

MAU34404: Quantum Mechanics II

Homework 1 due 11/02/2022

Ruaidhrí Campion
19333850
JS Theoretical Physics

Problem 1

1.

$$\begin{aligned}\delta H \Psi_0 &= \sum_n \langle \Psi_n | \delta H | \Psi_0 \rangle \Psi_n \\ \langle \Psi_n | \delta H | \Psi_0 \rangle &= \int_{-\infty}^{\infty} dx \frac{1}{\sqrt{2^n \pi n!}} H_n(x \sqrt{\alpha}) \Psi_0(x) \lambda x^4 \Psi_0(x) \\ &= \frac{\lambda \sqrt{\alpha}}{\sqrt{2^n \pi n!}} \int_{-\infty}^{\infty} dx H_n(x \sqrt{\alpha}) x^4 e^{-\alpha x^2}\end{aligned}$$

If n is odd, then the resulting Hermite polynomial will be a polynomial only in odd orders of x . This will result in a sum of integrals of odd functions on infinite bounds, which will vanish. Thus the only non-vanishing cases are when n is even. Since Ψ_4 can be expressed as Ψ_0 times a polynomial in x of degree 4, then $\delta H \Psi_0 = \lambda x^4 \Psi_0$ can be found by rearranging and scaling this expression of Ψ_4 . Thus the only non-vanishing cases are $n = 0, 2, 4$. We can find the coefficients of $\Psi_{0,2,4}$ by the method above or by directly integrating, the latter of which is done below.

$$\begin{aligned}\langle \Psi_0 | \delta H | \Psi_0 \rangle &= \frac{\lambda \sqrt{\alpha}}{\sqrt{2^0 \pi 0!}} \int_{-\infty}^{\infty} dx H_0(x \sqrt{\alpha}) x^4 e^{-\alpha x^2} \\ &= \lambda \sqrt{\frac{\alpha}{\pi}} \frac{\partial^2}{\partial \alpha^2} \left(\int_{-\infty}^{\infty} dx e^{-\alpha x^2} \right) \\ &= \lambda \sqrt{\frac{\alpha}{\pi}} \frac{\partial^2}{\partial \alpha^2} \left(\sqrt{\frac{\pi}{\alpha}} \right) \\ &= \lambda \sqrt{\frac{\alpha}{\pi}} \sqrt{\pi} \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \alpha^{-\frac{5}{2}} \\ &= \frac{3}{4} \frac{\lambda}{\alpha^2}\end{aligned}$$

$$\begin{aligned}\langle \Psi_2 | \delta H | \Psi_0 \rangle &= \frac{\lambda \sqrt{\alpha}}{\sqrt{2^2 \pi 2!}} \int_{-\infty}^{\infty} dx H_2(x \sqrt{\alpha}) x^4 e^{-\alpha x^2} \\ &= \frac{\lambda}{2} \sqrt{\frac{\alpha}{2\pi}} \int_{-\infty}^{\infty} dx (4\alpha x^2 - 2) x^4 e^{-\alpha x^2} \\ &= \lambda \sqrt{\frac{\alpha}{2\pi}} \left(-2\alpha \frac{\partial^3}{\partial \alpha^3} \left(\int_{-\infty}^{\infty} dx e^{-\alpha x^2} \right) - \frac{\partial^2}{\partial \alpha^2} \left(\int_{-\infty}^{\infty} dx e^{-\alpha x^2} \right) \right) \\ &= \lambda \sqrt{\frac{\alpha}{2\pi}} \left(-2\alpha \frac{\partial^3}{\partial \alpha^3} \left(\sqrt{\frac{\pi}{\alpha}} \right) - \frac{\partial^2}{\partial \alpha^2} \left(\sqrt{\frac{\pi}{\alpha}} \right) \right) \\ &= \lambda \sqrt{\frac{\alpha}{2\pi}} \left(-2\alpha \sqrt{\pi} \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \left(-\frac{5}{2} \right) \alpha^{-\frac{7}{2}} - \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \alpha^{-\frac{5}{2}} \right) \\ &= \frac{3}{\sqrt{2}} \frac{\lambda}{\alpha^2}\end{aligned}$$

