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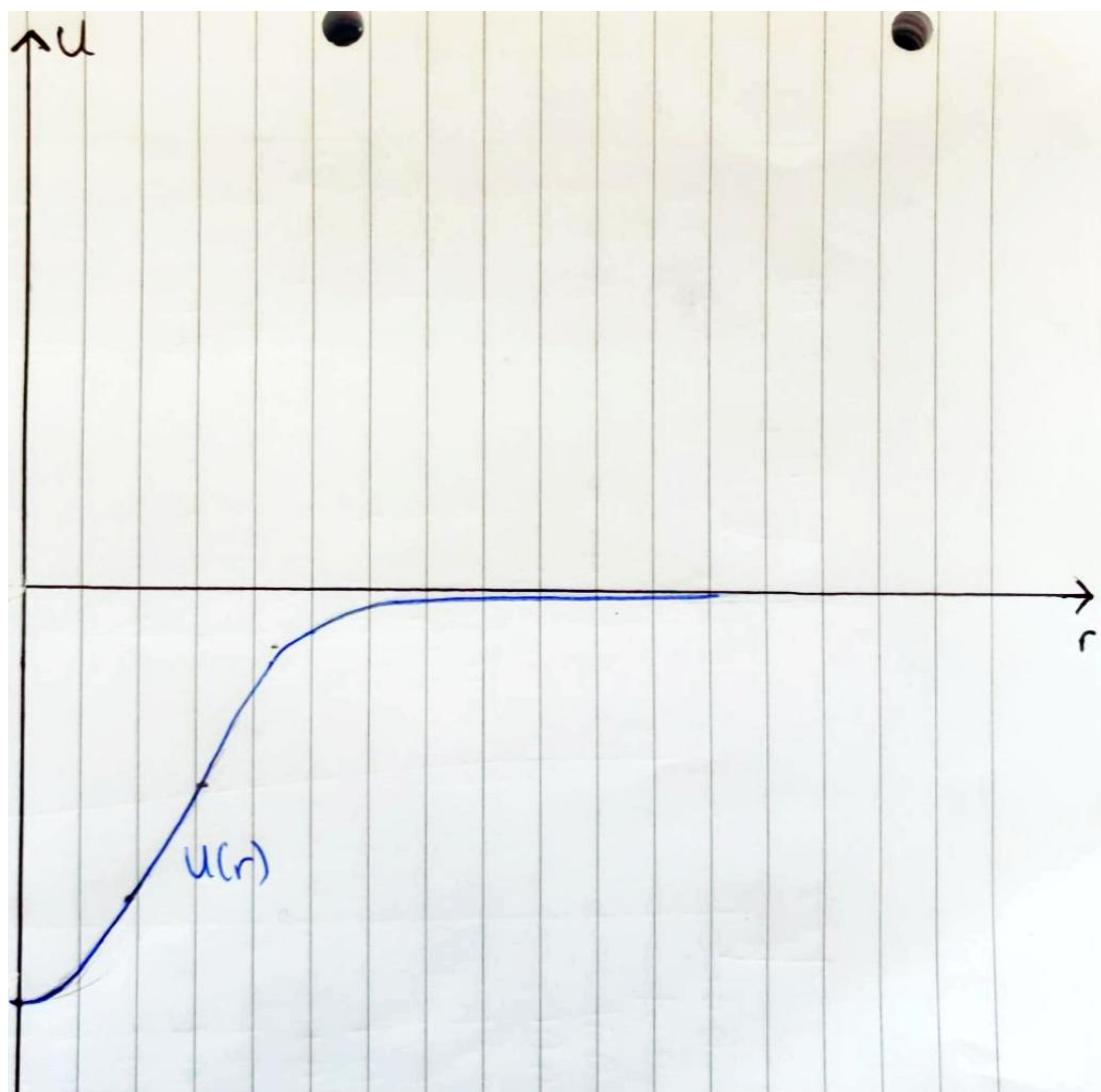
SF Theoretical Physics

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Problem 1

(a)



The field is attractive everywhere, as $F = -\frac{dU}{dr} < 0 \forall r > 0$.

(b)

$$\begin{aligned} U_{\text{eff}}(r) &= U(r) + \frac{M^2}{2mr^2} \\ &= \frac{M^2}{2mr^2} - \frac{\alpha^2}{2m(\beta^2 + r^2)} \end{aligned}$$



Finite motion is possible. The particle is bounded by the annulus of radii r_{\min} and r_{\max} . It may not necessarily be a closed loop.

