

Gibbs ?

capillary ?

disjoining ?

# PRESSES IN A FOAM

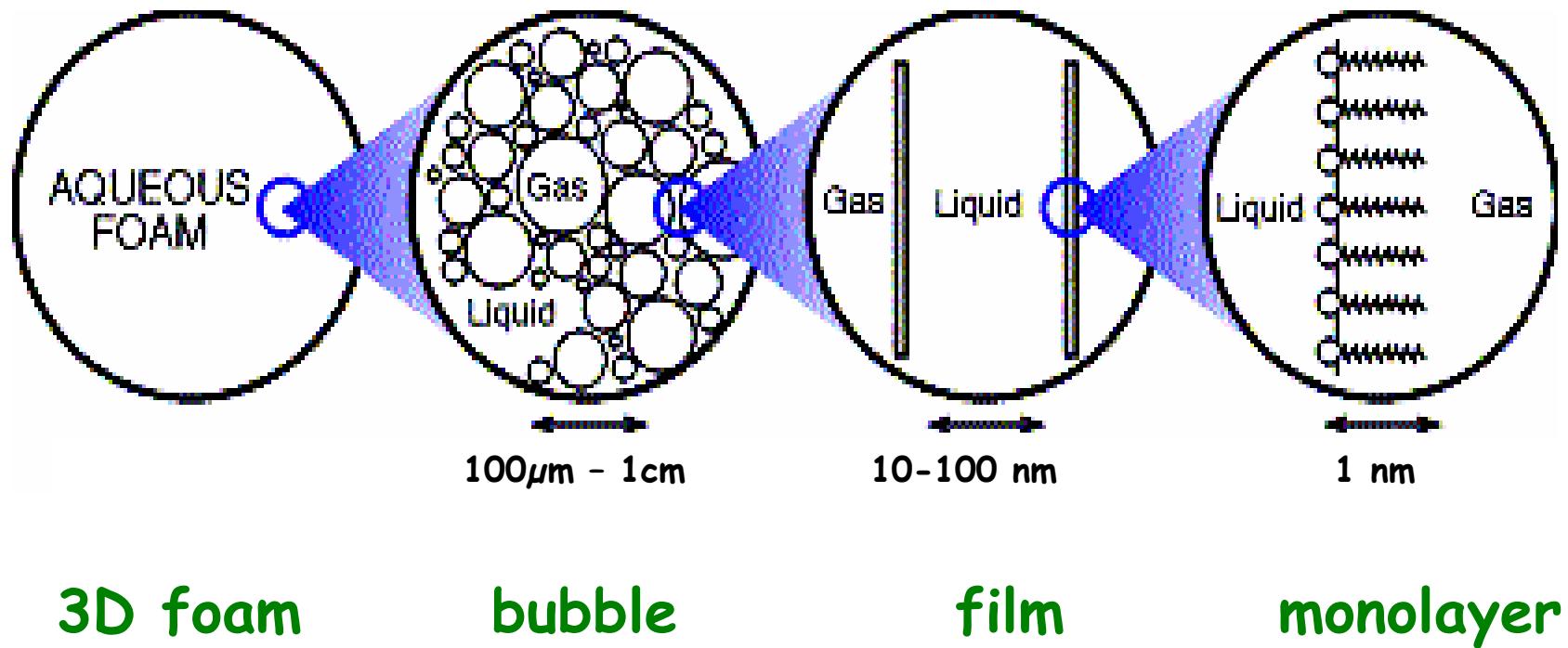
osmotic ?

Laplace ?

Langmuir ?

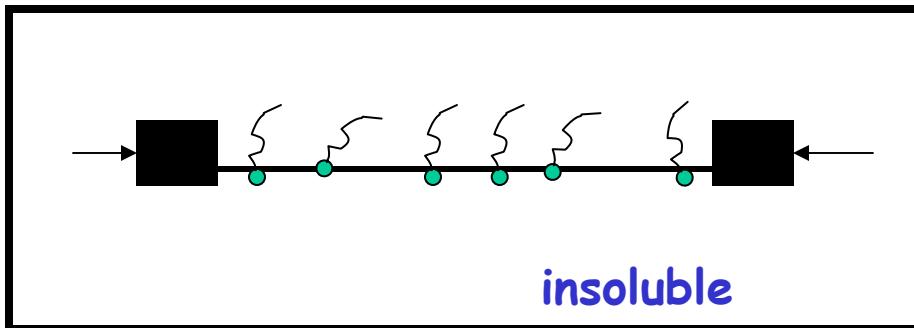
hydrostatic ?

different lengthscales...



