

MAU34403: Quantum Mechanics I

Homework 4 due 14/10/2021

Ruaidhrí Campion
19333850
JS Theoretical Physics

Problem 1

$$\begin{aligned}\mathbb{S}^\pm &= \mathbb{S}^1 \pm i\mathbb{S}^2 \\ &= S^1 \otimes I + I \otimes S^1 \pm i(S^2 \otimes I + I \otimes S^2)\end{aligned}$$

$$\begin{aligned}[\mathbb{S}^+, \mathbb{S}^-] &= [S^1 \otimes I + I \otimes S^1 + i(S^2 \otimes I + I \otimes S^2), S^1 \otimes I + I \otimes S^1 - i(S^2 \otimes I + I \otimes S^2)] \\ &= [S^1 \otimes I + I \otimes S^1, S^1 \otimes I + I \otimes S^1] + [S^1 \otimes I + I \otimes S^1, -i(S^2 \otimes I + I \otimes S^2)] \\ &\quad + [i(S^2 \otimes I + I \otimes S^2), S^1 \otimes I + I \otimes S^1] + [i(S^2 \otimes I + I \otimes S^2), -i(S^2 \otimes I + I \otimes S^2)] \\ &= 0 + 2i[S^2 \otimes I + I \otimes S^2, S^1 \otimes I + I \otimes S^1] + 0 \\ &= 2i([S^2 \otimes I, S^1 \otimes I] + [S^2 \otimes I, I \otimes S^1] + [I \otimes S^2, S^1 \otimes I] + [I \otimes S^2, I \otimes S^1]) \\ &= 2i(S^2 S^1 \otimes I - S^1 S^2 \otimes I + S^2 \otimes S^1 - S^1 \otimes S^2 + S^1 \otimes S^2 - S^1 \otimes S^2 + I \otimes S^2 S^1 - I \otimes S^1 S^2) \\ &= 2i((S^2 S^1 - S^1 S^2) \otimes I + I \otimes (S^2 S^1 - S^1 S^2)) \\ &= -2i([S^1, S^2] \otimes I + I \otimes [S^1, S^2]) \\ &= 2\hbar(S^3 \otimes I + I \otimes S^3) \\ &= 2\hbar\mathbb{S}^3\end{aligned}$$

$$\begin{aligned}[\mathbb{S}^\pm, \mathbb{S}^3] &= [S^1 \otimes I + I \otimes S^1 \pm i(S^2 \otimes I + I \otimes S^2), S^3 \otimes I + I \otimes S^3] \\ &= [S^1 \otimes I + I \otimes S^1, S^3 \otimes I + I \otimes S^3] + [\pm i(S^2 \otimes I + I \otimes S^2), S^3 \otimes I + I \otimes S^3] \\ &= [S^1 \otimes I, S^3 \otimes I] + [S^1 \otimes I, I \otimes S^3] + [I \otimes S^1, S^3 \otimes I] + [I \otimes S^1, I \otimes S^3] \\ &\quad \pm i([S^2 \otimes I, S^3 \otimes I] + [S^2 \otimes I, I \otimes S^3] + [I \otimes S^2, S^3 \otimes I] + [I \otimes S^2, I \otimes S^3]) \\ &= S^1 S^3 \otimes I - S^3 S^1 \otimes I + S^1 \otimes S^3 - S^1 \otimes S^3 + S^3 \otimes S^1 - S^3 \otimes S^1 + I \otimes S^1 S^3 - I \otimes S^3 S^1 \\ &\quad \pm i(S^2 S^3 \otimes I - S^3 S^2 \otimes I + S^2 \otimes S^3 - S^2 \otimes S^3 + S^3 \otimes S^2 - S^3 \otimes S^2 + I \otimes S^2 S^3 - I \otimes S^3 S^2) \\ &= (S^1 S^3 - S^3 S^1) \otimes I + I \otimes (S^1 S^3 - S^3 S^1) \pm i((S^2 S^3 - S^3 S^2) \otimes I + I \otimes (S^2 S^3 - S^3 S^2)) \\ &= [S^1, S^3] \otimes I + I \otimes [S^1, S^3] \pm i([S^2, S^3] \otimes I + I \otimes [S^2, S^3]) \\ &= -i\hbar(S^2 \otimes I + I \otimes S^2) \mp \hbar(S^1 \otimes I + I \otimes S^1) \\ &= -i\hbar\mathbb{S}^2 \mp \hbar\mathbb{S}^1 \\ &= \mp \hbar\mathbb{S}^\pm\end{aligned}$$

