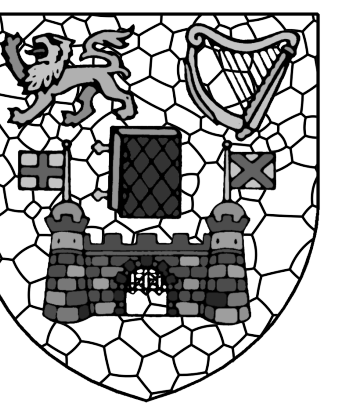


2D foams above the jamming transition: Deformation matters

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Simulations of 2D foams above the jamming point suggest a different form for the variation of the average contact number with packing fraction from that in soft disk systems.

The soft disk simulation

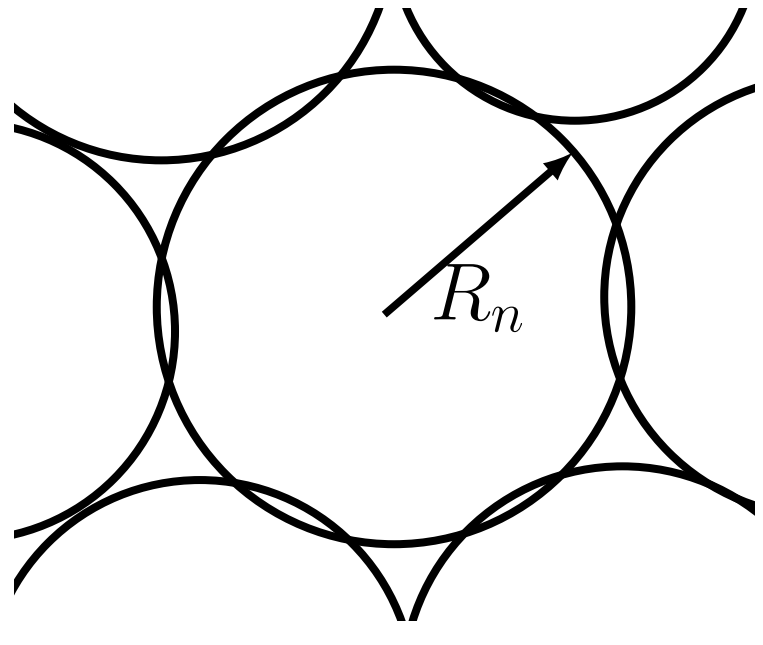


Fig. 1: A soft disk with corresponding overlaps.

Interaction and minimisation:

- **Random** initial configuration at packing fraction ϕ
- Interaction depends on overlap δ_{ij} :

$$U_{\text{soft}} = \frac{1}{2} \sum_{i < j}^N \delta_{ij}^2$$

- **Energy minimisation: Conjugate Gradient**

Deformation:

- Disks do **not** deform
- No cell area conservation

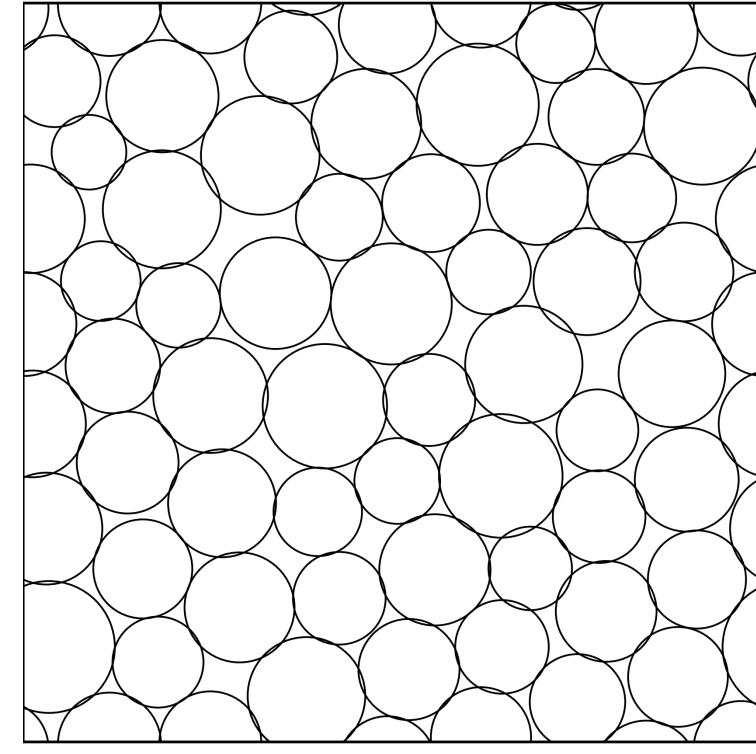


Fig. 2: Soft disk packing at $\phi = 0.90$.