... different pressures

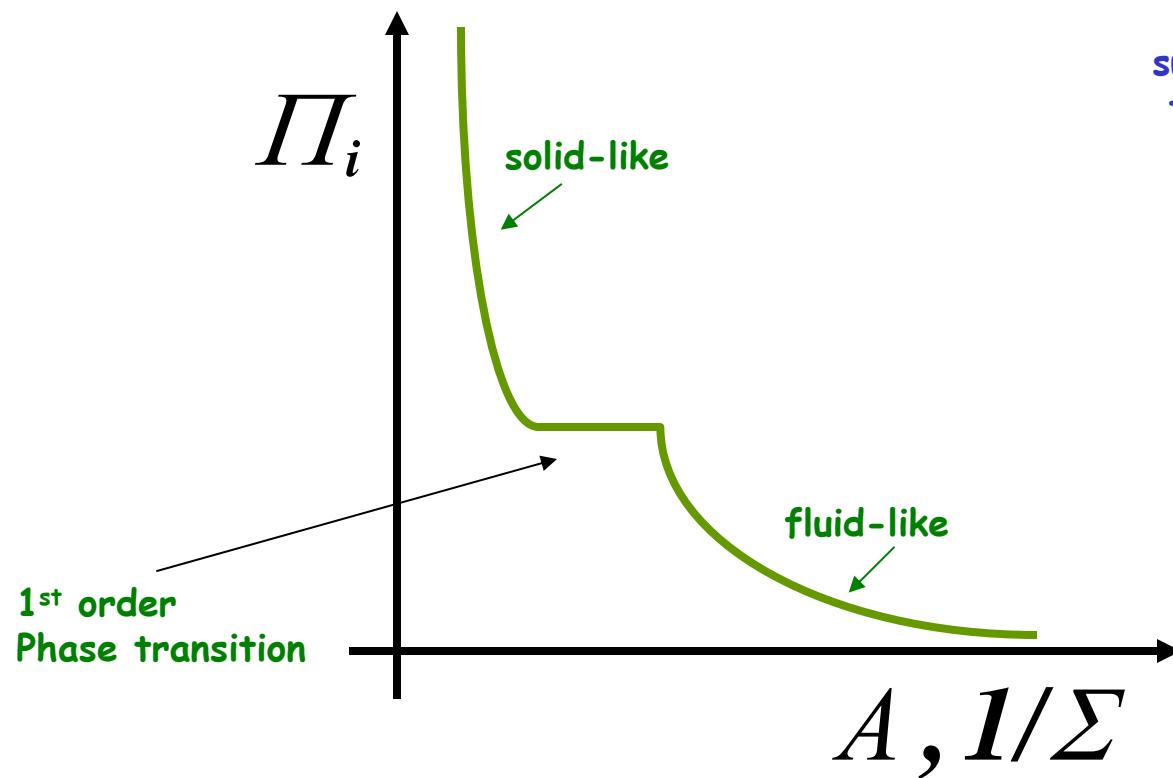
## Liquid interfaces



« Langmuir » monolayer

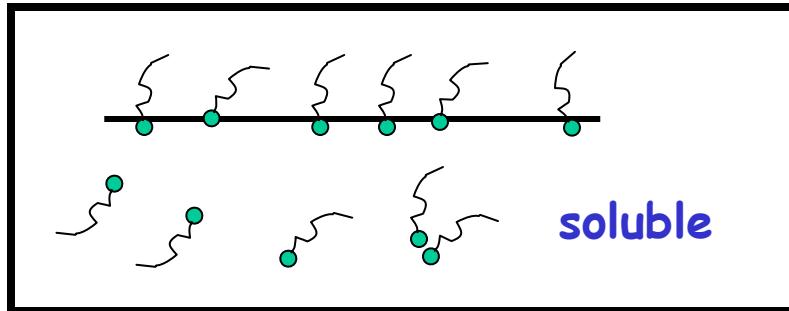
$$\Pi_i = \gamma_0 - \gamma$$

surface tension of  
the pure liquid



$A$  = area per molecule

$\Sigma$  = surface density



Liquid interface

Gibbs equation :

$$\Gamma = -\frac{1}{kT} \frac{\partial \gamma}{\partial \ln c}$$

A simple model (Langmuir) : equilibrium between adsorption and desorption

$$\Gamma / (\Gamma_\infty - \Gamma) = \frac{c}{a} \quad \Gamma = (a/c + 1)^{-1} \Gamma_\infty$$

With  $a$  = Szyszkowski concentration

With  $\Gamma_\infty$  = saturation surface concentration

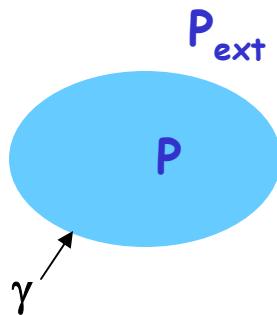
Finally, the « Langmuir » (interfacial) equation of state :

$$\Pi_i = -kT \Gamma_\infty \ln (1 - \Gamma/\Gamma_\infty)$$

More refinements are possible

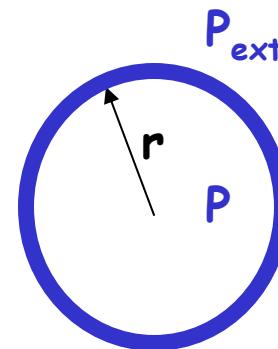
Though, it can be used dynamically

## curved interfaces and Laplace pressure



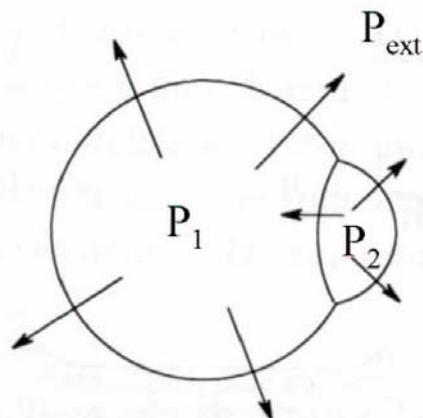
$$P - P_{ext} = \gamma (1/R_1 + 1/R_2)$$

