

MAU34406: Statistical Physics II

Homework 2 due 14/03/2022

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JS Theoretical Physics

Exercise 1

1.

$$\begin{aligned}\ln \Xi &= -\ln(1-z) + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{5}{2}}(z) \\ &= -\ln(1-e^{\beta\mu}) + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{5}{2}}(e^{\beta\mu})\end{aligned}$$

$$\begin{aligned}J &= -kT \ln \Xi \\ &= \beta^{-1} \ln(1-e^{\beta\mu}) - V(\text{deg.}) \left(\frac{2\pi m}{h^2} \right)^{\frac{3}{2}} \beta^{-\frac{5}{2}} \text{Li}_{\frac{5}{2}}(e^{\beta\mu})\end{aligned}$$

$$\begin{aligned}S &= -\frac{\partial J}{\partial T} \\ &= -\frac{\partial \beta}{\partial T} \frac{\partial J}{\partial \beta} \\ &= -\frac{\partial}{\partial T} \left(\frac{1}{kT} \right) \frac{\partial}{\partial \beta} \left[\beta^{-1} \ln(1-e^{\beta\mu}) - V(\text{deg.}) \left(\frac{2\pi m}{h^2} \right)^{\frac{3}{2}} \beta^{-\frac{5}{2}} \text{Li}_{\frac{5}{2}}(e^{\beta\mu}) \right] \\ &= \frac{1}{kT^2} \left[-\beta^{-2} \ln(1-e^{\beta\mu}) + \frac{\beta^{-1}}{1-e^{\beta\mu}} (-\mu e^{\beta\mu}) - V(\text{deg.}) \left(\frac{2\pi m}{h^2} \right)^{\frac{3}{2}} \left(-\frac{5\beta^{-\frac{7}{2}}}{2} \text{Li}_{\frac{5}{2}}(e^{\beta\mu}) + \beta^{-\frac{5}{2}} \frac{\partial e^{\beta\mu}}{\partial \beta} \frac{\partial \text{Li}_{\frac{5}{2}}(e^{\beta\mu})}{\partial e^{\beta\mu}} \right) \right] \\ &= k\beta^2 \left[-\beta^{-2} \ln(1-e^{\beta\mu}) - \beta^{-2} \frac{e^{\beta\mu} \beta \mu}{1-e^{\beta\mu}} - V(\text{deg.}) \left(\frac{2\pi m}{h^2} \right)^{\frac{3}{2}} \left(-\frac{5\beta^{-\frac{7}{2}}}{2} \text{Li}_{\frac{5}{2}}(e^{\beta\mu}) + \beta^{-\frac{7}{2}} e^{\beta\mu} \beta \mu \frac{\text{Li}_{\frac{3}{2}}(e^{\beta\mu})}{e^{\beta\mu}} \right) \right] \\ S &= -k \left[\ln(1-z) + \frac{z \ln(z)}{1-z} + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \left(\ln(z) \text{Li}_{\frac{3}{2}}(z) - \frac{5}{2} \text{Li}_{\frac{5}{2}}(z) \right) \right]\end{aligned}$$

In the case where $T \ll T_c$ we have $z \approx 1$.

$$\begin{aligned}\beta^{-\frac{3}{2}} \left(\ln(z) \text{Li}_{\frac{3}{2}}(z) - \frac{5}{2} \text{Li}_{\frac{5}{2}}(z) \right) &\approx \beta^{-\frac{3}{2}} \left(\ln(1) \text{Li}_{\frac{3}{2}}(1) - \frac{5}{2} \text{Li}_{\frac{5}{2}}(1) \right) \\ &\ll 1 \\ (\beta^{-\frac{3}{2}} \ll 1, \ln(1) = 0, \text{Li}_{\frac{3}{2}}(1) = \zeta\left(\frac{3}{2}\right) \approx 2.61, \text{Li}_{\frac{5}{2}}(1) = \zeta\left(\frac{5}{2}\right) \approx 1.34) \\ \lim_{z \rightarrow 1} \frac{z \ln(z)}{1-z} &= \lim_{z \rightarrow 1} \frac{\ln(z) + \frac{z}{z}}{-1} \\ &= -1 \\ \lim_{z \rightarrow 1} \ln(1-z) &= -\infty\end{aligned}$$

$$\implies S \approx -k \ln(1-z) \text{ for } T \ll T_c$$

In the case where $T \gg T_c$ we have $z \ll 1$.

$$\begin{aligned}
& -\ln(1-z) \ll 1 \\
& \text{Li}_r(z) = \sum_{n=1}^{\infty} \frac{z^n}{n^r} \\
& \approx z \\
& \Rightarrow \frac{5}{2} \text{Li}_{\frac{5}{2}}(z) \ll -\ln(z) \text{Li}_{\frac{3}{2}}(z) \\
& \Rightarrow S \approx -kz \ln(z) \left[\frac{1}{1-z} + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \right] \\
& \quad S \approx -kz \ln(z) V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \quad (\beta^{-\frac{3}{2}} \gg 1, \frac{1}{1-z} \approx 1)
\end{aligned}$$