$$[\mathbb{S}^+, \mathbb{S}^-] = 2\hbar\mathbb{S}^3 \quad [\mathbb{S}^+, \mathbb{S}^3] = -\hbar\mathbb{S}^+ \quad [\mathbb{S}^-, \mathbb{S}^3] = \hbar\mathbb{S}^-$$

$$\begin{aligned}\mathbb{S}^\pm &= S^1 \otimes I + I \otimes S^1 \pm i(S^2 \otimes I + I \otimes S^2) \\ &= \frac{\hbar}{2}(\sigma^1 \otimes I + I \otimes \sigma^1 \pm i(\sigma^2 \otimes I + I \otimes \sigma^2)) \\ \frac{2}{\hbar}\mathbb{S}^\pm &= \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \pm \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \pm \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\ &= \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \pm \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \pm \begin{pmatrix} 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & -1 & 0 \end{pmatrix} \\ &= \begin{pmatrix} 0 & 1 & 1 & 0 \\ 1 & 0 & 0 & 1 \\ 1 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 \end{pmatrix} \pm \begin{pmatrix} 0 & 1 & 1 & 0 \\ -1 & 0 & 0 & 1 \\ -1 & 0 & 0 & 1 \\ 0 & -1 & -1 & 0 \end{pmatrix}\end{aligned}$$

$$\begin{aligned}\mathbb{S}^+ &= \hbar \begin{pmatrix} 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix} & \mathbb{S}^- &= \hbar \begin{pmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 \end{pmatrix} \\ &= \hbar (|1\rangle\langle 2| + |1\rangle\langle 3| + |2\rangle\langle 4| + |3\rangle\langle 4|) & &= \hbar (|2\rangle\langle 1| + |3\rangle\langle 1| + |4\rangle\langle 2| + |4\rangle\langle 3|)\end{aligned}$$

$$\begin{aligned}|e_1\rangle &= |\uparrow\uparrow\rangle & |e_0\rangle &= \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) & |e_{-1}\rangle &= |\downarrow\downarrow\rangle \\ &= |1\rangle & &= \frac{1}{\sqrt{2}}(|2\rangle + |3\rangle) & &= |4\rangle\end{aligned}$$

$$\begin{aligned}\mathbb{S}^+ |e_1\rangle &= \hbar(0 + 0 + 0 + 0) & \mathbb{S}^+ |e_0\rangle &= \frac{\hbar}{\sqrt{2}}(|1\rangle + |1\rangle + 0 + 0) & \mathbb{S}^+ |e_{-1}\rangle &= \hbar(0 + 0 + |2\rangle + |3\rangle) \\ &= 0 & &= \hbar\sqrt{2}|e_1\rangle & &= \hbar\sqrt{2}|e_0\rangle\end{aligned}$$

$$\begin{aligned}\mathbb{S}^- |e_1\rangle &= \hbar(|2\rangle + |3\rangle + 0 + 0) & \mathbb{S}^- |e_0\rangle &= \frac{\hbar}{\sqrt{2}}(0 + 0 + |4\rangle + |4\rangle) & \mathbb{S}^- |e_{-1}\rangle &= \hbar(0 + 0 + 0 + 0) \\ &= \hbar\sqrt{2}|e_0\rangle & &= \hbar\sqrt{2}|e_{-1}\rangle & &= 0\end{aligned}$$

$$\mathbb{S}^\pm |e_{\pm 1}\rangle = 0 \quad \mathbb{S}^\pm |e_{\mp 1}\rangle = \hbar\sqrt{2}|e_0\rangle \quad \mathbb{S}^\pm |e_0\rangle = \hbar\sqrt{2}|e_{\pm 1}\rangle$$

The action of the introduced spin operators \mathbb{S}^\pm on the vectors $|e_m\rangle$ can be interpreted as a scaling of $\hbar\sqrt{2}$ and adding to or subtracting from the index m of the vector respectively, where we consider $|e_2\rangle = |e_{-2}\rangle = 0$.