For a single spherical bubble :



$$P - P_{ext} = 4\gamma / R$$

Two bubbles of different sizes, in contact :



$$P_2 > P_1 > P_{ext}$$

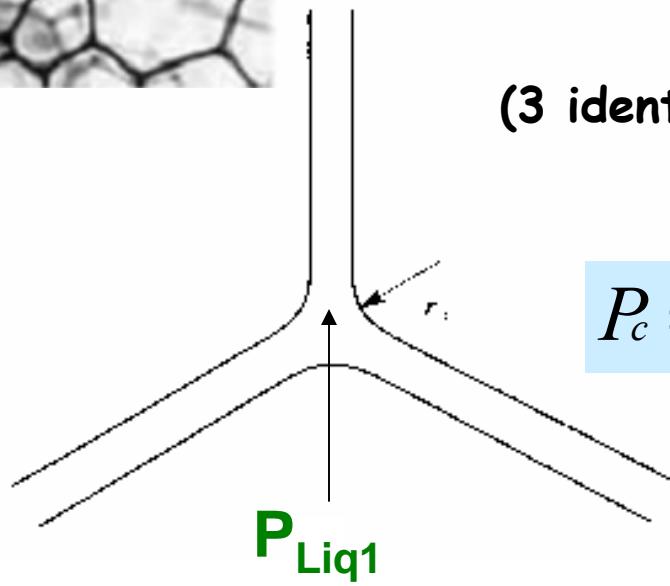
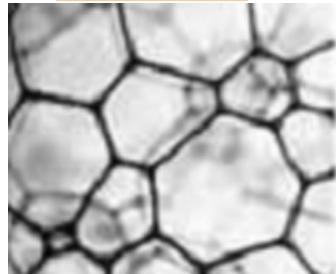
A useful lengthscale :

$$l_c = \sqrt{\frac{\gamma}{\rho g}}$$

$l < l_c$  : capillarity dominates gravity

## Liquid pressure in Plateau borders

DRY

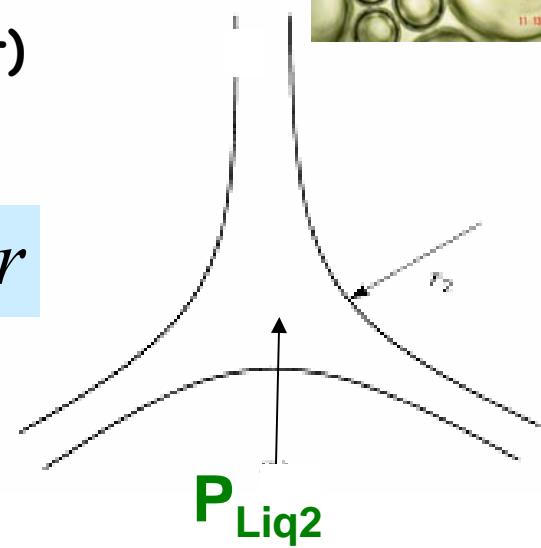
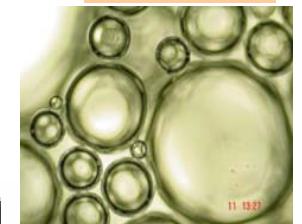


The Laplace pressure also describes the pressure inside the Plateau borders :

(3 identical bubbles at  $P_g = cst$ )

$$P_c = P_{gas} - P_{liquide} = \gamma/r$$

WET

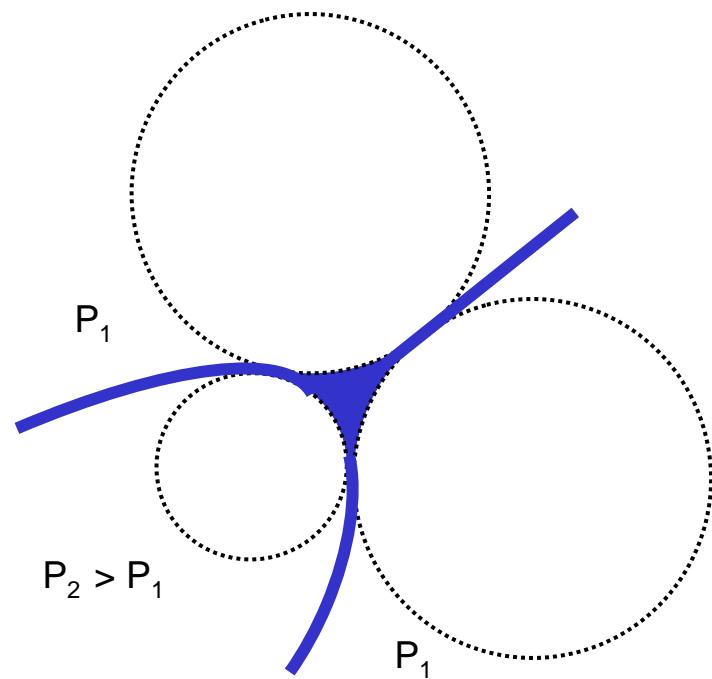


$$P_{Liq2} > P_{Liq1}$$



Gradients of liquid fraction = gradients of pressure = ...

Considering bubbles of different volumes...

