

MAU34405: Statistical Physics I

Homework 1 due 11/10/2021

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Exercise 1

$$\begin{aligned} T &= \left(\frac{\partial E}{\partial S} \right)_{V,N} & P &= - \left(\frac{\partial E}{\partial V} \right)_{T,N} & \mu &= \left(\frac{\partial E}{\partial N} \right)_{T,V} \\ T(S, V, N) &= \frac{2aS}{N} & P(S, V, N) &= \frac{2bV}{N} & \mu(S, V, N) &= -\frac{aS^2}{N^2} + \frac{bV^2}{N^2} \\ \mu &= -\frac{aS^2}{N^2} + \frac{bV^2}{N^2} \\ &= -\frac{1}{N} \left(a \frac{S^2}{N} - b \frac{V^2}{N} \right) \\ &= -\frac{E}{N} \end{aligned}$$

Exercise 2

1.

The container in question is thermally and chemically isolated, and so no particles are exchanged between the inside of the container and the outside, i.e. the number of particles N is conserved in the process. Since we are told that the gas undergoes free expansion, and that the container is thermally isolated, we know that there is no work done on or by the container by or on the surroundings, and also that there is no heat flow between the container and the surroundings, i.e. $dQ = dW = 0$. From this we can deduce that energy E is conserved, since $dE = dQ + dW = 0 + 0 = 0$.

2.

$$\begin{aligned} dS &= \left(\frac{\partial S}{\partial T} \right)_V dT + \left(\frac{\partial S}{\partial V} \right)_T dV && \text{(differential of } S(T, V)) \\ C_V \equiv \left(\frac{\partial E}{\partial T} \right)_V &= T \left(\frac{\partial S}{\partial T} \right)_V && \text{(alternative definition of specific heat capacity at constant volume)} \end{aligned}$$

$$\begin{aligned} \left(\frac{\partial S}{\partial V} \right)_T &= \left(\frac{\partial P}{\partial T} \right)_V && \text{(Maxwell's relations)} \\ &= \frac{kN}{V - bN} \end{aligned}$$

$$\implies dS = \frac{C_V}{T} dT + \frac{kN}{V - bN} dV \quad \text{(substituting)}$$

$$\begin{aligned} dE &= T dS - P dV + \sum_i \mu_i dN_i && \text{(differential form of first law)} \\ &= C_V dT + \left(\frac{kNT}{V - bN} - P \right) dV + 0 && \text{(substituting, isolated } \implies dN_i = 0) \\ &= \frac{3}{2} kN dT + a \frac{N^2}{V^2} dV && \text{(substituting given expressions)} \\ \implies E &= \frac{3}{2} kNT - a \frac{N^2}{V} + f(N) && \text{(integrating)} \end{aligned}$$

$$\begin{aligned}
dE &= dW + dQ = 0 + 0 = 0 \quad (\text{free expansion \& thermally isolated}) \\
\implies E_{\text{initial}} &= E_{\text{final}} \\
\implies \frac{3}{2} k N T_0 - a \frac{N^2}{\frac{1}{3} V_0} + f(N) &= \frac{3}{2} k N T_F - a \frac{N^2}{V_0} + f(N) \\
\implies \frac{3}{2} k N T_0 - 3a \frac{N^2}{V_0} &= \frac{3}{2} k N T_F - a \frac{N^2}{V_0} \\
\implies \textcolor{blue}{T_F} &= T_0 - \frac{4a N}{3k V_0}
\end{aligned}$$

Exercise 3

$$\begin{aligned}
dE &= T dS - P dV + \sum_i \mu_i dN_i \\
&= \frac{3AS^2}{NV} dS - \frac{AS^3}{NV^2} dV + 0 \quad (\text{assuming constant } N) \\
&= \frac{A}{N} d\left(\frac{S^3}{V}\right) \\
\implies E(S, V, N) &= \frac{AS^3}{NV} + f(N)
\end{aligned}$$

$$\begin{aligned}
E(\lambda S, \lambda V, \lambda N) &= \lambda E(S, V, N) \\
\implies \frac{\lambda^3 AS^3}{\lambda^2 NV} + f(\lambda N) &= \frac{\lambda AS^3}{NV} + \lambda f(N) \\
\implies f(\lambda N) &= \lambda f(N) \\
\implies f(N) &= \alpha N, \quad \text{for some constant } \alpha \\
\implies \textcolor{blue}{E(S, V, N)} &= \frac{AS^3}{NV} + \alpha N
\end{aligned}$$

2.