Problem 2

$$\begin{aligned}|\psi\rangle &= \frac{1}{\sqrt{10}}(2|\uparrow\uparrow\rangle - |\uparrow\downarrow\rangle + 2|\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle) \\ &= \frac{1}{\sqrt{10}}\left(2|e_1\rangle + |e_{-1}\rangle + \frac{1}{\sqrt{2}}|e_0\rangle - \frac{3}{\sqrt{2}}|f\rangle\right)\end{aligned}$$

(a)

After a measurement results in s_0 , we know the system must be in a state of a combination of $|e_0\rangle$ and $|f\rangle$, as these eigenvectors correspond to the eigenvalue s_0 . Thus the new state $|\phi\rangle$ is

$$\begin{aligned}|\phi\rangle &= \frac{|e_0\rangle\langle e_0|\psi\rangle + |f\rangle\langle f|\psi\rangle}{\sqrt{|\langle e_0|\psi\rangle|^2 + |\langle f|\psi\rangle|^2}} \\ \langle e_0|\psi\rangle &= \frac{1}{\sqrt{20}} & \langle f|\psi\rangle &= -\frac{3}{\sqrt{20}} \\ \implies |\langle e_0|\psi\rangle|^2 &= \frac{1}{20} & \implies |\langle f|\psi\rangle|^2 &= \frac{9}{20}\end{aligned}$$

$$\begin{aligned}|\phi\rangle &= \frac{\frac{1}{\sqrt{20}}|e_0\rangle - \frac{3}{\sqrt{20}}|f\rangle}{\sqrt{\frac{1}{20} + \frac{9}{20}}} \\ &= \frac{1}{\sqrt{10}}|e_0\rangle - \frac{3}{\sqrt{10}}|f\rangle \\ |\phi\rangle &= \frac{1}{\sqrt{5}}(-|\uparrow\downarrow\rangle + 2|\downarrow\uparrow\rangle)\end{aligned}$$

(b)

$$|\psi\rangle \xrightarrow{E_1} |\alpha\rangle \xrightarrow{s_0} |\beta\rangle$$

$$E_1 \implies |e_1\rangle, |e_0\rangle, |e_{-1}\rangle$$

$$|\alpha\rangle = \frac{|e_1\rangle\langle e_1|\psi\rangle + |e_0\rangle\langle e_0|\psi\rangle + |e_{-1}\rangle\langle e_{-1}|\psi\rangle}{\sqrt{|\langle e_1|\psi\rangle|^2 + |\langle e_0|\psi\rangle|^2 + |\langle e_{-1}|\psi\rangle|^2}}$$

$$\langle e_1|\psi\rangle = \frac{2}{\sqrt{10}}$$

$$\implies |\langle e_1|\psi\rangle|^2 = \frac{2}{5}$$

$$\langle e_0|\psi\rangle = \frac{1}{\sqrt{20}}$$

$$\implies |\langle e_0|\psi\rangle|^2 = \frac{1}{20}$$

$$\langle e_{-1}|\psi\rangle = \frac{1}{\sqrt{10}}$$

$$\implies |\langle e_{-1}|\psi\rangle|^2 = \frac{1}{10}$$

$$\implies |\alpha\rangle = \frac{\frac{2}{\sqrt{10}}|e_1\rangle + \frac{1}{\sqrt{20}}|e_0\rangle + \frac{1}{\sqrt{10}}|e_{-1}\rangle}{\sqrt{\frac{2}{5} + \frac{1}{20} + \frac{1}{10}}}$$

$$= \sqrt{\frac{8}{11}}|e_1\rangle + \sqrt{\frac{1}{11}}|e_0\rangle + \sqrt{\frac{2}{11}}|e_{-1}\rangle$$

$$s_0 \implies |e_0\rangle, |f\rangle$$

$$\langle e_0|\alpha\rangle = \sqrt{\frac{1}{11}}$$

$$\implies |\langle e_0|\alpha\rangle|^2 = \frac{1}{11}$$

$$\langle f|\alpha\rangle = 0$$

$$|\langle f|\alpha\rangle|^2 = 0$$

$$P_1(E_1) P_2(s_0) = \left(|\langle e_1|\psi\rangle|^2 + |\langle e_0|\psi\rangle|^2 + |\langle e_{-1}|\psi\rangle|^2 \right) \left(|\langle e_0|\alpha\rangle|^2 + |\langle f|\alpha\rangle|^2 \right)$$

$$= \left(\frac{2}{5} + \frac{1}{20} + \frac{1}{10} \right) \left(\frac{1}{11} + 0 \right)$$

