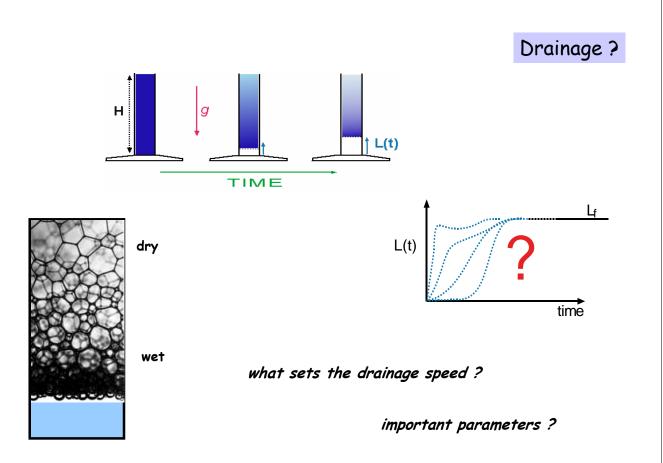
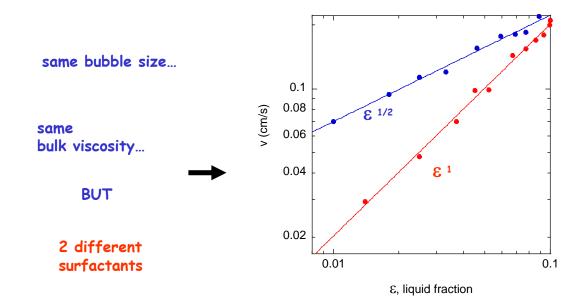
FOAM DRAINAGE - part 1

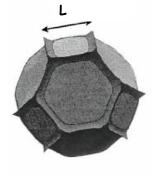




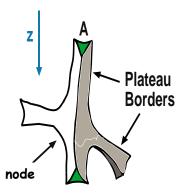
$$\rightarrow$$
 V = a ε^{b}

what sets a and b?

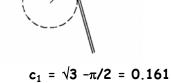
Geometry and liquid content (1)







A: Area of a Plateau border $A = c_1 r$



Volume of a Plateau border

$$V_{PB} \sim r^2 L$$

Volume of a node

$$V_{node} \sim r^3$$

$$V_{node} = c_n r^3$$
 With $c_n = 0.3$

With
$$c_n = 0.3$$

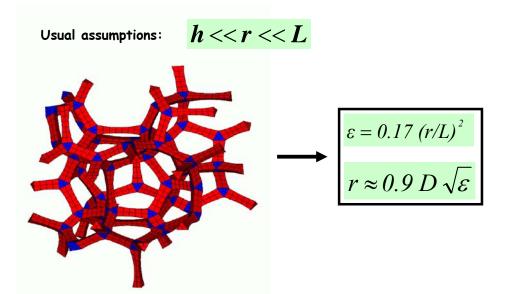
Volume of a film

$$V_{film} \sim hL^2$$

Geometry and liquid content (2)

Putting all together, and normalizing by the bubble gas volume $\sim L^3$:

$$\varepsilon = 0.17 (r/L)^2 + 0.2 (r/L)^3 + 1.2 h/L$$
 $c_2 = 0.17$ $c_3 = 0.2$



Drainage: like a porous material problem, but with a variable pore size, $r(\varepsilon)$...

Liquid velocity: scaling concepts

Porous media : $G = -\nabla p + \rho g = \mu v / k(\varepsilon)$

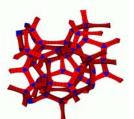
in the steady state, constant liquid fraction : $\nabla p = 0$

Power balance:

$$\rho g (r^2 L) v \sim \frac{\mu v^2}{r^2} (r^2 L)$$

dissipation in the PBs

$$\rho g (r^2 L) v \sim \frac{\mu v^2}{r^2} (r^2 L) \qquad v = K_c \frac{\rho g r^2}{\mu} = K_c \frac{\rho g L^2}{\mu} \varepsilon$$



$$\rho g (r^{2}L)v \sim \frac{\mu v^{2}}{r^{2}}(r^{3}) \qquad v = K_{n} \frac{\rho g r L}{\mu} = K_{n} \frac{\rho g L^{2}}{\mu} \sqrt{\varepsilon}$$
Permeability: K_{n}

Permeability : K_c

PB resistance :

 $R_c = 1 / (3c_2 K_c) \approx 2 / K_c$

dissipation in the nodes

resistance: $R_n = 2c_1/(3c_2^{1/2} c_n K_n) \approx 0.9/K_n$

(Koehler et al, Langmuir, 2000)

Experimental evidences, but further needs...