2.

$$\begin{aligned}
\langle N \rangle &= z \frac{\partial \ln \Xi}{\partial z} & \langle E \rangle &= \mu \langle N \rangle - \frac{\partial \ln \Xi}{\partial \beta} \\
&= \frac{z}{1-z} + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z) & &= \frac{3V}{2\beta} (\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{5}{2}}(z) \\
S &= -k \left\{ \ln(1-z) + \left[\frac{z}{1-z} + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z) \right] \ln(z) - \frac{5\beta}{3} \frac{3V}{2\beta} (\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{5}{2}}(z) \right\} \\
& \quad S = -k \left(\ln(1-z) + \langle N \rangle \ln(z) - \frac{5\beta}{3} \langle E \rangle \right)
\end{aligned}$$

3.

$$\begin{aligned}
N &= \langle n_0 \rangle + N_{\text{exc}} \\
&= \frac{z}{1-z} + V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z) \\
\Rightarrow \langle n_0 \rangle &= \frac{z}{1-z} \text{ and } N_{\text{exc}} = V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z)
\end{aligned}$$

In the case where $T \ll T_c$ we have $\langle n_0 \rangle \approx N$.

$$\begin{aligned}
\langle n_0 \rangle &= \frac{z}{1-z} \\
\Rightarrow z &= \frac{1}{1 + \frac{1}{\langle n_0 \rangle}} \\
&\approx \frac{1}{1 + \frac{1}{N}} \\
&= 1 - \frac{1}{N} + \mathcal{O}(N^{-2}) \\
\Rightarrow S &\approx -k \ln \left(1 - 1 + \frac{1}{N} \right) \\
& \quad S \approx k \ln(N)
\end{aligned}$$

In the case where $T \ll T_c$ we have $z \ll 1$, $N_{\text{exc}} \approx N$, and $\langle n_0 \rangle \ll 1$.

$$\begin{aligned}
N_{\text{exc}} &= V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z) \\
\Rightarrow N &\approx V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} z \\
\Rightarrow S &\approx -kN \ln(z)
\end{aligned}$$

$$\begin{aligned}
\langle n_0 \rangle &= N - N_{\text{exc}} \\
&= N - V(\text{deg.}) \left(\frac{2\pi m}{h^2 \beta} \right)^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z) \\
&= N - \left[\frac{h^2}{2\pi m k} \left(\frac{N}{V(\text{deg.}) \text{Li}_{\frac{3}{2}}(1)} \right)^{\frac{2}{3}} \right]^{-\frac{3}{2}} NT^{\frac{3}{2}} \frac{\text{Li}_{\frac{3}{2}}(z)}{\text{Li}_{\frac{3}{2}}(1)} \\
&\approx N \left[1 - \left(\frac{T}{T_c} \right)^{\frac{3}{2}} \frac{z}{\zeta(\frac{3}{2})} \right] \\
&\ll 1 \\
\Rightarrow z &\approx \zeta\left(\frac{3}{2}\right) \left(\frac{T_c}{T} \right)^{\frac{3}{2}}
\end{aligned}$$

$$\begin{aligned}
\Rightarrow S &\approx -kN \ln \left[\zeta\left(\frac{3}{2}\right) \left(\frac{T_c}{T} \right)^{\frac{3}{2}} \right] \\
&= kN \left[\frac{3}{2} \ln\left(\frac{T}{T_c}\right) - \ln \zeta\left(\frac{3}{2}\right) \right] \\
S &\approx kN \left[\frac{3}{2} \ln\left(\frac{T}{T_c}\right) \right]
\end{aligned}$$

$(T \gg T_c)$

Exercise 2

$$\begin{aligned}
P &\approx \frac{2 \langle E \rangle}{3V} \\
&= \frac{3}{2} \left(\frac{2\pi mk}{h^2} \right)^{\frac{3}{2}} T^{\frac{5}{2}} \text{Li}_{\frac{5}{2}}(z) \\
0 &= \left(\frac{\partial P}{\partial T} \right)_P \\
&= \frac{3}{2} \left(\frac{2\pi mk}{h^2} \right)^{\frac{3}{2}} \left[\frac{5T^{\frac{3}{2}}}{2} \text{Li}_{\frac{5}{2}}(z) + \frac{T^{\frac{5}{2}}}{z} z \left(\frac{\partial \text{Li}_{\frac{5}{2}}(z)}{\partial z} \right)_P \left(\frac{\partial z}{\partial T} \right)_P \right] \\
\Rightarrow 0 &= \frac{5T^{\frac{3}{2}}}{2} \text{Li}_{\frac{5}{2}}(z) + T^{\frac{5}{2}} \text{Li}_{\frac{3}{2}}(z) \frac{1}{z} \left(\frac{\partial z}{\partial T} \right)_P \\
\Rightarrow \frac{1}{z} \left(\frac{\partial z}{\partial T} \right)_P &= -\frac{5}{2T} \frac{\text{Li}_{\frac{5}{2}}(z)}{\text{Li}_{\frac{3}{2}}(z)}
\end{aligned}$$