$$P_1(E_1) P_2(s_0) = \frac{1}{20}$$

$$|\psi\rangle \xrightarrow{s_0} |\phi\rangle \xrightarrow{E_1} |\gamma\rangle$$

$$|\phi\rangle = \frac{1}{\sqrt{10}}|e_0\rangle - \frac{3}{\sqrt{10}}|f\rangle$$

$$E_1 \implies |e_1\rangle, |e_0\rangle, |e_{-1}\rangle$$

$$\langle e_1|\phi\rangle = \langle e_{-1}|\phi\rangle = 0$$

$$\implies |\langle e_1|\phi\rangle|^2 = |\langle e_{-1}|\phi\rangle|^2 = 0$$

$$\langle e_0|\phi\rangle = \frac{1}{\sqrt{10}}$$

$$\implies |\langle e_0|\phi\rangle|^2 = \frac{1}{10}$$

$$P_1(s_0) P_2(E_1) = \left(|\langle e_0|\psi\rangle|^2 + |\langle f|\psi\rangle|^2 \right) \left(|\langle e_1|\phi\rangle|^2 + |\langle e_0|\phi\rangle|^2 + |\langle e_{-1}|\phi\rangle|^2 \right)$$

$$= \left(\frac{1}{20} + \frac{9}{20} \right) \left(0 + \frac{1}{10} + 0 \right)$$

$$P_1(s_0) P_2(E_1) = \frac{1}{20}$$

The probabilities are equal since \mathbb{S}^3 and H are compatible, and so the order in which measurements are made does not affect the probability of the combination of measurements.

Problem 3

(i)

Getting \mathbb{S}^3 , H and $|\psi\rangle$ in a usable form

$$\begin{aligned}
 \mathbb{S}^3 &= S^3 \otimes I + I \otimes S^3 \\
 &= \frac{\hbar}{2} (\sigma^3 \otimes I + I \otimes \sigma^3) && \text{(tensor product notation)} \\
 &= \frac{\hbar}{2} \left(\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} + \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \right) \\
 &= \hbar (|1\rangle\langle 1| - |4\rangle\langle 4|) && \text{(bra-ket notation)}
 \end{aligned}$$

$$\begin{aligned}
 H &= \frac{3J}{4} + \frac{J}{\hbar^2} \sum_{\alpha=1}^3 S_1^\alpha S_2^\alpha \\
 &= \frac{3J}{4} + \frac{J}{\hbar^2} \sum_{\alpha=1}^3 (S^\alpha \otimes I) (I \otimes S^\alpha) \\
 &= \frac{3J}{4} + \frac{J}{4} \sum_{\alpha=1}^3 (\sigma^\alpha \otimes \sigma^\alpha) \\
 &= \frac{J}{4} \left(3 + \sum_{\alpha=1}^3 \sigma^\alpha \otimes \sigma^\alpha \right) && \text{(tensor product notation)} \\
 &= \frac{J}{4} (3I + \sigma^1 \otimes \sigma^1 + \sigma^2 \otimes \sigma^2 + \sigma^3 \otimes \sigma^3) \\
 &= \frac{J}{4} \left(\begin{pmatrix} 3 & 0 & 0 & 0 \\ 0 & 3 & 0 & 0 \\ 0 & 0 & 3 & 0 \\ 0 & 0 & 0 & 3 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \right) \\
 &= \frac{J}{2} (2|1\rangle\langle 1| + |2\rangle\langle 2| + |2\rangle\langle 3| + |3\rangle\langle 2| + |3\rangle\langle 3| + 2|4\rangle\langle 4|) && \text{(bra-ket notation)}
 \end{aligned}$$

$$|\uparrow\uparrow\rangle = |1\rangle$$

$$|\uparrow\downarrow\rangle = |2\rangle$$

$$|\downarrow\uparrow\rangle = |3\rangle$$

$$|\downarrow\downarrow\rangle = |4\rangle$$

$$\implies |\psi\rangle = \frac{1}{\sqrt{10}} (2|1\rangle - |2\rangle + 2|3\rangle + |4\rangle)$$