The Plat simulation for 2D foam

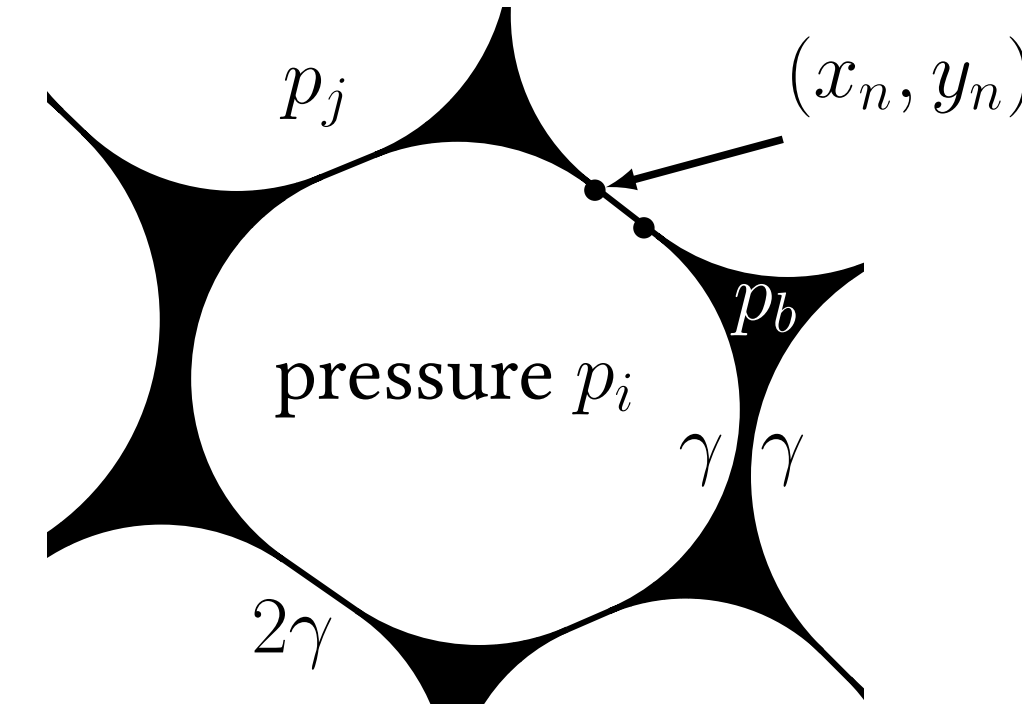


Fig. 3: A deformed bubble in Plat [1, 2].

Variables:

- Vertex coordinates (x_n, y_n)
- Cell and Plateau border pressures

Constraints for equilibrium:

- Const. cell area
- Circular arcs meet tangentially at a ver-

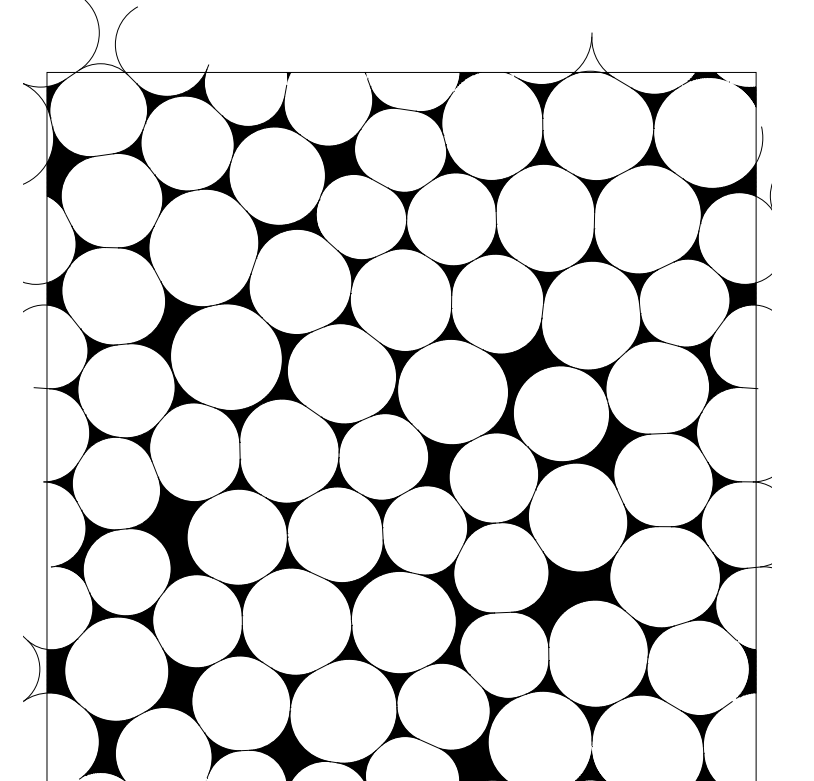


Fig. 4: Plat foam at $\phi = 0.90$.

tex (x_n, y_n)

- **Laplace-Young law** is fulfilled for cells and Plateau borders

Deformation:

- Bubbles change shape upon contact

The Morse-Witten theory for 2D bubbles

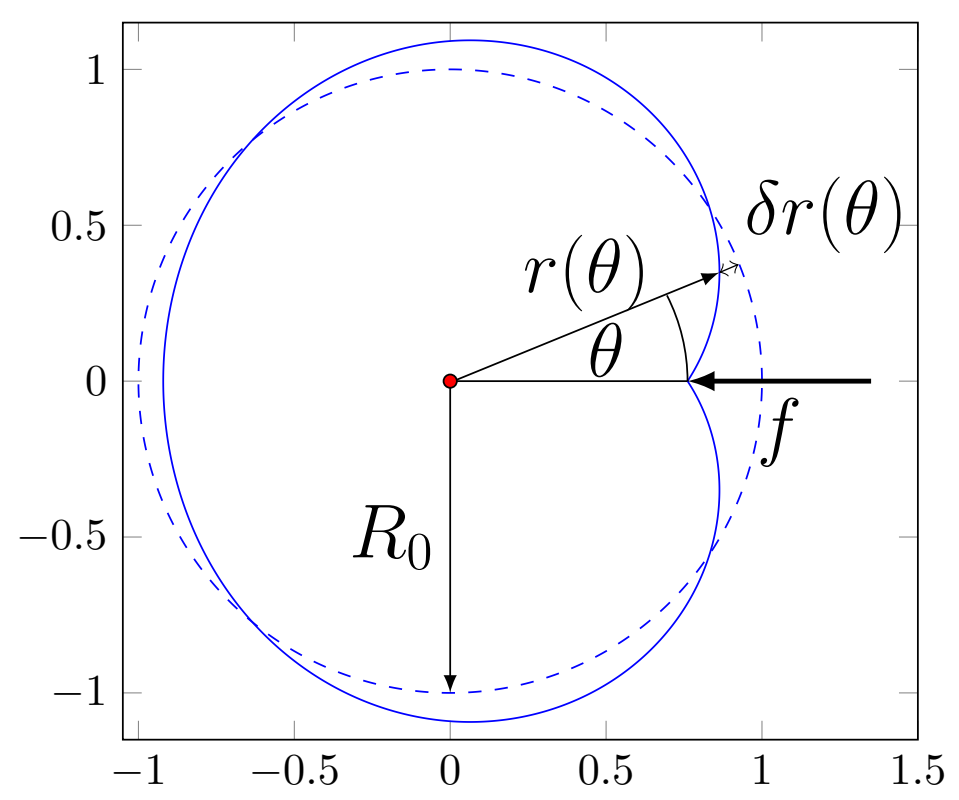


Fig. 5: Deformation of a 2D bubble with point force.

$$r(\theta) = R_0(1 + \delta r(\theta))$$

D. Weaire, R. Höhler, and S. Hutzler. *Bubble-bubble interactions in a 2d foam, close to the wet limit*. Adv. Colloid Interface Sci. accepted, (2017). <https://dx.doi.org/10.1016/j.cis.2017.07.004>.

What is the shape $\delta r(\theta)$ of a **single, incompressible 2D bubble**, subject to a single contact force f ?