$$\begin{aligned}
N &= \langle N \rangle \\
&\approx V(\text{deg.}) \left(\frac{2\pi mk}{h^2} \right)^{\frac{3}{2}} T^{\frac{3}{2}} \text{Li}_{\frac{3}{2}}(z) \quad (T \gg T_c \Rightarrow z \ll 1) \\
0 &= \left(\frac{\partial N}{\partial T} \right)_N \\
&= \left(\frac{\partial N}{\partial T} \right)_v \\
&= V(\text{deg.}) \left(\frac{2\pi mk}{h^2} \right)^{\frac{3}{2}} \left[\frac{3T^{\frac{1}{2}}}{2} \text{Li}_{\frac{3}{2}}(z) + \frac{T^{\frac{3}{2}}}{z} z \left(\frac{\partial \text{Li}_{\frac{3}{2}}(z)}{\partial z} \right)_v \left(\frac{\partial z}{\partial T} \right)_v \right] \\
\Rightarrow 0 &= \frac{3T^{\frac{1}{2}}}{2} \text{Li}_{\frac{3}{2}}(z) + T^{\frac{3}{2}} \text{Li}_{\frac{1}{2}}(z) \frac{1}{z} \left(\frac{\partial z}{\partial T} \right)_v \\
\Rightarrow \frac{1}{z} \left(\frac{\partial z}{\partial T} \right)_v &= -\frac{3}{2T} \frac{\text{Li}_{\frac{3}{2}}(z)}{\text{Li}_{\frac{1}{2}}(z)}
\end{aligned}$$

$$\begin{aligned}
C_P &= T \left(\frac{\partial S}{\partial T} \right)_{N,P} & C_V &= T \left(\frac{\partial S}{\partial T} \right)_{N,V} \\
&= T \left(\frac{\partial S}{\partial z} \right)_N \left(\frac{\partial z}{\partial T} \right)_P & &= T \left(\frac{\partial S}{\partial z} \right)_N \left(\frac{\partial z}{\partial T} \right)_v
\end{aligned}$$

$$\begin{aligned}
\gamma &= \frac{C_P}{C_V} \\
&= \frac{T \left(\frac{\partial S}{\partial z} \right)_N \left(\frac{\partial z}{\partial T} \right)_P}{T \left(\frac{\partial S}{\partial z} \right)_N \left(\frac{\partial z}{\partial T} \right)_v} \\
&= \frac{\frac{1}{z} \left(\frac{\partial z}{\partial T} \right)_P}{\frac{1}{z} \left(\frac{\partial z}{\partial T} \right)_v} \\
&= \frac{5 \text{Li}_{\frac{5}{2}}(z) \text{Li}_{\frac{1}{2}}(z)}{3 \left(\text{Li}_{\frac{3}{2}}(z) \right)^2}
\end{aligned}$$

Exercise 3

$$\begin{aligned}
\ln \Xi &= -\ln(1-z) - (\text{deg.}) \frac{V_d (2m)^{\frac{d}{2}}}{h^d} \int_0^\infty \epsilon^{\frac{d}{2}-1} \ln(1 - ze^{-\beta\epsilon}) d\epsilon \\
N_{\text{exc}} &= z \frac{\partial \ln \Xi}{\partial z} \\
&= 0 - z (\text{deg.}) \frac{V_d (2m)^{\frac{d}{2}}}{h^d} \int_0^\infty \epsilon^{\frac{d}{2}-1} \frac{-ze^{-\beta\epsilon}}{1 - ze^{-\beta\epsilon}} d\epsilon \\
&= (\text{deg.}) \frac{V_d (2m)^{\frac{d}{2}}}{h^d} \int_0^\infty \epsilon^{\frac{d}{2}-1} \frac{ze^{-\beta\epsilon}}{1 - ze^{-\beta\epsilon}} d\epsilon \\
&= (\text{deg.}) \frac{V_d (2m)^{\frac{d}{2}}}{h^d} \int_0^\infty d(\beta\epsilon) \frac{(\beta\epsilon)^{\frac{d}{2}-1}}{z^{-1}e^{\beta\epsilon} - 1} \beta^{-\frac{d}{2}} \\
&= (\text{deg.}) \frac{V_d (2m)^{\frac{d}{2}}}{h^d} \beta^{-\frac{d}{2}} \Gamma\left(\frac{d}{2}\right) \text{Li}_{\frac{d}{2}}(z)
\end{aligned}$$

For Bose-Einstein condensation to occur we need an upper bound on the number of particles occupying excited states when $z \approx 1$. However, for $z \approx 1$ we have $N_{\text{exc}} \propto \text{Li}_{\frac{d}{2}}(1) = \zeta(\frac{d}{2})$, which does not converge for $d = 1, 2$. Thus, Bose-Einstein condensation cannot occur for $d = 1, 2$.