Finding the anti-commutator

$$\begin{aligned}
[\mathbb{S}^3, H]_+ &= 2 \mathbb{S}^3 H && (\mathbb{S}^3 \text{ and } H \text{ commute} \implies \mathbb{S}^3 H = H \mathbb{S}^3) \\
&= 2 \frac{\hbar}{2} (\sigma^3 \otimes I + I \otimes \sigma^3) \frac{J}{4} \left(3 + \sum_{\alpha=1}^3 \sigma^\alpha \otimes \sigma^\alpha \right) && (\text{using tensor product notation}) \\
&= \frac{J\hbar}{4} \left(3 (\sigma^3 \otimes I + I \otimes \sigma^3) + \sum_{\alpha=1}^3 (\sigma^3 \sigma^\alpha \otimes \sigma^\alpha + \sigma^\alpha \otimes \sigma^3 \sigma^\alpha) \right) \\
&= \frac{J\hbar}{4} (3 (\sigma^3 \otimes I + I \otimes \sigma^3) + \sigma^3 \sigma^1 \otimes \sigma^1 + \sigma^1 \otimes \sigma^3 \sigma^1 + \sigma^3 \sigma^2 \otimes \sigma^2 + \sigma^2 \otimes \sigma^3 \sigma^2 + \sigma^3 \sigma^3 \otimes \sigma^3 + \sigma^3 \otimes \sigma^3 \sigma^3) \\
&= \frac{J\hbar}{4} (3 (\sigma^3 \otimes I + I \otimes \sigma^3) + i \sigma^2 \otimes \sigma^1 + \sigma^1 \otimes i \sigma^2 - i \sigma^1 \otimes \sigma^2 - \sigma^2 \otimes i \sigma^1 + I \otimes \sigma^3 + \sigma^3 \otimes I) \\
&= \frac{J\hbar}{4} (4 (\sigma^3 \otimes I + I \otimes \sigma^3) + 0 + 0) \\
&= 2J (S^3 \otimes I + I \otimes S^3) \\
&= 2J \mathbb{S}^3
\end{aligned}$$

Finding the expected values

$$\begin{aligned}
\langle \mathbb{S}^3 \rangle_\psi &= \langle \psi | \mathbb{S}^3 | \psi \rangle \\
&= \frac{1}{\sqrt{10}} (2 \langle 1 | - \langle 2 | + 2 \langle 3 | + \langle 4 |) \hbar (|1\rangle \langle 1| - |4\rangle \langle 4|) |\psi\rangle && (\text{using bra-ket notation}) \\
&= \frac{\hbar}{\sqrt{10}} (2 \langle 1 | - \langle 4 |) \frac{1}{\sqrt{10}} (2 |1\rangle - |2\rangle + 2 |3\rangle + |4\rangle) \\
&= \frac{\hbar}{10} (4 - 1) \\
&= \frac{3\hbar}{10}
\end{aligned}$$

$$\begin{aligned}
\langle H \rangle_\psi &= \langle \psi | H | \psi \rangle \\
&= \frac{1}{\sqrt{10}} (2 \langle 1 | - \langle 2 | + 2 \langle 3 | + \langle 4 |) \frac{J}{2} (2 |1\rangle \langle 1| + |2\rangle \langle 2| + |2\rangle \langle 3| + |3\rangle \langle 2| + |3\rangle \langle 3| + 2 |4\rangle \langle 4|) |\psi\rangle && (\text{using bra-ket notation}) \\
&= \frac{J}{2\sqrt{10}} (4 \langle 1 | - \langle 2 | - \langle 3 | + 2 \langle 2 | + 2 \langle 3 | + 2 \langle 4 |) \frac{1}{\sqrt{10}} (2 |1\rangle - |2\rangle + 2 |3\rangle + |4\rangle) \\
&= \frac{J}{20} (8 - 1 + 2 + 2) \\
&= \frac{11J}{20}
\end{aligned}$$

Finding the uncertainties

$$\begin{aligned}
\Delta(\mathbb{S}^3)^2 &= \langle s | s \rangle && (\text{where } |s\rangle \equiv (\mathbb{S}^3 - \langle \mathbb{S}^3 \rangle) |\psi\rangle) \\
&= \langle \psi | (\mathbb{S}^3 - \langle \mathbb{S}^3 \rangle) (\mathbb{S}^3 - \langle \mathbb{S}^3 \rangle) | \psi \rangle && (\text{since } (\mathbb{S}^3)^\dagger = \mathbb{S}^3) \\
&= \langle \psi | \mathbb{S}^3 \mathbb{S}^3 | \psi \rangle - 2 \langle \mathbb{S}^3 \rangle \langle \psi | \mathbb{S}^3 | \psi \rangle + \langle \mathbb{S}^3 \rangle^2 \langle \psi | \psi \rangle \\
&= \frac{\hbar}{\sqrt{10}} (2 \langle 1 | - \langle 4 |) \frac{\hbar}{\sqrt{10}} (2 |1\rangle - |4\rangle) - 2 \langle \mathbb{S}^3 \rangle^2 + \langle \mathbb{S}^3 \rangle^2 \\
&= \frac{\hbar^2}{10} (4 + 1) - \frac{9\hbar^2}{100} \\
&= \frac{41\hbar^2}{100}
\end{aligned}$$