Linear differential equation for the $\delta r(\theta)$:

$$-\left(\frac{d^2}{d\theta^2}\right) \delta r(\theta) = A + \frac{f}{\pi} \cos \theta$$

Solution for $\delta r(\theta)$:

$$\delta r(\theta) = \frac{f}{2\pi} g(\theta), \quad g(\theta) = (\pi - \theta) \sin \theta - \frac{\cos \theta}{2} - 1$$

- δr **linear** in $f \Rightarrow$ superposition for multiple f

The Morse-Witten (MW) simulation model

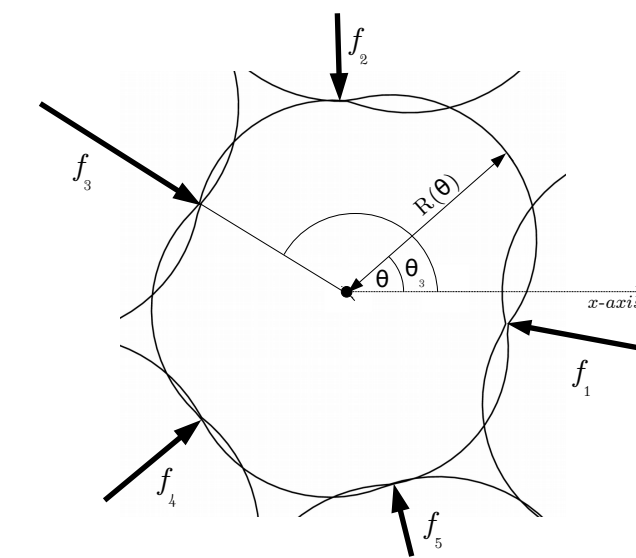


Fig. 6: Bubble deformation with multiple forces.

Shape for multiple forces:

$$\delta r(\theta) = \frac{1}{2\pi} \sum_c f_c g(\theta_c - \theta)$$

Equilibration procedure

1. Identify contacts, update δr
2. Calc. forces from δr , inverting Eq.
3. Balance forces at contacts

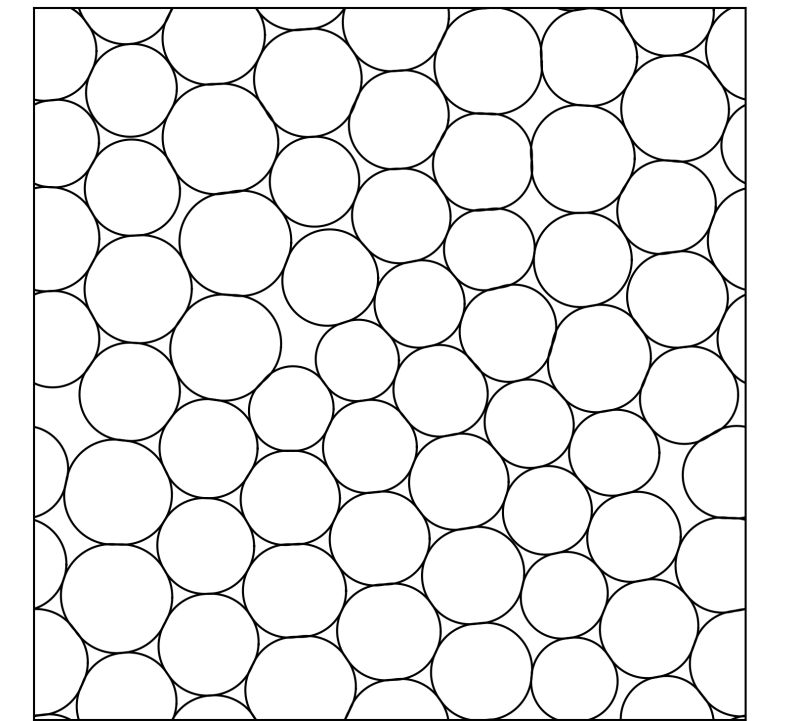


Fig. 7: MW foam at $\phi = 0.90$.