Experimental evidence of the different regimes :

$$v = K_c \frac{\rho g r^2}{\mu} = K_c \frac{\rho g L^2}{\mu} \varepsilon$$

$$v = K_n \frac{\rho g \, rL}{\mu} = K_n \frac{\rho g \, L^2}{\mu} \, \sqrt{\varepsilon}$$

which experimental conditions set the drainage regime?

It depends on the bubble size

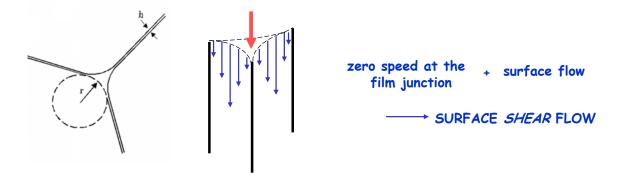
It depends on the nature of the surfactant

which values for the Kc (or Rc) and Kn (or Rn)?

Within one regime, they are actually not constant values!

Coupling bulk and surface flow...

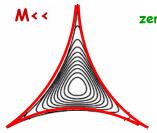
First, considering the flow inside a Plateau border (PB):



M: interfacial « mobility » parameter (Leonard and Lemlich, 1965)

$$M = \frac{\mu r}{\mu_s} = 0.9 \frac{\mu D\sqrt{\varepsilon}}{\mu_s}$$
 with μ_s = surface shear viscosity

Plateau border resistance and interfacial mobility

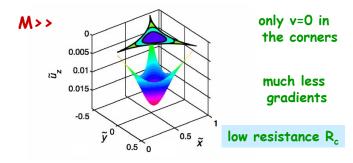


zero speed on the walls

high gradients

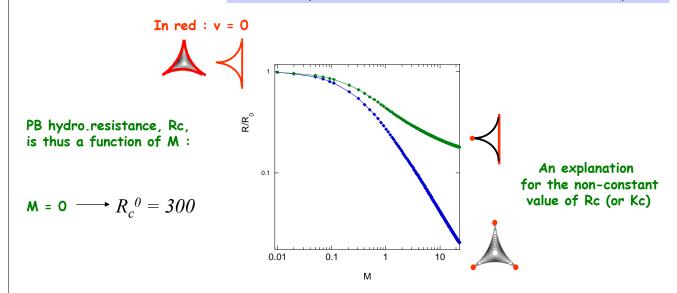
high resistance R_c

$$M = \frac{\mu r}{\mu_s} = 0.9 \frac{\mu D\sqrt{\varepsilon}}{\mu_s}$$



(Koehler et al, JCIS, 2005)

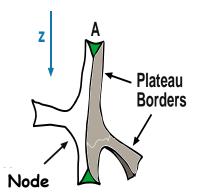
Control parameter and interfacial mobility (3)



numerical formula in : N'Guyen (JCIS, 2002)

Analytical formula in : Koehler , Hilgenfeldt, Stone (JCIS, 2005)

... a simplified form valid for M \leq 10 : $~R_c=~R_c^{~0}\,/\,(1\,+\,2.4\,M)$



interfacial mobility and node resistance

node hydrodynamic resistance, Rn, must thus also be a function of M:

$$R_n = R_n^0 f(M)$$
?

(Cox et al, J. Phys. Cond. Matter, 2001)

M <<1

a node in a Poiseuille flow, with rigid walls?

- Node = correction to the PB length
- R_n / R_c ~ r / L
- As ε , or r, or M decreases : node resistance decreases and becomes smaller than R_c

M >> 1

a node with fluid interfaces?

- R_n must become large
- Simulations : $R_c \sim R_n$

All together...