$$\begin{aligned}
\Delta H^2 &= \langle h|h \rangle && (\text{where } |h\rangle \equiv (H - \langle H \rangle) |\psi\rangle) \\
&= \langle \psi | (H - \langle H \rangle) (H - \langle H \rangle) |\psi \rangle && (\text{since } H^\dagger = H) \\
&= \langle \psi | H H |\psi \rangle - 2\langle H \rangle \langle \psi | H |\psi \rangle + \langle H \rangle^2 \langle \psi | \psi \rangle \\
&= \frac{J}{2\sqrt{10}} (4|1\rangle + |2\rangle + |3\rangle + 2|4\rangle) \frac{J}{2\sqrt{10}} (4|1\rangle + |2\rangle + |3\rangle + 2|4\rangle) - 2\langle H \rangle^2 + \langle H \rangle^2 \\
&= \frac{J^2}{40} (16 + 1 + 1 + 4) - \frac{169J^2}{400} \\
&= \frac{99J^2}{400}
\end{aligned}$$

Checking the relation

$$\begin{aligned}
\Delta(\mathbb{S}^3)^2 \Delta H^2 &= \frac{41\hbar^2}{100} \frac{99J^2}{400} \left(\frac{1}{2} \langle [\mathbb{S}^3, H]_+ \rangle - \langle \mathbb{S}^3 \rangle \langle H \rangle \right)^2 - \frac{1}{4} \langle [\mathbb{S}^3, H] \rangle^2 = \left(\frac{1}{2} \langle 2J \mathbb{S}^3 \rangle - \frac{3\hbar}{10} \frac{11J}{20} \right)^2 - \frac{1}{4} \langle 0 \rangle^2 \\
&= \frac{4059J^2\hbar^2}{40000} \\
&\implies \Delta(\mathbb{S}^3)^2 \Delta H^2 > \left(\frac{1}{2} \langle [\mathbb{S}^3, H]_+ \rangle - \langle \mathbb{S}^3 \rangle \langle H \rangle \right)^2 - \frac{1}{4} \langle [\mathbb{S}^3, H] \rangle^2
\end{aligned}$$

(ii)

Getting \mathbb{S}^1 in a usable form

$$\begin{aligned}
\mathbb{S}^1 &= S^1 \otimes I + I \otimes S^1 \\
&= \frac{\hbar}{2} (\sigma^1 \otimes I + I \otimes \sigma^1) && (\text{tensor product notation}) \\
&= \frac{\hbar}{2} \left(\begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \right) \\
&= \frac{\hbar}{2} (|1\rangle\langle 2| + |1\rangle\langle 3| + |2\rangle\langle 1| + |2\rangle\langle 4| + |3\rangle\langle 1| + |3\rangle\langle 4| + |4\rangle\langle 2| + |4\rangle\langle 3|) && (\text{bra-ket notation})
\end{aligned}$$

Finding the commutator

$$\begin{aligned}
[\mathbb{S}^3, \mathbb{S}^1] &= \frac{\hbar^2}{4} ((\sigma^3 \otimes I + I \otimes \sigma^3) (\sigma^1 \otimes I + I \otimes \sigma^1) - (\sigma^1 \otimes I + I \otimes \sigma^1) (\sigma^3 \otimes I + I \otimes \sigma^3)) \\
&= \frac{\hbar^2}{4} (\sigma^3 \sigma^1 \otimes I + \sigma^3 \otimes \sigma^1 + \sigma^1 \otimes \sigma^3 + I \otimes \sigma^3 \sigma^1 - \sigma^1 \sigma^3 \otimes I - \sigma^1 \otimes \sigma^3 - \sigma^3 \otimes \sigma^1 - I \otimes \sigma^1 \sigma^3) \\
&= \frac{\hbar^2}{4} (i\sigma^2 \otimes I + I \otimes i\sigma^2 + i\sigma^2 \otimes I + I \otimes i\sigma^2) \\
&= \frac{i\hbar^2}{2} (\sigma^2 \otimes I + I \otimes \sigma^2) \\
&= \frac{i\hbar^2}{2} \left(\begin{pmatrix} 0 & 0 & -i & 0 \\ 0 & 0 & 0 & -i \\ i & 0 & 0 & 0 \\ 0 & i & 0 & 0 \end{pmatrix} + \begin{pmatrix} 0 & -i & 0 & 0 \\ i & 0 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 0 & 0 & i & 0 \end{pmatrix} \right) \\
&= \frac{\hbar^2}{2} \begin{pmatrix} 0 & 1 & 1 & 0 \\ -1 & 0 & 0 & 1 \\ -1 & 0 & 0 & 0 \\ 0 & -1 & -1 & 0 \end{pmatrix} \\
&= \frac{\hbar^2}{2} (|1\rangle\langle 2| + |1\rangle\langle 3| - |2\rangle\langle 1| + |2\rangle\langle 4| - |3\rangle\langle 1| + |3\rangle\langle 4| - |4\rangle\langle 2| - |4\rangle\langle 3|)
\end{aligned}$$