4. Move bubble centers in force direction
5. Recalculate δr with Eq.

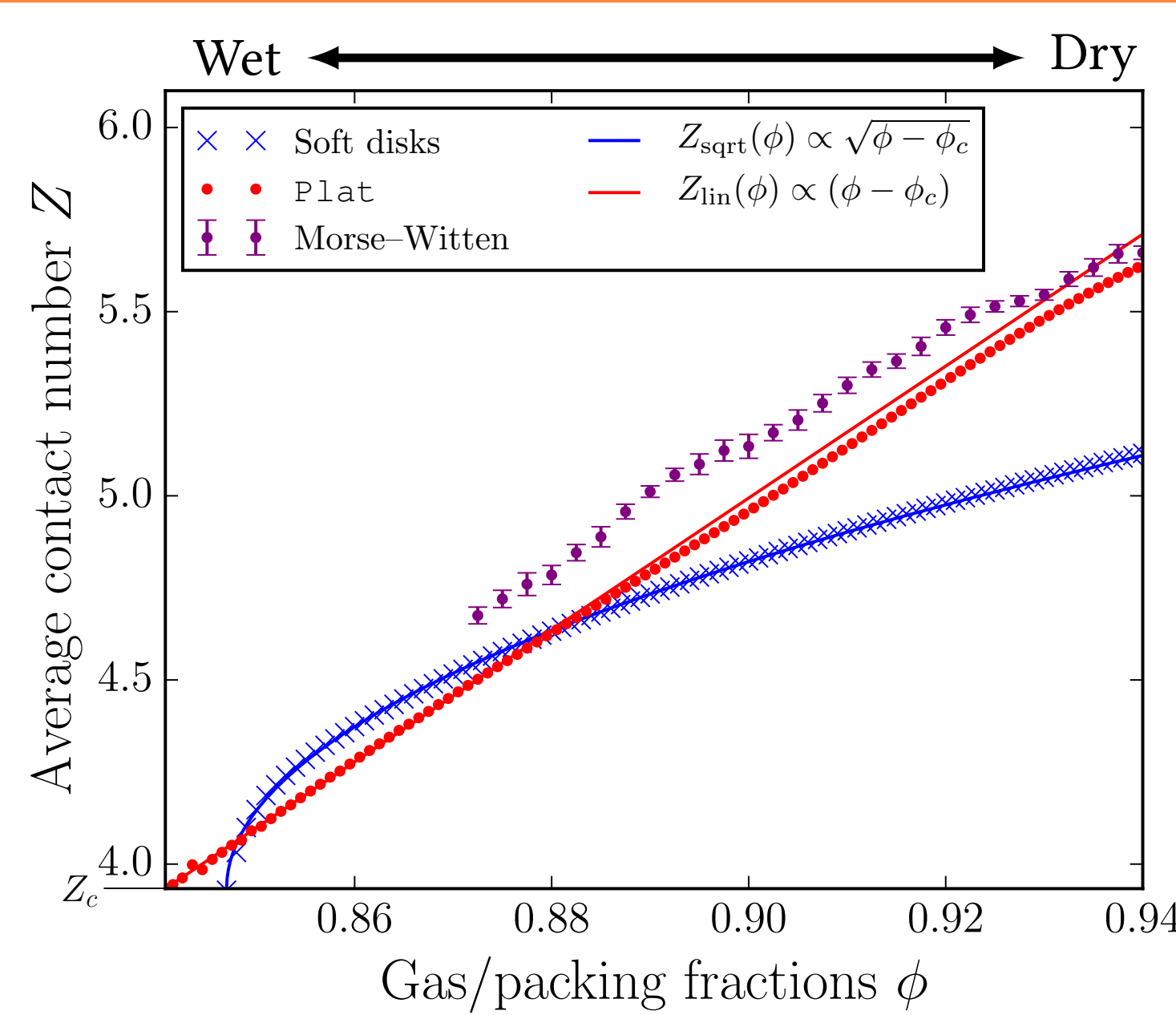
Constraints for equilibrium

- $\delta r(\theta)$ consistent with f_i
- force balance at films
- force balance on each bubble

The average contact number with increasing packing fraction $Z(\phi)$

Soft disk simulation:

- $Z - Z_c \propto \sqrt{\phi - \phi_c}$
- Well known behaviour in soft disk systems [11, 9, 13]
- Square-root scaling also observed in experimental data for photo-elastic disks [7]
- Power law relationship $Z(\phi)$ for a confined bubble raft experiment [6], *however*:
 - Identification of contacting bubbles?
 - Definition of liquid fraction?