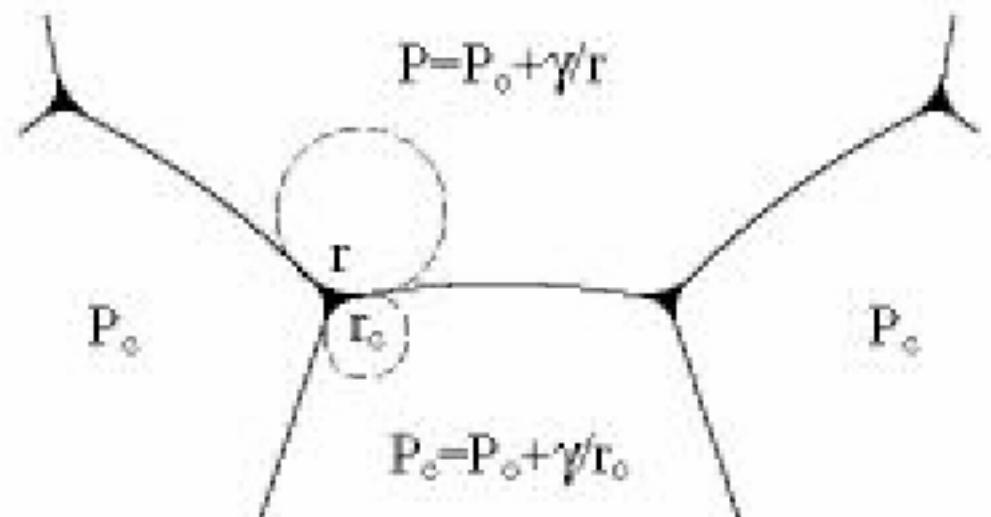


differences in thin film curvature

...as well as...

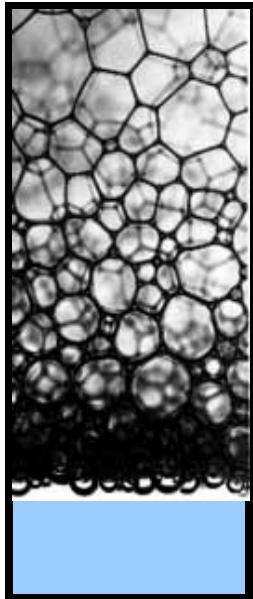
differences in Plateau borders radii

$$P_c = P_{gas} - P_{liquide} = \gamma/r$$



## Equilibrium under gravity and hydrostatic pressure

$$P_c = P_{gas} - P_{liquide} = \gamma/r$$



Hydrostatic pressure :  $P_{liq} = P_0 + \rho g z$

$$\rightarrow r(z) = \frac{\gamma}{(P_g - P_0 - \rho g z)}$$

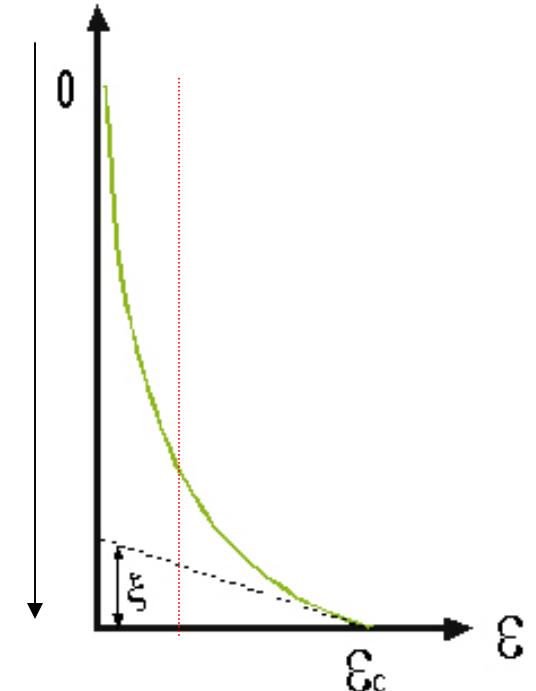
using  $\varepsilon = c (r/D)^2$

$$\rightarrow \varepsilon(z) = \frac{c\gamma^2}{D^2} \frac{1}{(P_g - P_0 - \rho g z)^2}$$

With the bottom boundary condition :