Finding the anti-commutator

$$\begin{aligned}
[\mathbb{S}^3, \mathbb{S}^1]_+ &= \frac{\hbar^2}{4} ((\sigma^3 \otimes I + I \otimes \sigma^3) (\sigma^1 \otimes I + I \otimes \sigma^1) + (\sigma^1 \otimes I + I \otimes \sigma^1) (\sigma^3 \otimes I + I \otimes \sigma^3)) \\
&= \frac{\hbar^2}{4} (\sigma^3 \sigma^1 \otimes I + \sigma^3 \otimes \sigma^1 + \sigma^1 \otimes \sigma^3 + I \otimes \sigma^3 \sigma^1 + \sigma^1 \sigma^3 \otimes I + \sigma^1 \otimes \sigma^3 + \sigma^3 \otimes \sigma^1 + I \otimes \sigma^1 \sigma^3) \\
&= \frac{\hbar^2}{4} (i \sigma^2 \otimes I + I \otimes i \sigma^2 - i \sigma^2 \otimes I - I \otimes i \sigma^2 + 2 \sigma^3 \otimes \sigma^1 + 2 \sigma^1 \otimes \sigma^3) \\
&= \frac{\hbar^2}{2} (\sigma^3 \otimes \sigma^1 + \sigma^1 \otimes \sigma^3) \\
&= \frac{\hbar^2}{2} \left(\begin{pmatrix} 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \end{pmatrix} \right) \\
&= \frac{\hbar^2}{2} \begin{pmatrix} 0 & 1 & 1 & 0 \\ 1 & 0 & 0 & -1 \\ 1 & 0 & 0 & -1 \\ 0 & -1 & -1 & 0 \end{pmatrix} \\
&= \frac{\hbar^2}{2} (|1\rangle\langle 2| + |1\rangle\langle 3| + |2\rangle\langle 1| - |2\rangle\langle 4| + |3\rangle\langle 1| - |3\rangle\langle 4| - |4\rangle\langle 2| - |4\rangle\langle 3|)
\end{aligned}$$

Finding the expected values

$$\begin{aligned}
\langle \mathbb{S}^1 \rangle_\psi &= \langle \psi | \mathbb{S}^1 | \psi \rangle \\
&= \frac{1}{\sqrt{10}} (2 \langle 1 | - \langle 2 | + 2 \langle 3 | + \langle 4 |) \frac{\hbar}{2} (|1\rangle\langle 2| + |1\rangle\langle 3| + |2\rangle\langle 1| + |2\rangle\langle 4| + |3\rangle\langle 1| + |3\rangle\langle 4| + |4\rangle\langle 2| + |4\rangle\langle 3|) |\psi\rangle \\
&= \frac{\hbar}{2\sqrt{10}} (2 \langle 2 | + 2 \langle 3 | - \langle 1 | - \langle 4 | + 2 \langle 1 | + 2 \langle 4 | + \langle 2 | + \langle 3 |) |\psi\rangle \\
&= \frac{\hbar}{2\sqrt{10}} (\langle 1 | + 3 \langle 2 | + 3 \langle 3 | + \langle 4 |) \frac{1}{\sqrt{10}} (2|1\rangle - |2\rangle + 2|3\rangle + |4\rangle) \\
&= \frac{\hbar}{20} (2 - 3 + 6 + 1) \\
&= \frac{3\hbar}{10}
\end{aligned}$$

$$\begin{aligned}
\langle [\mathbb{S}^3, \mathbb{S}^1] \rangle_\psi &= \langle \psi | [\mathbb{S}^3, \mathbb{S}^1] | \psi \rangle \\
&= \frac{1}{\sqrt{10}} (2 \langle 1 | - \langle 2 | + 2 \langle 3 | + \langle 4 |) \frac{\hbar^2}{2} (|1\rangle\langle 2| + |1\rangle\langle 3| - |2\rangle\langle 1| + |2\rangle\langle 4| - |3\rangle\langle 1| + |3\rangle\langle 4| - |4\rangle\langle 2| - |4\rangle\langle 3|) |\psi\rangle \\
&= \frac{\hbar^2}{2\sqrt{10}} (2 \langle 2 | + 2 \langle 3 | + \langle 1 | - \langle 4 | - 2 \langle 1 | + 2 \langle 4 | - \langle 2 | - \langle 3 |) |\psi\rangle \\
&= \frac{\hbar^2}{2\sqrt{10}} (-\langle 1 | + \langle 2 | + \langle 3 | + \langle 4 |) \frac{1}{\sqrt{10}} (2|1\rangle - |2\rangle + 2|3\rangle + |4\rangle) \\
&= \frac{\hbar^2}{20} (-2 - 1 + 2 + 1) \\
&= 0
\end{aligned}$$

