H}_0 &= -\frac{\hbar^2}{2ma^2} \left(\hat{T}_1 + \hat{T}_1^\dagger - 2\hat{\mathbb{1}} \right) \\
&= -\frac{\hbar^2}{2ma^2} \left(\hat{T}_1 + \hat{T}_1^{-1} - 2\hat{\mathbb{1}} \right) \\
&= -\frac{\hbar^2}{2ma^2} \left[\exp\left(\frac{i\hat{P}a}{\hbar}\right) + \exp\left(-\frac{i\hat{P}a}{\hbar}\right) - 2\hat{\mathbb{1}} \right] \\
&= -\frac{\hbar^2}{2ma^2} \left[\sum_{k=0}^{\infty} \frac{\left(\frac{ia}{\hbar}\right)^k}{k!} \hat{P}^k + \sum_{k=0}^{\infty} \frac{\left(-\frac{ia}{\hbar}\right)^k}{k!} \hat{P}^k - 2\hat{\mathbb{1}} \right] \\
&= -\frac{\hbar^2}{2ma^2} \left[\sum_{k=0}^{\infty} \frac{\left(\frac{a}{\hbar}\right)^k}{k!} (i^k + (-i)^k) \hat{P}^k - 2\hat{\mathbb{1}} \right] \\
&= -\frac{\hbar^2}{2ma^2} \left[(1+1)\hat{\mathbb{1}} + \frac{a}{\hbar} (i-i)\hat{P} + \frac{a^2}{2\hbar^2} (-1-1)\hat{P}^2 + \mathcal{O}(a^3) - 2\hat{\mathbb{1}} \right] \\
&= -\frac{\hbar^2}{2ma^2} \left[2\hat{\mathbb{1}} + 0 - \frac{a^2}{\hbar^2} \hat{P}^2 + \mathcal{O}(a^3) - 2\hat{\mathbb{1}} \right] \\
&= \frac{\hat{P}^2}{2m} + \mathcal{O}(a) \\
\Rightarrow \lim_{a \rightarrow 0} \hat{H}_0 &= \frac{\hat{P}^2}{2m}
\end{aligned}$$

3.8)

$$\begin{aligned}
\hat{T}_1 |\theta\rangle &= \left(\sum_{j=-\infty}^{\infty} |x_{j-1}\rangle \langle x_j| \right) \left(\sum_{k=-\infty}^{\infty} e^{i\theta k} |x_k\rangle \right) \\
&= \sum_{j,k=-\infty}^{\infty} e^{i\theta k} |x_{j-1}\rangle \langle x_j | x_k\rangle \\
&= \sum_{j,k=-\infty}^{\infty} e^{i\theta k} |x_{j-1}\rangle \delta_{jk} \\
&= \sum_{j=-\infty}^{\infty} e^{i\theta j} |x_{j-1}\rangle \\
&= \sum_{l=-\infty}^{\infty} e^{i\theta(l+1)} |x_l\rangle \quad (l = j - 1) \\
&= e^{i\theta} \sum_{l=-\infty}^{\infty} e^{i\theta l} |x_l\rangle \\
&= e^{i\theta} |\theta\rangle
\end{aligned}$$

Thus $|\theta\rangle$ are eigenstates of $\hat{T}_1$ with eigenvalues $e^{i\theta}$, $\theta \in [-\pi, \pi]$.

$$\begin{aligned}
\hat{T}_1^\dagger (\hat{T}_1 |\theta\rangle) &= \hat{T}_1^\dagger e^{i\theta} |\theta\rangle = e^{i\theta} \hat{T}_1^\dagger |\theta\rangle \\
\text{also } (\hat{T}_1^\dagger \hat{T}_1) |\theta\rangle &= \hat{1} |\theta\rangle \\
\implies \hat{T}_1^\dagger |\theta\rangle &= e^{-i\theta} |\theta\rangle
\end{aligned}$$

Thus $|\theta\rangle$ are also eigenstates of $\hat{T}_1^\dagger$ with eigenvalues $e^{-i\theta}$.

$$\begin{aligned}
\hat{H}_0 |\theta\rangle &= -\frac{\hbar^2}{2ma^2} (\hat{T}_1 |\theta\rangle + \hat{T}_1^\dagger |\theta\rangle - 2\hat{1} |\theta\rangle) \\
&= -\frac{\hbar^2}{2ma^2} (e^{i\theta} + e^{-i\theta} - 2) |\theta\rangle
\end{aligned}$$

Thus $|\theta\rangle$ are eigenstates of $\hat{H}_0$ with eigenvalues $-\frac{\hbar^2}{2ma^2} (e^{i\theta} + e^{-i\theta} - 2) = \frac{\hbar^2}{ma^2} (1 - \cos \theta) = \frac{2\hbar^2}{ma^2} \sin^2 \frac{\theta}{2}$.

3.9)

$$\begin{aligned}
\hat{H}_1 &= -\frac{\hbar^2}{2ma^2} \left[\hat{T}_1 + \hat{T}_1^\dagger + \alpha (\hat{T}_1^2 - 2\hat{T}_1 + \hat{1}) + \alpha^* (\hat{T}_1^{\dagger 2} - 2\hat{T}_1^\dagger + \hat{1}) - 2\hat{1} \right] \\
&= -\frac{\hbar^2}{2ma^2} \left[(\alpha + \alpha^* - 2) \hat{1} + (1 - 2\alpha) \hat{T}_1 + (1 - 2\alpha^*) \hat{T}_1^\dagger + \alpha \hat{T}_1^2 + \alpha^* \hat{T}_1^{\dagger 2} \right] \\
\implies \hat{H}_1 |\theta\rangle &= -\frac{\hbar^2}{2ma^2} \left[(\alpha + \alpha^* - 2) + (1 - 2\alpha) e^{i\theta} + (1 - 2\alpha^*) e^{-i\theta} + \alpha \hat{T}_1 e^{i\theta} + \alpha^* \hat{T}_1^\dagger e^{-i\theta} \right] |\theta\rangle \\
&= -\frac{\hbar^2}{2ma^2} \left[(\alpha + \alpha^* - 2) + (1 - 2\alpha) e^{i\theta} + (1 - 2\alpha^*) e^{-i\theta} + \alpha e^{2i\theta} + \alpha^* e^{-2i\theta} \right] |\theta\rangle \\
&= -\frac{\hbar^2}{2ma^2} \left[(e^{i\theta} + e^{-i\theta} - 2) + \alpha (e^{2i\theta} - 2e^{i\theta} + 1) + \alpha^* (e^{-2i\theta} - 2e^{-i\theta} + 1) \right] |\theta\rangle \\
&= -\frac{\hbar^2}{2ma^2} \left[(e^{i\theta} + e^{-i\theta} - 2) + \alpha e^{i\theta} (e^{i\theta} + e^{-i\theta} - 2) + \alpha^* e^{-i\theta} (e^{i\theta} + e^{-i\theta} - 2) \right] |\theta\rangle \\
&= -\frac{\hbar^2}{2ma^2} (\alpha e^{i\theta} + \alpha^* e^{-i\theta} + 1) (e^{i\theta} + e^{-i\theta} - 2) |\theta\rangle
\end{aligned}$$

Thus $|\theta\rangle$ is an eigenstate of $\hat{H}_1$ with eigenvalues $-\frac{\hbar^2}{2ma^2} (\alpha e^{i\theta} + \alpha^* e^{-i\theta} + 1) (e^{i\theta} + e^{-i\theta} - 2)$.