Infinite motion is also possible. The particle can travel from infinity and get deflected around the origin, and continue to infinity.

$$U_{\text{eff}}(r_{\text{turning}}) = E$$

$$\text{Finite motion} \implies r_{\min} = \frac{1}{2} \sqrt{\frac{M^2 - \alpha^2}{Em} - 2\beta^2 + \frac{\sqrt{8EmM^2\beta^2 + (\alpha^2 + 2Em\beta^2 - M^2)^2}}{Em}}$$

$$\text{and } r_{\max} = \frac{1}{2} \sqrt{\frac{M^2 - \alpha^2}{Em} - 2\beta^2 - \frac{\sqrt{8EmM^2\beta^2 + (\alpha^2 + 2Em\beta^2 - M^2)^2}}{Em}}$$

$$\text{Infinite motion} \implies r_{\min} = \frac{1}{2} \sqrt{\frac{M^2 - \alpha^2}{Em} - 2\beta^2 - \frac{\sqrt{8EmM^2\beta^2 + (\alpha^2 + 2Em\beta^2 - M^2)^2}}{Em}}$$

Assuming $\left[\{r \geq 0 \& m > 0 \& \alpha > 0 \& \beta > 0 \& e < 0\}, \text{Solve} \left[\frac{M^2}{2m * r^2} - \frac{\alpha^2}{2m * (\beta^2 + r^2)} = e, r \right] \right]$

$$\left\{ \left\{ r \rightarrow -\frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 - \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\}, \left\{ r \rightarrow \frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 - \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\}, \right. \\ \left. \left\{ r \rightarrow -\frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 + \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\}, \left\{ r \rightarrow \frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 + \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\} \right\}$$

Assuming $\left[\{r \geq 0 \& m > 0 \& \alpha > 0 \& \beta > 0 \& e > 0\}, \text{Solve} \left[\frac{M^2}{2m * r^2} - \frac{\alpha^2}{2m * (\beta^2 + r^2)} = e, r \right] \right]$

$$\left\{ \left\{ r \rightarrow -\frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 - \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\}, \left\{ r \rightarrow \frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 - \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\}, \right. \\ \left. \left\{ r \rightarrow -\frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 + \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\}, \left\{ r \rightarrow \frac{1}{2} \sqrt{\frac{M^2}{em} - \frac{\alpha^2}{em} - 2\beta^2 + \frac{\sqrt{8emM^2\beta^2 + (-M^2 + \alpha^2 + 2em\beta^2)^2}}{em}} \right\} \right\}$$

(c)

$$\begin{aligned}
 \frac{dU_{\text{eff}}(r)}{dr} \Big|_{r=r_0} &= 0 \\
 \implies r_0 &= \beta \sqrt{\frac{M}{\alpha - M}} \\
 E &= U_{\text{eff}}(r_0) \\
 &= \frac{M(-M + \alpha)}{2m\beta^2} - \frac{\alpha^2}{2m \left(\beta^2 + \frac{M\beta^2}{-M + \alpha} \right)} \\
 E &= -\frac{(\alpha - M)^2}{2m\beta^2}
 \end{aligned}$$

$$\text{Ellipse} \implies 2m\pi ab = TM$$

$$\text{Circle} \implies 2m\pi r_0^2 = TM$$

$$\begin{aligned}
 T &= \frac{2m\pi r_0^2}{M} \\
 T &= \frac{2m\pi\beta^2}{\alpha - M}
 \end{aligned}$$

$$\mathbf{U}_{\text{eff}}[\mathbf{r}_-] := \frac{\mathsf{M}^2}{2 \mathsf{m} * \mathbf{r}^2} - \frac{\alpha^2}{2 \mathsf{m} * (\beta^2 + \mathbf{r}^2)} ;$$

$$\mathbf{Solve}[\mathbf{D}[\mathbf{U}_{\text{eff}}[\mathbf{r}], \mathbf{r}] == 0, \mathbf{r}]$$

$$\left\{ \left\{ \mathbf{r} \rightarrow -\frac{\sqrt{\mathsf{M}} \beta}{\sqrt{-\mathsf{M} - \alpha}} \right\}, \left\{ \mathbf{r} \rightarrow \frac{\sqrt{\mathsf{M}} \beta}{\sqrt{-\mathsf{M} - \alpha}} \right\}, \left\{ \mathbf{r} \rightarrow -\frac{\sqrt{\mathsf{M}} \beta}{\sqrt{-\mathsf{M} + \alpha}} \right\}, \left\{ \mathbf{r} \rightarrow \frac{\sqrt{\mathsf{M}} \beta}{\sqrt{-\mathsf{M} + \alpha}} \right\} \right\}$$

$$\mathbf{r}_\theta = \frac{\sqrt{\mathsf{M}} \beta}{\sqrt{-\mathsf{M} + \alpha}} ;$$

$$\mathbf{U}_{\text{eff}}[\mathbf{r}_\theta]$$

$$\frac{\mathsf{M} (-\mathsf{M} + \alpha)}{2 \mathsf{m} \beta^2} - \frac{\alpha^2}{2 \mathsf{m} \left(\beta^2 + \frac{\mathsf{M} \beta^2}{-\mathsf{M} + \alpha} \right)}$$