Plat simulation:

- $Z - Z_c \propto \phi - \phi_c$
- Periodic system: $Z_c = 4(1 - 1/N)$
- data averaged over 10,000 simulations of $N = 60$ bubble systems
- hints of this in earlier simulations (Plat [1] and lattice gas model [12])

Morse-Witten (MW) model:

- Results are *preliminary*
- $Z - Z_c \approx \phi - \phi_c$
- data averaged over 5 simulations of $N = 100$ bubble systems

The separation between bubbles at the critical packing fraction

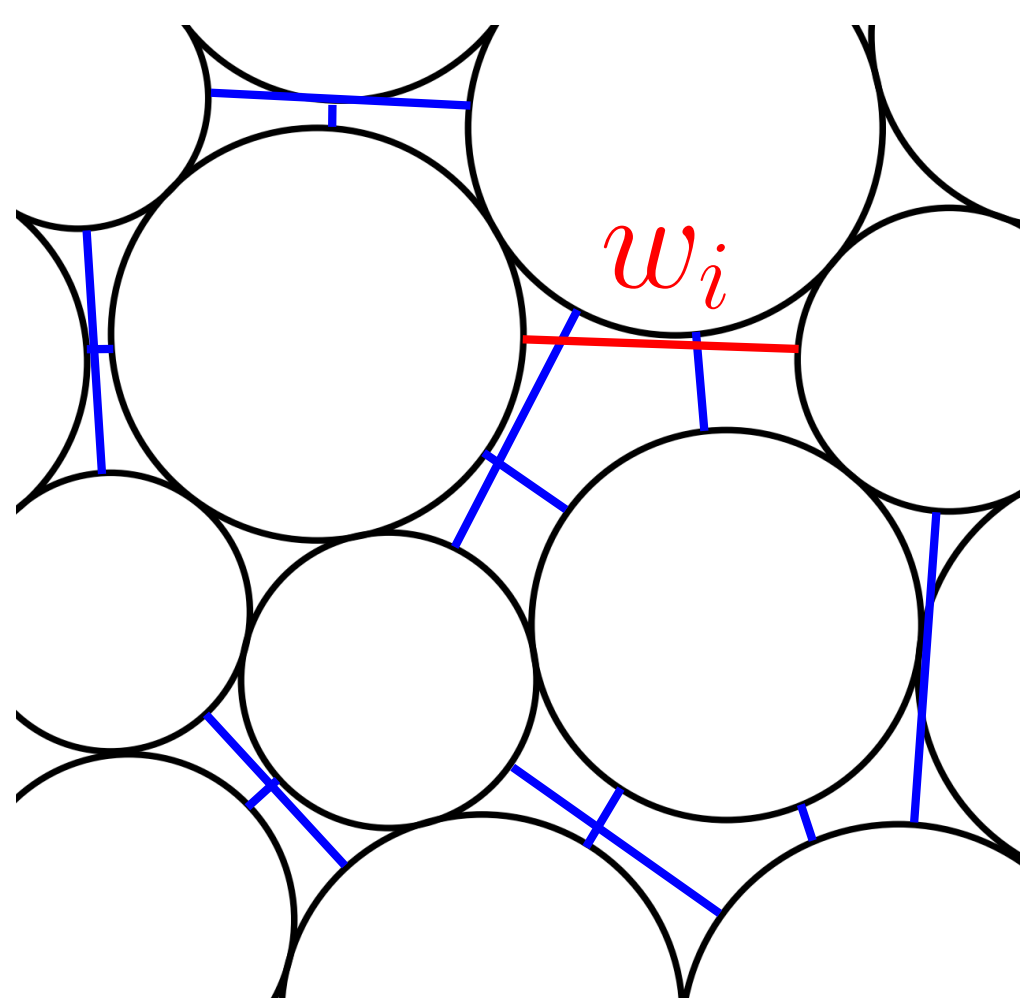


Fig. 8: Separations w_i between different disks/bubbles.

Is $Z(\phi)$ determined by the distribution separation between bubbles?

Crude argument [10]: For affine compressions at ϕ_c , $Z(\phi) - Z_c$ is given by the radial integral over the distribution of separation $f(w)$ in the limit $\epsilon \approx w/D \rightarrow 0$

$$Z(\phi) - Z_c = 2\pi \int_0^{\epsilon D} dw f(w)(w + D), \quad \text{compression: } \epsilon = \frac{\phi - \phi_c}{2\phi_c}$$

ϕ_c : critical packing fraction. D : average diameter.