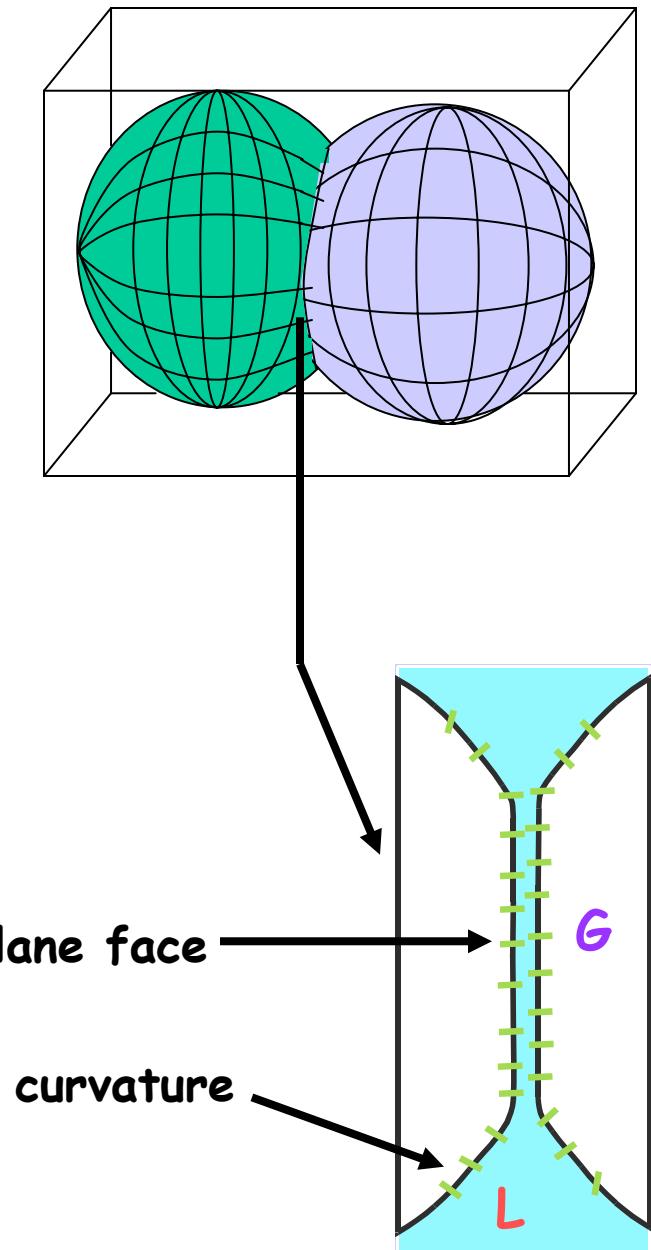
$$\varepsilon_c = 0.36$$

$$\rightarrow P_g - P_0 = \rho g H + \sqrt{\frac{c}{\varepsilon_c}} \frac{\gamma}{D}$$

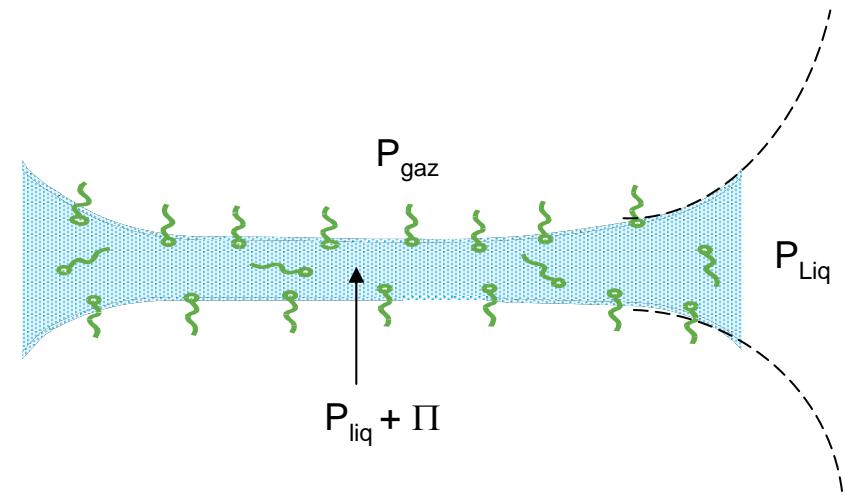


Height of the  
« wet layer »:

$$\xi = \sqrt{\frac{\gamma}{\rho g R}}$$



Something wrong in the thin films ?

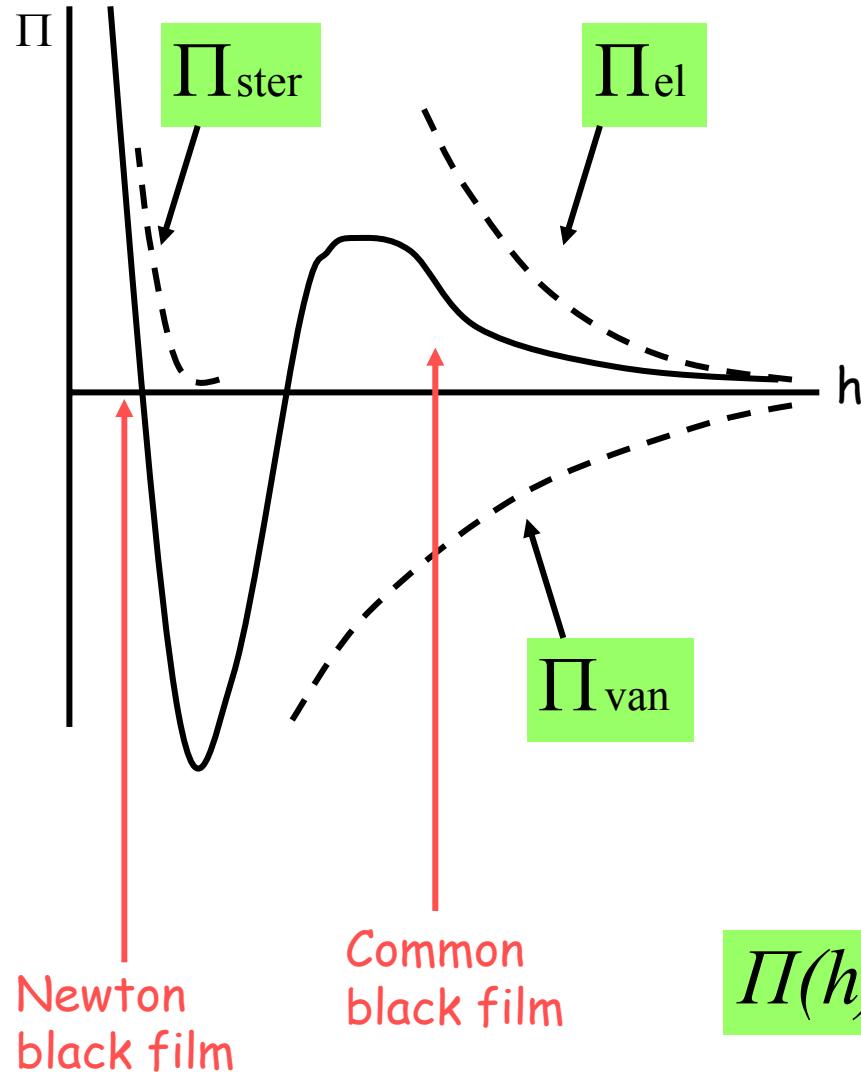


$$\Pi_d + P_{liquide} = P_{gas}$$

$$\Pi_d = P_{gas} - P_{liquide} = P_c$$

At equilibrium

which forces between surfaces?