(d)

$$\begin{aligned}
 E &< 0 \\
 \phi(r) &= \text{constant} \\
 \mathbf{M} &= 0 \text{ as } \vec{r} \parallel \vec{v} \text{ and } \vec{M} = \vec{r} \times m\vec{v} \\
 \implies U_{\text{eff}}(r) &= U(r)
 \end{aligned}$$

$$E = U(A), A \equiv \text{amplitude}$$

$$E = -\frac{\alpha^2}{2m(\beta^2 + A^2)}$$

$$\begin{aligned}
A &= \sqrt{\frac{-\alpha^2 - 2Em\beta^2}{2Em}} \\
T &= \sqrt{2m} \int_{x_a}^{x_b} \frac{dx}{\sqrt{E - U(x)}} \\
&= 2\sqrt{2m} \int_A^0 \frac{dr}{\sqrt{\frac{\alpha^2}{2m(\beta^2+r^2)} - \frac{\alpha^2}{2m(\beta^2+A^2)}}} \\
T &= \frac{4m\beta\sqrt{\beta^2 + A^2}}{\alpha} \text{EllipticE}\left(-\frac{A^2}{\beta^2}\right) \\
&\approx \frac{2m\pi\beta^2}{\alpha} \text{ for small } A
\end{aligned}$$

For small amplitudes, the period does not depend on the amplitude. This makes sense, as the motion will be approximately simple harmonic, for which the period does not depend on amplitude.

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Clear[r, e]
Solve[e == -\frac{\alpha^2}{2 m * (\beta^2 + r^2)}, r]
{{r \rightarrow -\frac{\sqrt{-\alpha^2 - 2 e m \beta^2}}{\sqrt{2} \sqrt{e} \sqrt{m}}}, {r \rightarrow \frac{\sqrt{-\alpha^2 - 2 e m \beta^2}}{\sqrt{2} \sqrt{e} \sqrt{m}}}}

Clear[r, e];
A = \frac{\sqrt{-\alpha^2 - 2 e m \beta^2}}{\sqrt{2} \sqrt{e} \sqrt{m}};
2 \sqrt{2 m} Assuming[{\alpha > 0 && \beta > 0 && m > 0 && e < 0 && -\alpha^2 - 2 e * m * \beta^2 > 0}, Integrate[\frac{1}{\sqrt{\frac{\alpha^2}{2 m * (\beta^2 + r^2)} - \frac{\alpha^2}{2 m * (\beta^2 + A^2)}}}, {r, A, 0}]]
2 \sqrt{2} \sqrt{m} \beta \text{EllipticE}\left[1 + \frac{\alpha^2}{2 e m \beta^2}\right] if 2 e m + \alpha^2 > 0

Clear[e, A];
e = -\frac{\alpha^2}{2 m * (\beta^2 + A^2)};
2 \sqrt{2} \sqrt{m} \beta \text{EllipticE}\left[1 + \frac{\alpha^2}{2 e m \beta^2}\right] // FullSimplify
4 \sqrt{m} \beta \text{EllipticE}\left[-\frac{A^2}{\beta^2}\right]
\sqrt{\frac{\alpha^2}{m (A^2 + \beta^2)}}

Clear[A];
Series[\frac{4 \sqrt{m} \beta \text{EllipticE}\left[-\frac{A^2}{\beta^2}\right]}{\sqrt{\frac{\alpha^2}{m (A^2 + \beta^2)}}}, {A, 0, 0}] // Normal // FullSimplify
2 \sqrt{m} \pi \beta
\sqrt{\frac{\alpha^2}{m \beta^2}}

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Problem 2

(a)