Exercise 4

$$\begin{aligned}
\mathcal{W}(1, p) dp &= (d-2) \frac{VS_{d-2}}{h^{d-1}} p^{d-2} dp \\
&= (d-2) \frac{VS_{d-2}}{(hc)^{d-1}} \epsilon^{d-2} d\epsilon & (\epsilon = cp) \\
&= (d-2) \frac{VS_{d-2}}{(hc)^{d-1}} (\hbar\omega)^{d-2} d(\hbar\omega) & (\epsilon = \hbar\omega) \\
\implies \mathcal{W}(1, \omega) d\omega &= (d-2) \frac{VS_{d-2}}{(2\pi c)^{d-1}} \omega^{d-2} d\omega
\end{aligned}$$

$$\begin{aligned}
E(\epsilon) d\epsilon &= \langle n(\epsilon) \rangle \epsilon \mathcal{W}(1, \epsilon) d\epsilon \\
&= \frac{d-2}{e^{\beta\epsilon} - 1} \frac{VS_{d-2}}{(hc)^{d-1}} \epsilon^{d-1} d\epsilon \\
&= \frac{d-2}{e^{\beta\hbar\omega} - 1} \frac{VS_{d-2}}{(hc)^{d-1}} (\hbar\omega)^{d-1} \hbar d\omega & (\epsilon = \hbar\omega) \\
&= \frac{d-2}{e^{\beta\hbar\omega} - 1} \frac{\hbar VS_{d-2}}{(2\pi c)^{d-1}} \omega^{d-1} d\omega
\end{aligned}$$

$$\begin{aligned}
\langle E \rangle &= \int_0^\infty E(\epsilon) d\epsilon \\
&= \frac{\hbar(d-2)VS_{d-2}}{(2\pi c)^{d-1}} \int_0^\infty \frac{\omega^{d-1}}{e^{\beta\hbar\omega} - 1} d\omega \\
&= \frac{\hbar(d-2)VS_{d-2}}{(2\pi c)^{d-1}} \int_0^\infty (\beta\hbar)^{-d} \frac{(\beta\hbar\omega)^{d-1}}{e^{\beta\hbar\omega} - 1} d(\beta\hbar\omega) \\
&= \frac{\beta^{-d}(d-2)VS_{d-2}}{(2\pi\hbar c)^{d-1}} \Gamma(d) \text{Li}_d(1) \\
&= \frac{k^d(d-2)VS_{d-2}}{(hc)^{d-1}} \Gamma(d) \zeta(d) T^d \\
\langle \epsilon \rangle &= \frac{\langle E \rangle}{V} \\
\langle \epsilon \rangle &= \frac{k^d(d-2)S_{d-2}}{(hc)^{d-1}} \Gamma(d) \zeta(d) T^d
\end{aligned}$$

$$-PV = J$$

$$\begin{aligned}
P &= -\frac{J}{V} \\
&= \frac{kT}{V} \ln \Xi \\
&= -\frac{1}{\beta V} \int_0^\infty \mathcal{W}(1, \omega) \ln(1 - e^{-\beta\hbar\omega}) d\omega \\
&= -\frac{d-2}{\beta} \frac{S_{d-2}}{(2\pi c)^{d-1}} \int_0^\infty \omega^{d-2} \ln(1 - e^{-\beta\hbar\omega}) d\omega \\
&= -\frac{d-2}{\beta} \frac{S_{d-2}}{(2\pi c)^{d-1}} \left(\frac{\omega^{d-1} \ln(1 - e^{-\beta\hbar\omega})}{d-1} \Big|_0^\infty - \int_0^\infty \frac{\omega^{d-1}}{(d-1)(1 - e^{-\beta\hbar\omega})} (-e^{-\beta\hbar\omega}(-\beta\hbar)) d\omega \right) \\
&\quad (u = \ln(1 - e^{-\beta\hbar\omega}), dv = \omega^{d-2} d\omega) \\
&= \frac{1}{d-1} \frac{\hbar(d-2)S_{d-2}}{(2\pi c)^{d-1}} \int_0^\infty \frac{\omega^{d-1}}{e^{\beta\hbar\omega} - 1} d\omega \\
&= \frac{1}{d-1} \frac{\langle E \rangle}{V} \\
P &= \frac{\langle \epsilon \rangle}{d-1}
\end{aligned}$$