$$\begin{aligned}
\langle [\mathbb{S}^3, \mathbb{S}^1]_+ \rangle &= \langle \psi | [\mathbb{S}^3, \mathbb{S}^1]_+ | \psi \rangle \\
&= \frac{1}{\sqrt{10}} (2 \langle 1| - \langle 2| + 2 \langle 3| + \langle 4|) \frac{\hbar^2}{2} (|1\rangle \langle 2| + |1\rangle \langle 3| + |2\rangle \langle 1| - |2\rangle \langle 4| + |3\rangle \langle 1| - |3\rangle \langle 4| - |4\rangle \langle 2| - |4\rangle \langle 3|) |\psi\rangle \\
&= \frac{\hbar^2}{2\sqrt{10}} (2 \langle 2| + 2 \langle 3| - \langle 1| + \langle 4| + 2 \langle 1| - 2 \langle 4| - \langle 2| - \langle 3|) |\psi\rangle \\
&= \frac{\hbar^2}{2\sqrt{10}} (\langle 1| + \langle 2| + \langle 3| - \langle 4|) \frac{1}{\sqrt{10}} (2|1\rangle - |2\rangle + 2|3\rangle + |4\rangle) \\
&= \frac{\hbar^2}{20} (2 - 1 + 2 - 1) \\
&= \frac{\hbar^2}{10}
\end{aligned}$$

Finding the uncertainties

$$\begin{aligned}
\Delta(\mathbb{S}^1)^2 &= \langle s | s \rangle && \text{(where } |s\rangle \equiv (\mathbb{S}^1 - \langle \mathbb{S}^1 \rangle) |\psi\rangle) \\
&= \langle \psi | (\mathbb{S}^1 - \langle \mathbb{S}^1 \rangle) (\mathbb{S}^1 - \langle \mathbb{S}^1 \rangle) | \psi \rangle && \text{(since } (\mathbb{S}^1)^\dagger = \mathbb{S}^3) \\
&= \langle \psi | \mathbb{S}^1 \mathbb{S}^1 | \psi \rangle - 2\langle \mathbb{S}^1 \rangle \langle \psi | \mathbb{S}^1 | \psi \rangle + \langle \mathbb{S}^1 \rangle^2 \langle \psi | \psi \rangle \\
&= \frac{\hbar}{2\sqrt{10}} (\langle 1| + 3 \langle 2| + 3 \langle 3| + \langle 4|) \frac{\hbar}{2\sqrt{10}} (|1\rangle + 3|2\rangle + 3|3\rangle + |4\rangle) - 2\langle \mathbb{S}^1 \rangle^2 + \langle \mathbb{S}^1 \rangle^2 \\
&= \frac{\hbar^2}{40} (1 + 9 + 9 + 1) - \frac{9\hbar^2}{100} \\
&= \frac{41\hbar^2}{100}
\end{aligned}$$

Checking the relation

$$\begin{aligned}
\Delta(\mathbb{S}^3)^2 \Delta(\mathbb{S}^1)^2 &= \frac{41\hbar^2}{100} \frac{41\hbar^2}{100} & \left(\frac{1}{2} \langle [\mathbb{S}^3, H]_+ \rangle - \langle \mathbb{S}^3 \rangle \langle H \rangle \right)^2 - \frac{1}{4} \langle [\mathbb{S}^3, H] \rangle^2 = \left(\frac{1}{2} \frac{\hbar^2}{10} - \frac{3\hbar}{10} \frac{3\hbar}{10} \right)^2 - \frac{0}{4} \\
&= \frac{1681\hbar^4}{10000} & & = \frac{\hbar^4}{625} \\
\implies \Delta(\mathbb{S}^3)^2 \Delta H^2 &> \left(\frac{1}{2} \langle [\mathbb{S}^3, H]_+ \rangle - \langle \mathbb{S}^3 \rangle \langle H \rangle \right)^2 - \frac{1}{4} \langle [\mathbb{S}^3, H] \rangle^2
\end{aligned}$$