$$\begin{aligned}
\mathcal{L} &= \frac{m}{2}v^2 - U(r) && \text{Lagrangian for the particle} \\
&= \frac{m}{2}(\dot{r}^2 + r^2\dot{\phi}^2) - U(r) && \text{in polar coordinates} \\
M &= mr^2\dot{\phi} = \text{constant} && \dot{\phi} \text{ is cyclic} \\
\implies \dot{\phi} &= \frac{M}{mr^2} \\
E &= T + U \\
&= \frac{m}{2}(\dot{r}^2 + r^2\dot{\phi}^2) + U(r) \\
&= \frac{m}{2}\dot{r}^2 + \frac{M^2}{2mr^2} + U(r) \\
&= \frac{m}{2}\left(\frac{dr}{dt}\right)^2 + U_{\text{eff}}(r) && U_{\text{eff}}(r) \equiv U(r) + \frac{M^2}{2mr^2} \\
\frac{dr}{dt} &= \pm\sqrt{\frac{2(E - U_{\text{eff}}(r))}{m}} && \text{positive } t \implies +, \text{ negative } t \implies - \\
&= \text{sign}(t)\sqrt{\frac{2(E - U_{\text{eff}}(r))}{m}} \\
\frac{d\phi}{dr} &= \frac{d\phi}{dt} \frac{dt}{dr} \\
&= \dot{\phi}\left(\frac{dr}{dt}\right)^{-1} \\
&= \frac{M}{mr^2}\text{sign}(t)\sqrt{\frac{m}{2(E - U_{\text{eff}}(r))}} \\
&= \frac{M}{\sqrt{2m}}\text{sign}(t)\frac{1}{r^2\sqrt{E - U_{\text{eff}}(r)}} \\
\phi &= \frac{M}{\sqrt{2m}} \int_{r_a}^{r_b} \frac{dr}{r^2\sqrt{E - U_{\text{eff}}(r)}} \\
\phi_0 &= \frac{M}{\sqrt{2m}} \int_{r_{\min}}^{\infty} \frac{dr}{r^2\sqrt{E - U_{\text{eff}}(r)}} \\
\chi &\equiv |\pi - 2\phi_0| \\
&= \left| \pi - M\sqrt{\frac{2}{m}} \int_{r_{\min}}^{\infty} \frac{dr}{r^2\sqrt{E - U_{\text{eff}}(r)}} \right|
\end{aligned}$$

(b)

$$\begin{aligned}
E &= \frac{m}{2}v_{\infty}^2 + U_{\infty} && \text{where } v_{\infty} \equiv \text{velocity at infinity} \\
M &= m\rho v_{\infty} && \text{and } U_{\infty} \equiv U(r = \infty) \\
\implies \chi &= \left| \pi - m\rho v_{\infty} \sqrt{\frac{2}{m}} \int_{r_{\min}}^{\infty} \frac{dr}{r^2\sqrt{\frac{m}{2}v_{\infty}^2 + U_{\infty} - U_{\text{eff}}(r)}} \right| \\
&= \left| \pi - 2\rho \int_{r_{\min}}^{\infty} \frac{dr}{r^2\sqrt{1 - \frac{\rho^2}{r^2} - \frac{2}{mv_{\infty}^2}(U(r) - U_{\infty})}} \right|
\end{aligned}$$

(c)

(i)

$$\begin{aligned}
U_{\text{eff}}(r) &= U(r) + \frac{M^2}{2mr^2} \\
&= \frac{\alpha + \beta r^2}{r^2} + \frac{M^2}{2mr^2} \\
U_{\text{eff}}(r_{\min}) &= E \\
\implies r_{\min} &= \sqrt{\frac{M^2 + 2m\alpha}{2Em - 2m\beta}} \\
\phi(r) &= \frac{M}{\sqrt{2m}} \int_{r_{\min}}^r \frac{dr}{r^2 \sqrt{E - U_{\text{eff}}(r)}} \\
&= \frac{M}{\sqrt{M^2 + 2m\alpha}} \left(\frac{\pi}{2} + \text{arccsc} \left(r \sqrt{\frac{2m(E - \beta)}{M^2 + 2m\alpha}} \right) \right) \\
\implies r(\phi) &= \sqrt{\frac{M^2 + 2m\alpha}{2m(E - \beta)}} \sec \left(\phi \frac{\sqrt{M^2 + 2m\alpha}}{M} \right) \\
t(r) &= \pm \sqrt{\frac{m}{2}} \int \frac{dr}{\sqrt{E - U_{\text{eff}}(r)}} \\
&= \sqrt{\frac{m}{2}} \int_{r_{\min}}^r \frac{dr}{\sqrt{E - \left(\frac{\alpha + \beta r^2}{r^2} + \frac{M^2}{2mr^2} \right)}} \\
&= \frac{r}{E - \beta} \sqrt{\frac{Em}{2} - \frac{M^2 + 2m(\alpha + r^2\beta)}{4r^2}} \\
\implies r(t) &= \sqrt{\frac{M^2 + 2m\alpha + 4t^2(E - \beta)^2}{2m(E - \beta)}} \\
\text{and } \phi(t) &= \frac{M}{\sqrt{M^2 + 2m\alpha}} \left(\pi - \text{arcsec} \left(\sqrt{\frac{4t^2(E - \beta)^2}{M^2 + 2m\alpha} + 1} \right) \right)
\end{aligned}$$