$$\begin{aligned}
\langle \Psi_4 | \delta H | \Psi_0 \rangle &= \frac{\lambda \sqrt{\alpha}}{\sqrt{2^4 \pi} 4!} \int_{-\infty}^{\infty} dx H_4(x \sqrt{\alpha}) x^4 e^{-\alpha x^2} \\
&= \frac{\lambda}{48} \sqrt{\frac{6\alpha}{\pi}} \int_{-\infty}^{\infty} dx (16\alpha^2 x^4 - 48\alpha x^2 + 12) x^4 e^{-\alpha x^2} \\
&= \frac{\lambda}{12} \sqrt{\frac{6\alpha}{\pi}} \left(4\alpha^2 \frac{\partial^4}{\partial \alpha^4} \left(\int_{-\infty}^{\infty} dx e^{-\alpha x^2} \right) + 12\alpha \frac{\partial^3}{\partial \alpha^3} \left(\int_{-\infty}^{\infty} dx e^{-\alpha x^2} \right) + 3 \frac{\partial^2}{\partial \alpha^2} \left(\int_{-\infty}^{\infty} dx e^{-\alpha x^2} \right) \right) \\
&= \frac{\lambda}{12} \sqrt{\frac{6\alpha}{\pi}} \left(4\alpha^2 \frac{\partial^4}{\partial \alpha^4} \left(\sqrt{\frac{\pi}{\alpha}} \right) + 12\alpha \frac{\partial^3}{\partial \alpha^3} \left(\sqrt{\frac{\pi}{\alpha}} \right) + 3 \frac{\partial^2}{\partial \alpha^2} \left(\sqrt{\frac{\pi}{\alpha}} \right) \right) \\
&= \frac{\lambda}{12} \sqrt{\frac{6\alpha}{\pi}} \left(4\alpha^2 \sqrt{\pi} \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \left(-\frac{5}{2} \right) \left(-\frac{7}{2} \right) \alpha^{-\frac{9}{2}} \right. \\
&\quad \left. + 12\alpha \sqrt{\pi} \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \left(-\frac{5}{2} \right) \alpha^{-\frac{7}{2}} + 3 \sqrt{\pi} \left(-\frac{1}{2} \right) \left(-\frac{3}{2} \right) \alpha^{-\frac{5}{2}} \right) \\
&= \frac{\sqrt{6}}{2} \frac{\lambda}{\alpha^2}
\end{aligned}$$

$$\Rightarrow \delta H \Psi_0 = \frac{\lambda}{\alpha^2} \left(\frac{3}{4} \Psi_0 + \frac{3}{\sqrt{2}} \Psi_2 + \frac{\sqrt{6}}{2} \Psi_4 \right)$$

2.

$$\begin{aligned}
\delta_1 \Psi_0 &= \sum_{n \neq 0} \frac{\langle \Psi_n | \delta H | \Psi_0 \rangle}{E_0 - E_n} \Psi_n \\
&= \sum_{n=2,4} \frac{\langle \Psi_n | \delta H | \Psi_0 \rangle}{-n \hbar \omega} \Psi_n \\
&= -\frac{\lambda}{\hbar \omega \alpha^2} \left(\frac{3}{2\sqrt{2}} \Psi_2 + \frac{\sqrt{6}}{8} \Psi_4 \right) \\
\delta_1 \Psi_0 &= -\frac{\lambda}{2 \hbar \omega \alpha^2} \left(\frac{3}{\sqrt{2}} \Psi_2 + \frac{\sqrt{6}}{4} \Psi_4 \right)
\end{aligned}$$

3.

$$\begin{aligned}
\delta_2 E_0 &= \sum_{n \neq 0} \frac{|\langle \Psi_n | \delta_1 H | \Psi_0 \rangle|^2}{E_0 - E_n} + \langle \Psi_0 | \delta_2 H | \Psi_0 \rangle \\
\delta H &= \lambda x^4 \Rightarrow \begin{cases} \delta_1 H = \delta H \\ \delta_2 H = 0 \end{cases} \\
\Rightarrow \delta_2 E_0 &= \sum_{n=2,4} \frac{|\langle \Psi_n | \delta H | \Psi_0 \rangle|^2}{-n \hbar \omega} \\
&= -\frac{1}{2 \hbar \omega} \frac{18 \lambda^2}{4 \alpha^4} - \frac{1}{4 \hbar \omega} \frac{6 \lambda^2}{4 \alpha^4} \\
&= -\frac{21 \lambda^2}{8 \hbar \omega \alpha^4}
\end{aligned}$$

The ratio $\delta_1 E_0 : \delta_2 E_0$ between the first and second order changes in E_0 is $\frac{3\lambda}{4\alpha^2} : -\frac{21\lambda^2}{8\hbar\omega\alpha^4} = 1 : -\frac{7\lambda}{2\hbar\omega\alpha^2}$.