$$\begin{aligned}
T &= \frac{3AS^2}{NV}, \quad P = \frac{AS^3}{NV^2} \\
\implies V &= \frac{3AS^2}{NT}, \quad V^2 = \frac{AS^3}{NP} \\
\implies \left(\frac{3AS^2}{NT}\right)^2 &= \frac{AS^3}{NP} \\
\implies S &= \frac{NT^2}{9AP} \\
\implies V &= \frac{NT^3}{27AP^2}
\end{aligned}$$

$$\begin{aligned}
\implies d\mu &= -\frac{T^2}{9AP} dT + \frac{T^3}{27AP^2} \\
&= -\frac{1}{A} d\left(\frac{T^3}{27P}\right) \\
\implies \mu &= -\frac{T^3}{27AP} + \mu_0
\end{aligned}$$

$$\begin{aligned}
E &= TS - PV + \mu N \\
&= \frac{3AS^3}{NV} - \frac{AS^3}{NV} - \frac{NT^3}{27AP} + \mu_0 N \\
&= \frac{2AS^3}{NV} - \frac{N}{27A} \frac{27A^3 S^6}{N^3 V^3} \frac{NV^2}{AS^3} + \mu_0 N \\
&= \frac{2AS^3}{NV} - \frac{AS^3}{NV} + \mu_0 N \\
E(S, V, N) &= \frac{AS^3}{NV} + \mu_0 N \implies \alpha = \mu_0
\end{aligned}$$

Exercise 4

$$\begin{aligned}
dS &= \frac{1}{T} dE + \frac{P}{T} dV - \frac{\mu}{T} dN \\
&= \frac{1}{\sqrt{a E^{\frac{3}{2}} V^{-1} N^{-\frac{1}{2}}}} dE + \frac{2E}{V \sqrt{a E^{\frac{3}{2}} V^{-1} N^{-\frac{1}{2}}}} dV \quad (\text{assuming constant } N) \\
&= a^{-\frac{1}{2}} E^{-\frac{3}{4}} V^{\frac{1}{2}} N^{\frac{1}{4}} dE + 2a^{-\frac{1}{2}} E^{\frac{1}{4}} V^{-\frac{1}{2}} N^{\frac{1}{4}} dV \\
&= a^{-\frac{1}{2}} N^{\frac{1}{4}} d(4E^{\frac{1}{4}} V^{\frac{1}{2}}) \\
\implies S(E, V, N) &= 4a^{-\frac{1}{2}} E^{\frac{1}{4}} V^{\frac{1}{2}} N^{\frac{1}{4}} + f(N)
\end{aligned}$$

$$\begin{aligned}
S(\lambda E, \lambda V, \lambda N) &= \lambda S(E, V, N) \\
\implies 4a^{-\frac{1}{2}} \lambda^{\frac{1}{4}} E^{\frac{1}{4}} \lambda^{\frac{1}{2}} V^{\frac{1}{2}} \lambda^{\frac{1}{4}} N^{\frac{1}{4}} + f(\lambda N) &= 4\lambda a^{-\frac{1}{2}} E^{\frac{1}{4}} V^{\frac{1}{2}} N^{\frac{1}{4}} + \lambda f(N) \\
\implies f(\lambda N) &= \lambda f(N) \\
\implies f(N) &= \beta N, \quad \text{for some constant } \beta \\
\implies S(E, V, N) &= 4a^{-\frac{1}{2}} E^{\frac{1}{4}} V^{\frac{1}{2}} N^{\frac{1}{4}} + \beta N
\end{aligned}$$

Exercise 5

$$\begin{aligned}
\left. \frac{\partial J}{\partial L} \right|_T &= \frac{aT}{L_0} \left[1 + 2 \left(\frac{L_0}{L} \right)^3 \right] \\
\implies J(T, L) &= aT \int \left(\frac{1}{L_0} + \frac{2aL_0^2}{L^3} \right) dL \\
&= aT \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right) + f(T) \\
J(T, L_0) &= aT - aT + f(T) \\
&\equiv 0 \implies f(T) = 0 \\
\implies J(T, L) &= aT \left(\frac{L}{L_0} - \frac{L_0^2}{L^2} \right)
\end{aligned}$$

$$\begin{aligned}
dF &\equiv dE - T dS - S dT \\
&= J dL - S dT \quad (\text{since } dE = T dS + J dL) \\
\implies J &= \left(\frac{\partial F}{\partial L} \right)_T, \quad -S = \left(\frac{\partial F}{\partial T} \right)_L
\end{aligned}$$

$$\begin{aligned}\left(\frac{\partial J}{\partial T}\right)_L &= \left(\frac{\partial}{\partial T} \left(\frac{\partial F}{\partial L}\right)_T\right)_L \\ &= \left(\frac{\partial}{\partial L} \left(\frac{\partial F}{\partial T}\right)_L\right)_T \\ &= -\left(\frac{\partial S}{\partial L}\right)_T\end{aligned}$$

$$\begin{aligned}dE &\equiv T dS + J dL \\ \implies \left(\frac{\partial E}{\partial L}\right)_T &= T \left(\frac{\partial S}{\partial L}\right)_T + J \\ &= -T \left(\frac{\partial J}{\partial T}\right)_L + J \\ &= -T \left(a \left(\frac{L}{L_0} - \frac{L_0^2}{L^2}\right)\right) + J \\ &= -J + J = 0 \\ dE &= \left(\frac{\partial E}{\partial L}\right)_T dL + \left(\frac{\partial E}{\partial T}\right)_L dT \\ &= 0 + C_L dT \\ \implies E(T) &= C_L T + \alpha\end{aligned}$$