$$\begin{aligned}
M &= m\rho v_{\infty} \\
E &= \frac{m}{2} v_{\infty}^2 + U_{\infty} \\
&= \frac{m}{2} v_{\infty}^2 + U(r = \infty) \\
&= \frac{m}{2} v_{\infty}^2 + \beta \\
\implies r(\phi) &= \rho \sqrt{1 + \frac{2\alpha}{m\rho^2 v_{\infty}^2}} \sec \left(\phi \sqrt{1 + \frac{2\alpha}{m\rho^2 v_{\infty}^2}} \right) \\
\text{and } r(t) &= \rho \sqrt{1 + \frac{2\alpha}{m\rho^2 v_{\infty}^2} + \frac{v_{\infty}^2 t^2}{\rho^2}} \\
\text{and } \phi(t) &= \frac{1}{\sqrt{1 + \frac{2\alpha}{m\rho^2 v_{\infty}^2}}} \left(\pi - \text{arcsec} \left(\frac{1}{\rho} \sqrt{\frac{v_{\infty}^2 t^2}{1 + \frac{2\alpha}{m\rho^2 v_{\infty}^2}} + \rho^2} \right) \right)
\end{aligned}$$

Assuming [$\{\alpha > 0 \& \beta > 0 \& m > 0\}$, Solve [$\frac{\alpha + \beta * r^2}{r^2} + \frac{M^2}{2 m * r^2} = e$, r]]

$$\left\{ \left\{ r \rightarrow -\frac{\sqrt{-M^2 - 2 m \alpha}}{\sqrt{-2 e m + 2 m \beta}} \right\}, \left\{ r \rightarrow \frac{\sqrt{-M^2 - 2 m \alpha}}{\sqrt{-2 e m + 2 m \beta}} \right\} \right\}$$

Clear[r, e]

$\frac{M}{\sqrt{2 m}}$ Assuming [$\{\alpha > 0 \& \beta > 0 \& m > 0\}$, Integrate [$\frac{1}{q^2 \sqrt{e - (\frac{\alpha+\beta*q^2}{q^2} + \frac{M^2}{2 m * q^2})}}$, {q, $\frac{\sqrt{-M^2 - 2 m \alpha}}{\sqrt{-2 e m + 2 m \beta}}$, r}]]

$$\frac{M \left(\pi + 2 \text{ArcCsc} \left[\sqrt{2} r \sqrt{\frac{m (e-\beta)}{M^2+2 m \alpha}} \right] \right)}{\sqrt{2} \sqrt{m} \sqrt{\frac{2 M^2}{m} + 4 \alpha}} \quad \text{if condition} \oplus$$

Assuming [$\{\alpha > 0 \& \beta > 0 \& m > 0\}$, Solve [$\phi == \frac{M \left(\pi + 2 \text{ArcCsc} \left[\sqrt{2} r \sqrt{\frac{m (e-\beta)}{M^2+2 m \alpha}} \right] \right)}{\sqrt{2} \sqrt{m} \sqrt{\frac{2 M^2}{m} + 4 \alpha}}$, r]]

$$\left\{ \left\{ r \rightarrow \frac{\csc \left[\frac{-M \pi + 2 \sqrt{m} \sqrt{\frac{M^2+2 m \alpha}{m}} \phi}{2 M} \right]}{\sqrt{2} \sqrt{\frac{m (e-\beta)}{M^2+2 m \alpha}}} \right\} \right\}$$

Clear[r, e]

$\sqrt{\frac{m}{2}}$ Assuming [$\{\alpha > 0 \& \beta > 0 \& m > 0 \& e > 0\}$, Integrate [$\frac{1}{\sqrt{e - (\frac{\alpha+\beta*q^2}{q^2} + \frac{M^2}{2 m * q^2})}}$, {q, $\frac{\sqrt{-M^2 - 2 m \alpha}}{\sqrt{-2 e m + 2 m \beta}}$, r}]]

$$\frac{\sqrt{m} r \sqrt{e - \frac{M^2+2 m (\alpha+r^2 \beta)}{2 m r^2}}}{\sqrt{2} (e - \beta)} \quad \text{if condition} \oplus$$

Assuming [$\{\alpha > 0 \& \beta > 0 \& m > 0 \& e > 0\}$, Solve [$t == \frac{\sqrt{m} r \sqrt{e - \frac{M^2+2 m (\alpha+t^2 \beta)}{2 m r^2}}}{\sqrt{2} (e - \beta)}$, r]] // FullSimplify

$$\left\{ \left\{ r \rightarrow -\frac{\sqrt{M^2 + 2 m \alpha + 4 t^2 (e - \beta)^2}}{\sqrt{2} \sqrt{m} \sqrt{e - \beta}} \right\}, \left\{ r \rightarrow \frac{\sqrt{M^2 + 2 m \alpha + 4 t^2 (e - \beta)^2}}{\sqrt{2} \sqrt{m} \sqrt{e - \beta}} \right\} \right\}$$