Soft disks:

$$f(w) \propto \left(\frac{w}{D}\right)^{-\frac{1}{2}} \Rightarrow Z - Z_c \propto \sqrt{\phi - \phi_c}$$

2D foams:

$$f(w) = \text{const.} \Rightarrow Z - Z_c \propto \phi - \phi_c$$

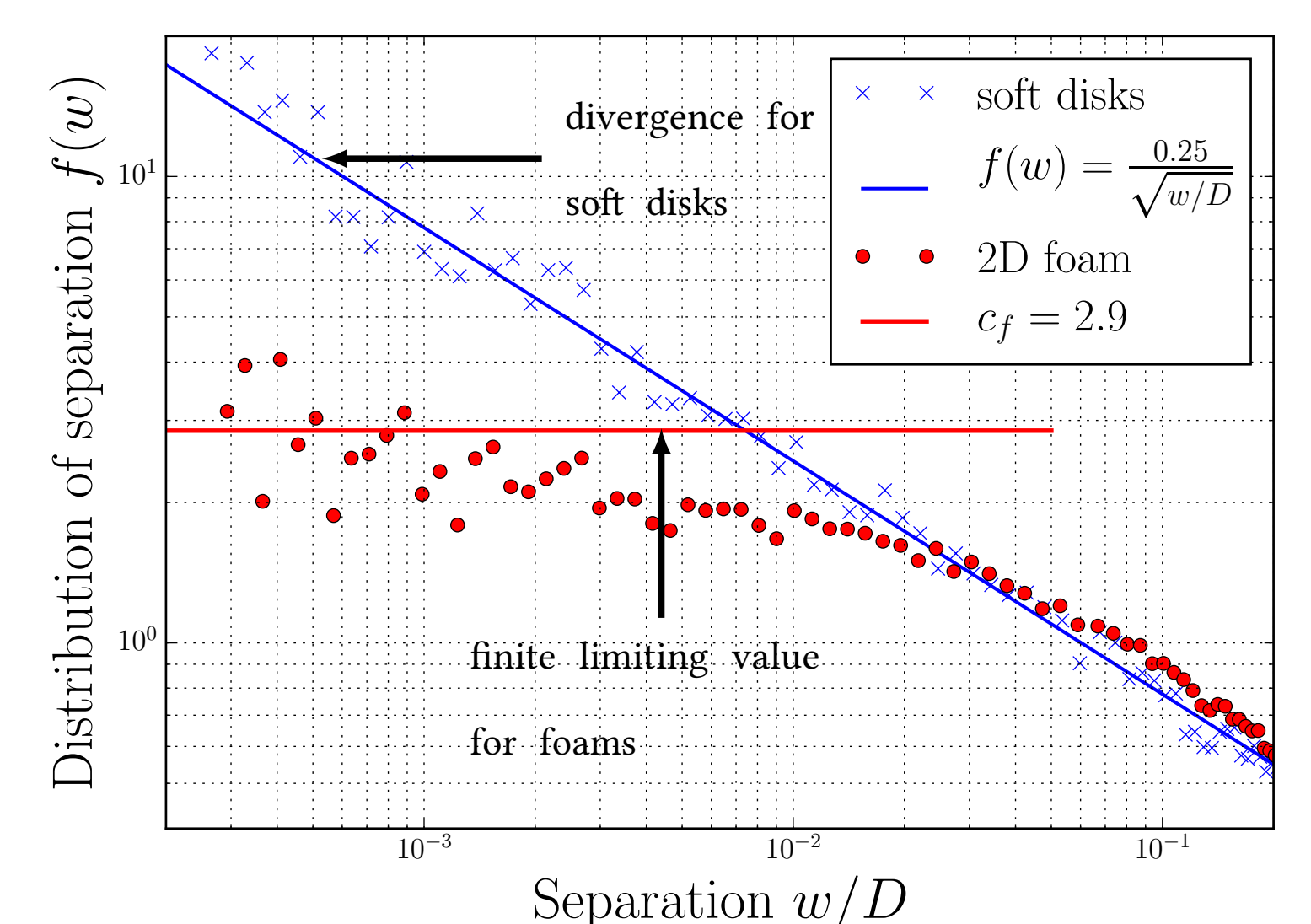


Fig. 9: Distribution of separation.

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