A simple additive model of hydrodynamic resistances mounted in series :



$$v = \frac{\rho g L^{2}}{\mu} \varepsilon \left(\frac{1}{1/K_{c}(M) + \sqrt{\varepsilon}/K_{n}(M)} \right)$$

$$= \frac{\rho g L^{2}}{\mu} \varepsilon \frac{1}{(1/2)R_{c}(M) + (1/0.9)\sqrt{\varepsilon}R_{n}(M)}$$

One recovers the two extreme limits

In the intermediate range, this can also be written as : v ~ ϵ^{β} , with ~1/2 < β < 1

Experimentally:

Bubble diameter : from 150 μ m to 8 mm

factor 800 on the drainage speed

Bulk viscosity

Interfacial properties:

casein SDS/DOH

SDS

More than 3 orders of magnitude in M

Low M, main dissipation in the Plateau borders:

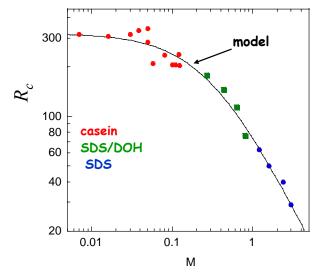
Expt. :
$$\longrightarrow v = \frac{1}{R_c} \frac{\rho g d^2}{\mu} \varepsilon^1$$

Model:

$$R_c = R_c^0 / (1 + 2.4 M)$$

 $R_c^0 = 300$ $M = 0.9 \frac{\mu d\sqrt{\varepsilon}}{\mu_s}$

Comparing data to the model : a way to estimate $\mu_{\rm s}$

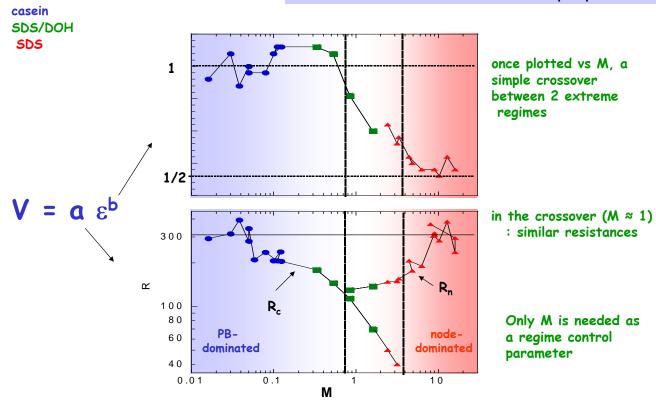


$$\mu_{\rm s} = 10^{-2} \, {\rm g/s}$$
 $\mu_{\rm s} = 1.8 \, 10^{-3} \, {\rm g/s}$ $\mu_{\rm s} = 7 \, 10^{-5} \, {\rm g/s}$ Casein SDS / DOH SDS

very good values of μ_s . in agreement with other studies. validates the assumptions.

All the available data: a final simple picture...

...valid on more than 3 orders of magnitude of M!



No significant transport in the thin films

simulations: films can be important only if: $h/L \ge 10^{-2}$ and M >> 1Koehler, Hilgenfeldt, Stone (JCIS, 2005)

Experimentally:

Large bubbles (D> a few mm)

Large $r: low P_c: low \Pi_d$ (~ 10 Pa)

High flow speeds

High M

- Consistent with directly observed thickness:

 h > a few hundred of nm
- \longrightarrow But, still $V_{film} \leftrightarrow V_{PB}$

small bubbles (D < 1mm)

Low r : high P_c : high Π_d (> 500 Pa)

low flow speeds

low M

Consistent with coarsening rate measurements: h = tens of nm

 \longrightarrow And, still $V_{film} \leftrightarrow V_{PB}$

What to remember?

drainage speed results from a balance between viscous dissipation in Plateau borders and in nodes: each depending on the coupling between the bulk and the surface flows



$$M = \frac{\mu r}{\mu_s}$$

This coupling is described by M, which include the surface SHEAR viscosity. M is also the control parameter which provide the drainage regime.

In the most simplified version:

$$V = (1/R)(\rho g/\mu)L^2\varepsilon^b$$

Taking: <R> = 250

 $\langle R \rangle = 250$ $\langle b \rangle$ between $\frac{1}{2}$ and 1

 \longrightarrow Very reasonable value of v, and of T_drainage : H/V

Next time : gas effect in drainage ??