Clear[r, t];

$$r = \frac{\sqrt{M^2 + 2 m \alpha + 4 t^2 (e - \beta)^2}}{\sqrt{2} \sqrt{m} \sqrt{e - \beta}};$$

$$\frac{M \left(\pi + 2 \text{ArcCsc} \left[\sqrt{2} r \sqrt{\frac{m (e-\beta)}{M^2+2 m \alpha}} \right] \right)}{\sqrt{2} \sqrt{m} \sqrt{\frac{2 M^2}{m} + 4 \alpha}} \quad \text{// FullSimplify}$$

$$\frac{M \left(\pi - \text{ArcSec} \left[\frac{\sqrt{M^2 + 2 m \alpha + 4 t^2 (e - \beta)^2} \sqrt{\frac{m (e-\beta)}{M^2+2 m \alpha}}}{\sqrt{m} \sqrt{e - \beta}} \right] \right)}{\sqrt{m} \sqrt{\frac{M^2}{m} + 2 \alpha}}$$

(ii)

$$\begin{aligned}
 \phi_0 &= \frac{M}{\sqrt{2m}} \int_{r_{\min}}^{\infty} \frac{dr}{r^2 \sqrt{E - \left(\frac{\alpha + \beta r^2}{r^2} + \frac{M^2}{2mr^2} \right)}} \\
 &= \frac{M\pi}{2\sqrt{M^2 + 2m\alpha}} \\
 \chi &= |\pi - 2\phi_0| \\
 \chi &= \pi \left| \frac{M}{\sqrt{M^2 + 2m\alpha}} - 1 \right| \\
 &= \pi \left| \frac{1}{\sqrt{1 + \frac{2\alpha}{m\rho^2 v_\infty^2}}} - 1 \right|
 \end{aligned}$$

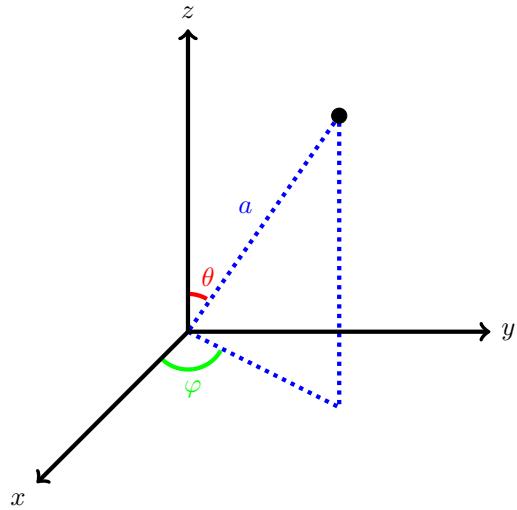
$$\frac{M}{\sqrt{2m}} \text{ Assuming } [\{\alpha > 0 \& \beta > 0 \& m > 0\}, \text{Integrate}\left[\frac{1}{q^2 \sqrt{e - \left(\frac{\alpha + \beta q^2}{q^2} + \frac{M^2}{2m q^2}\right)}}, \{q, \frac{\sqrt{-M^2 - 2m\alpha}}{\sqrt{-2e + m + 2m\beta}}, \text{Infinity}\}\right]]$$

$$\frac{M\pi}{\sqrt{2} \sqrt{m} \sqrt{\frac{2M^2}{m} + 4\alpha}} \quad \text{if } \sqrt{-m(M^2 + 2m\alpha)} \sqrt{-e + \beta} > 0$$

Problem 3

(a)

$$x^2 + y^2 + z^2 = a^2 \\ \implies r = a$$



$$x = a \sin \theta \cos \varphi$$

$$\dot{x} = a\dot{\theta} \cos \theta \cos \varphi - a\dot{\varphi} \sin \theta \sin \varphi$$

$$y = a \sin \theta \sin \varphi$$

$$\dot{y} = a\dot{\theta} \cos \theta \sin \varphi + a\dot{\varphi} \sin \theta \cos \varphi$$

$$z = a \cos \theta$$

$$\dot{z} = -a\dot{\theta} \sin \theta$$

$$\mathcal{L}_\lambda = \frac{ma^2}{2} \left(\dot{\theta}^2 + \dot{\varphi}^2 \sin^2 \theta - w^2 \sin^2 \theta \right) + \lambda (r^2 - a^2)$$

$$\varphi \text{ is cyclic} \implies M = ma^2 \sin^2 \theta \dot{\varphi} = \text{constant}$$