$$\begin{aligned}\text{Free expansion in adiabatic environment } \implies dE &= 0 \\ \implies E(T_{\text{initial}}) &= E(T_{\text{final}}) \\ \implies C_L T_{\text{initial}} + \alpha &= C_L T_{\text{final}} + \alpha \\ \implies \textcolor{blue}{T_{\text{initial}}} &= \textcolor{blue}{T_{\text{final}}}\end{aligned}$$

$$\begin{aligned}-\left(\frac{\partial S}{\partial L}\right)_T &= \left(\frac{\partial J}{\partial T}\right)_L \\ \implies S &= a \int_{L_{\text{initial}}}^{L_0} \left(\frac{L_0^2}{L^2} - \frac{L}{L_0}\right) dL \\ &= a \left(-\frac{L_0^2}{L} - \frac{L^2}{2L_0}\right) \Big|_{L_{\text{initial}}}^{L_0} \\ &= a \left(-L_0 - \frac{L_0}{2} + \frac{L_0^2}{L_{\text{initial}}} + \frac{L_{\text{initial}}^2}{2L_0}\right) \\ S &= a L_0 \left(\frac{L_{\text{initial}}^2}{2L_0^2} + \frac{L_0}{L_{\text{initial}}} - \frac{3}{2}\right)\end{aligned}$$

$$\begin{aligned}\text{Label } f(x) &\equiv \frac{x^2}{2} + \frac{1}{x} \\ f\left(\frac{L_{\text{initial}}}{L_0}\right) &= \frac{L_{\text{initial}}^2}{2L_0^2} + \frac{L_0}{L_{\text{initial}}}, \quad f(1) = \frac{3}{2} \\ \implies S &= a L_0 \left(f\left(\frac{L_{\text{initial}}}{L_0}\right) - f(1)\right)\end{aligned}$$

$$f'(1) = 1 - \frac{1}{1^2} = 0 \quad f''(1) = 1 + \frac{1}{1^3} = 3 > 0$$

Thus f has a local minimum at $x = 1$. Since $f \rightarrow \infty$ as $x \rightarrow 0$ and $x \rightarrow \infty$, $x = 1$ is a global minimum for positive x . Thus for any $L_{\text{initial}} \neq L_0$ we have that $f\left(\frac{L_{\text{initial}}}{L_0}\right) > f(1)$. Since a and L_0 are positive constants, we must also have $S = a L_0 \left(f\left(\frac{L_{\text{initial}}}{L_0}\right) - f(1)\right) > 0$, and thus the entropy increases in this process.

Exercise 6

$$\begin{aligned}
 PV^k = a &\implies P = aV^{-k} \\
 \text{Also, } P &= -\left(\frac{\partial E}{\partial V}\right)_{S,N} \\
 \implies dE|_{S,N} &= -aV^{-k}dV \\
 \implies E &= -\frac{aV^{-k+1}}{-k+1} + f_1(S, N) \\
 &= \frac{PV}{k-1} + f_2\left(\frac{S}{N}, N\right)
 \end{aligned}$$

$$\begin{aligned}
 E\left(P, \lambda V, \frac{S}{N}, \lambda N\right) &= \lambda E\left(P, V, \frac{S}{N}, N\right) \\
 \implies \frac{\lambda PV}{k-1} + f_2\left(\frac{S}{N}, \lambda N\right) &= \frac{\lambda PV}{k-1} + \lambda f_2\left(\frac{S}{N}, N\right) \\
 \implies f_2\left(\frac{S}{N}, N\right) &= N f_3\left(\frac{S}{N}\right)
 \end{aligned}$$

$$\begin{aligned}
 PV^k &= a(S, N) \\
 &= f_4\left(\frac{S}{N}, N\right) \\
 P(\lambda V)^k &= \lambda^k PV^k \\
 &= \lambda^k f_4\left(\frac{S}{N}, N\right) \\
 \implies f_4\left(\frac{S}{N}, N\right) &= N^k f_5\left(\frac{S}{N}\right) \\
 \implies f_5\left(\frac{S}{N}\right) &= \frac{PV^k}{N^k} \\
 \implies \frac{S}{N} &= f_5^{-1}\left(\frac{PV^k}{N^k}\right) \\
 \\
 \implies E &= \frac{PV}{k-1} + N f_3\left(f_5^{-1}\left(\frac{PV^k}{N^k}\right)\right) \\
 &= \frac{PV}{k-1} + N f\left(\frac{PV^k}{N^k}\right) \quad (\text{where } f \equiv f_3 \circ f_5^{-1})
 \end{aligned}$$