✓ electrostatic repulsion :

$$\Pi_{el} = 64 kT \rho_\infty \gamma^2 \exp(-kh)$$

✓ dipolar interaction  
(Van Der Waals):

$$\Pi_{van} = -\frac{A_h}{6 \pi h^3}$$

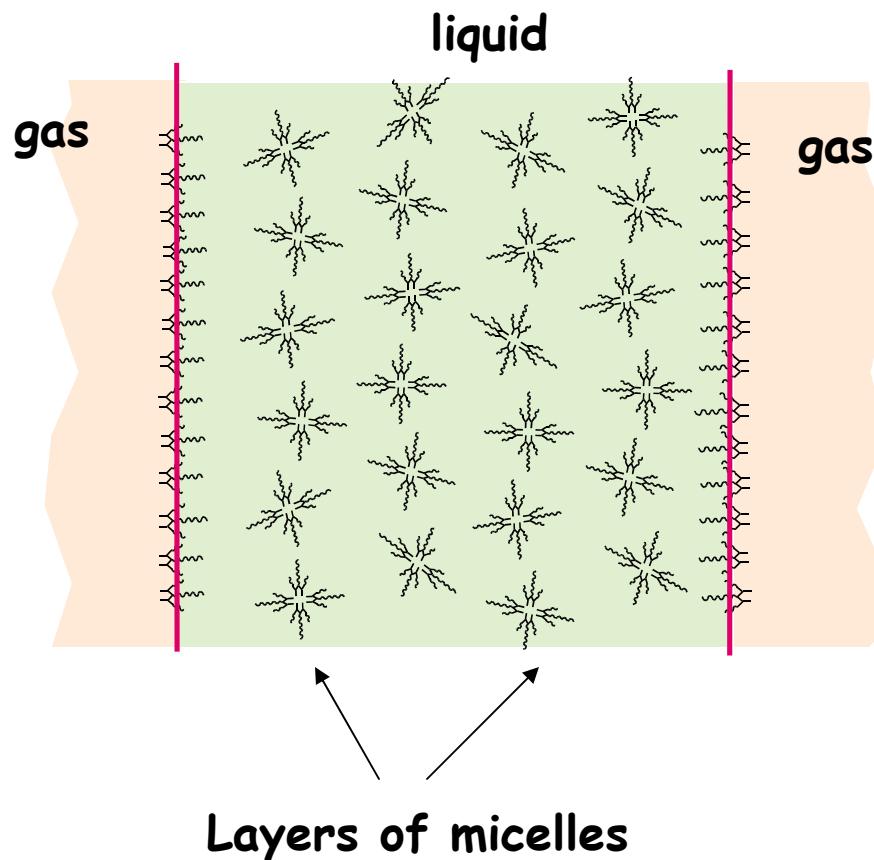
✓ steric contribution :  
confinement

$$\Pi(h) = \Pi_{el} + \Pi_{van} + \Pi_{ster}$$

(DLVO model)

Typical  $h$ , and  $\Pi$  are...

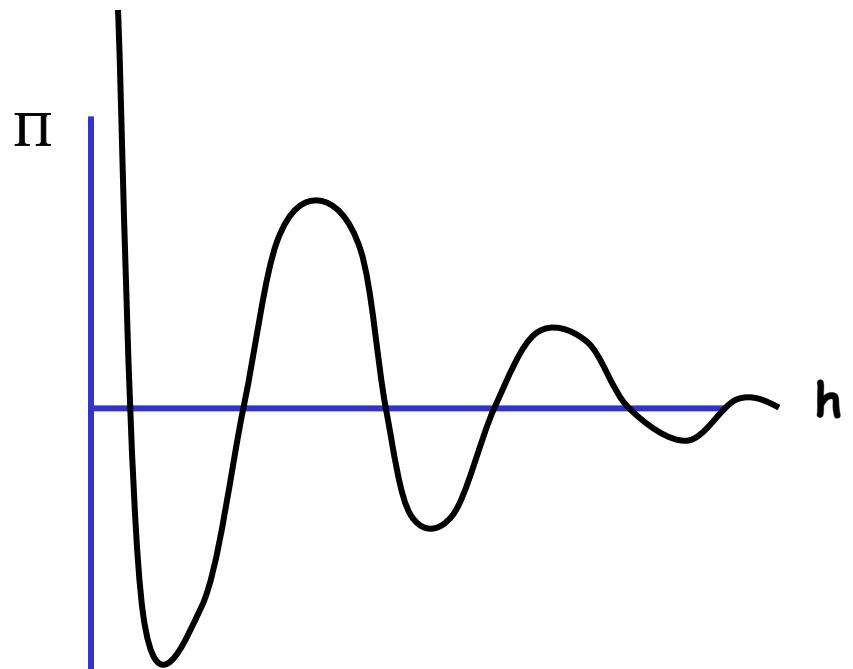
## Another type of force : supramolecular forces