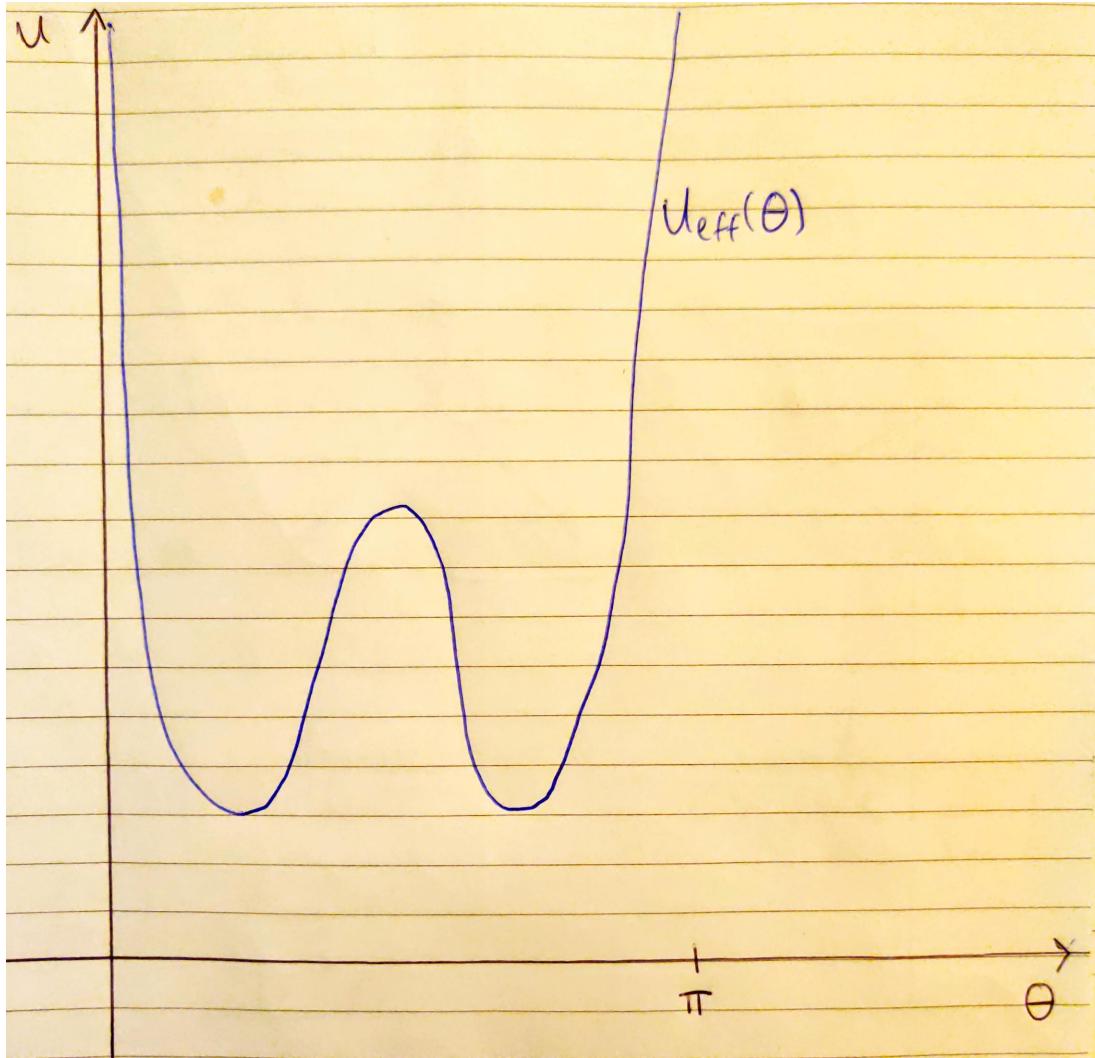
$$\implies \dot{\varphi}^2 = \frac{M^2}{m^2 a^4 \sin^4 \theta}$$

$$\implies \mathcal{L}_\lambda = \frac{ma^2 \dot{\theta}^2}{2} + \frac{M^2}{2ma^2 \sin^2 \theta} - \frac{mw^2 a^2 \sin^2 \theta}{2} + \lambda (r^2 - a^2)$$

(b)

$$U(\theta) = \frac{mw^2 a^2 \sin^2 \theta}{2}$$

$$U_{\text{eff}}(\theta) = \frac{mw^2 a^2 \sin^2 \theta}{2} + \frac{M^2}{2ma^2 \sin^2 \theta}$$



$$U_{\text{eff}}\left(\frac{\pi}{2}\right) = \frac{mw^2 a^2}{2} + \frac{M^2}{2ma^2}$$

$$U_{\text{eff}}(\rho) = \frac{mw^2 \rho^2}{2} + \frac{M^2}{2m\rho^2}$$

$$\rho_{\text{turning}} = a \sin \theta_{\text{turning}}$$

$$= \sqrt{\frac{M}{mw}}$$

$$U_{\text{eff,min}} = U_{\text{eff}}(\theta_{\text{turning}})$$

$$= Mw$$

If $Mw < E < \frac{mw^2 a^2}{2} + \frac{M^2}{2ma^2}$ then the motion is finite. The particle can move on the sphere of radius a , but only between two pairs of values for z .

If $E > \frac{mw^2a^2}{2} + \frac{M^2}{2ma^2}$ then the motion is also finite. The particle can move on the sphere of radius a , but only between two values for z .

$$Mw < E < \frac{mw^2a^2}{2} + \frac{M^2}{2ma^2} \implies \theta_1 = \arcsin\left(\sqrt{\frac{E}{a^2mw^2} - \frac{\sqrt{a^4m^2(E^2 - M^2w^2)}}{a^4m^2w^2}}\right)$$

$$\theta_2 = \arcsin\left(\sqrt{\frac{E}{a^2mw^2} + \frac{\sqrt{a^4m^2(E^2 - M^2w^2)}}{a^4m^2w^2}}\right)$$

$$\theta_3 = \pi - \arcsin\left(\sqrt{\frac{E}{a^2mw^2} + \frac{\sqrt{a^4m^2(E^2 - M^2w^2)}}{a^4m^2w^2}}\right)$$

$$\theta_4 = \pi - \arcsin\left(\sqrt{\frac{E}{a^2mw^2} - \frac{\sqrt{a^4m^2(E^2 - M^2w^2)}}{a^4m^2w^2}}\right)$$

$$E > \frac{mw^2a^2}{2} + \frac{M^2}{2ma^2} \implies \theta_1 = \arcsin\left(\sqrt{\frac{E}{a^2mw^2} - \frac{\sqrt{a^4m^2(E^2 - M^2w^2)}}{a^4m^2w^2}}\right)$$

$$\theta_2 = \pi - \arcsin\left(\sqrt{\frac{E}{a^2mw^2} - \frac{\sqrt{a^4m^2(E^2 - M^2w^2)}}{a^4m^2w^2}}\right)$$

```
Clear[e, a, m, w, M, θ, ρ];
Assuming[{m > 0 && w > 0}, Solve[D[m w^2 ρ^2/2 + M^2/(2 m ρ^2), ρ] == 0, ρ]]
{{ρ → -I Sqrt[m]/Sqrt[w]}, {ρ → -I Sqrt[m]/Sqrt[w]}, {ρ → I Sqrt[m]/Sqrt[w]}, {ρ → I Sqrt[m]/Sqrt[w]}}

Clear[e, a, m, w, M, θ, ρ];
Assuming[{m > 0 && w > 0 && a > 0 && θ < π < e}, Solve[e == m w^2 a^2 (Sin[θ])^2/2 + M^2/(2 m a^2 (Sin[θ])^2), θ]]
{{θ → -ArcSin[Sqrt[e/a^2 m w^2 - Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → π - ArcSin[Sqrt[e/a^2 m w^2 - Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → ArcSin[Sqrt[e/a^2 m w^2 - Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → π + ArcSin[Sqrt[e/a^2 m w^2 - Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → -ArcSin[Sqrt[e/a^2 m w^2 + Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → π - ArcSin[Sqrt[e/a^2 m w^2 + Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → ArcSin[Sqrt[e/a^2 m w^2 + Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}, {θ → π + ArcSin[Sqrt[e/a^2 m w^2 + Sqrt[a^4 m^2 (e^2 - M^2 w^2)]/a^4 m^2 w^2]] + 2 π c1, if c1 ∈ ℤ}}
```

(c)

$$\begin{aligned} E &= U_{\text{eff,min}} \\ &= Mw \\ 2m\pi ab &= TM \text{ for an ellipse} \\ \implies T &= \frac{2m\pi R^2}{M} \\ R &= \rho_{\text{turning}} \\ &= \sqrt{\frac{M}{mw}} \\ \implies T &= \frac{2\pi}{w} \end{aligned}$$

$$\begin{aligned} E &= Mw \\ &= \frac{m\omega^2}{2} \frac{M}{mw} \\ \implies w &= \frac{\omega}{\sqrt{2}} \\ \implies T &= \frac{2\sqrt{2}\pi}{\omega} \end{aligned}$$

For small ω , T tends to infinity. This makes sense as the time it takes for the particle to travel in a circle will increase as the angular speed of the particle decreases.

For large ω , T tends to 0. This makes sense as the time it takes for the particle to travel in a circle will decrease as the angular speed of the particle increases.