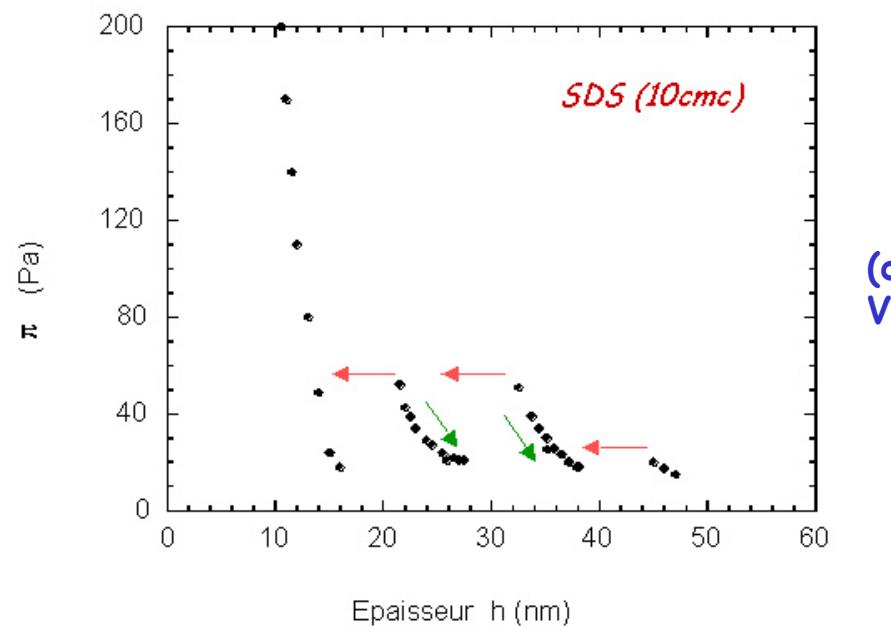
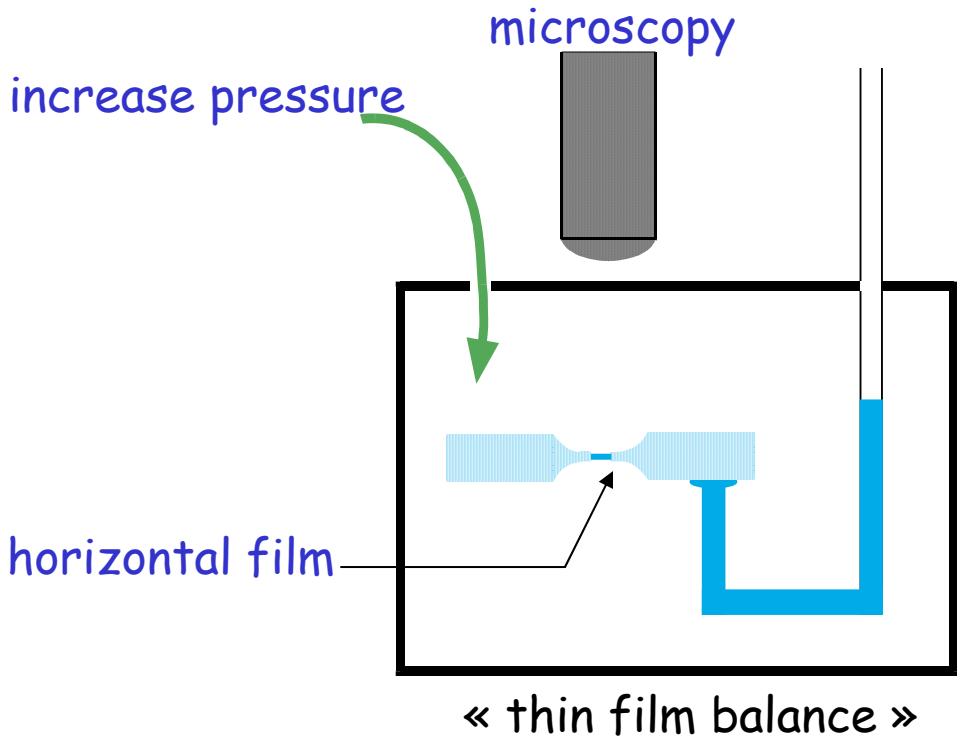
confinement effect in the film

Arrangements in layers

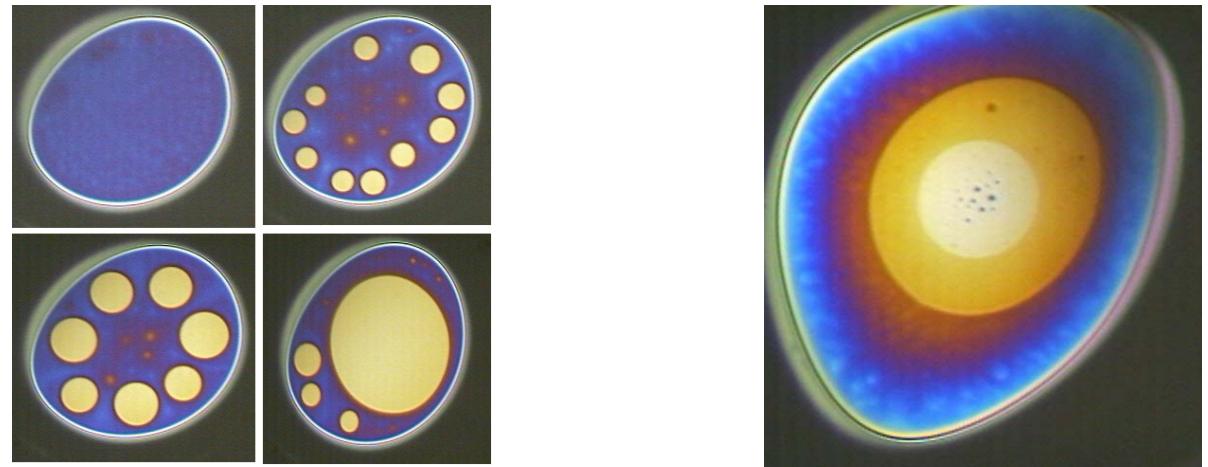
Oscillation in the disjoining pressure



implies thickness steps,  
and a thinning layer by layer



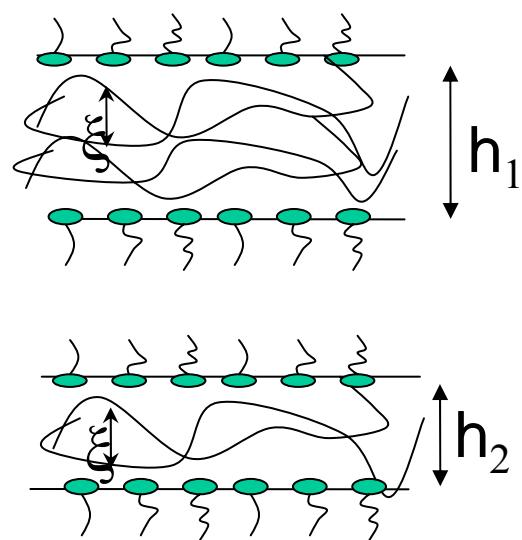
Other types of film stratification,  
at high thickness :



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Stratification with surfactant-polyelectrolytes systems :

Confinement effects  
in thin films



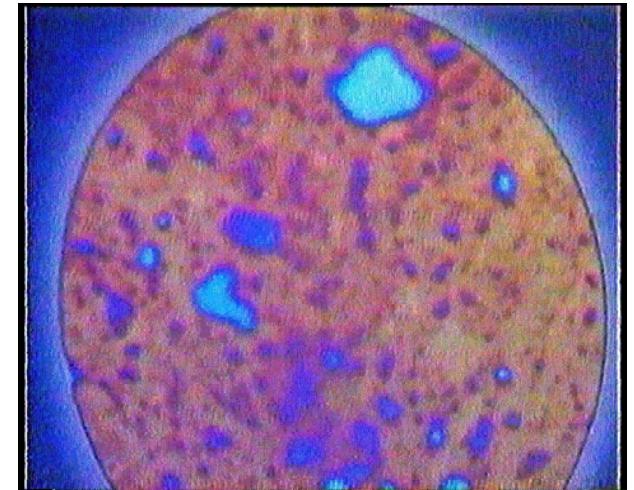
layer thickness  
corresponds  
to the bulk mesh size (!?)

$$\xi \sim h_1 - h_2$$

high concentration of polyelectrolytes and surfactant :

→ non-homogeneous and non-flat films

gelified structure : high local viscosity

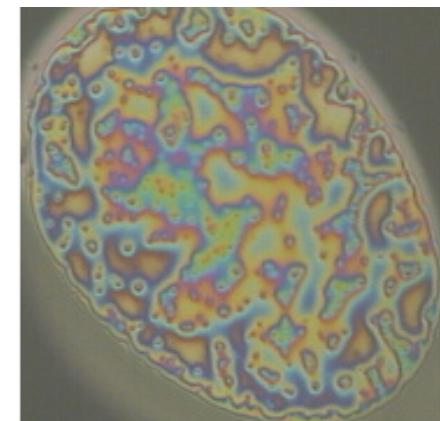
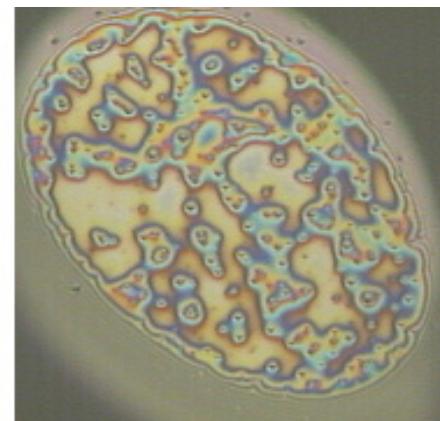
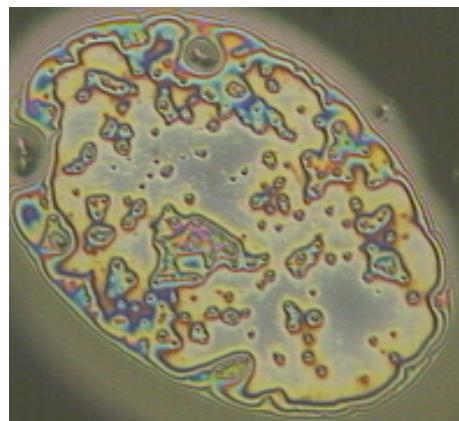


Origins of stability ?

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Increasing the protein concentration →

casein  
films



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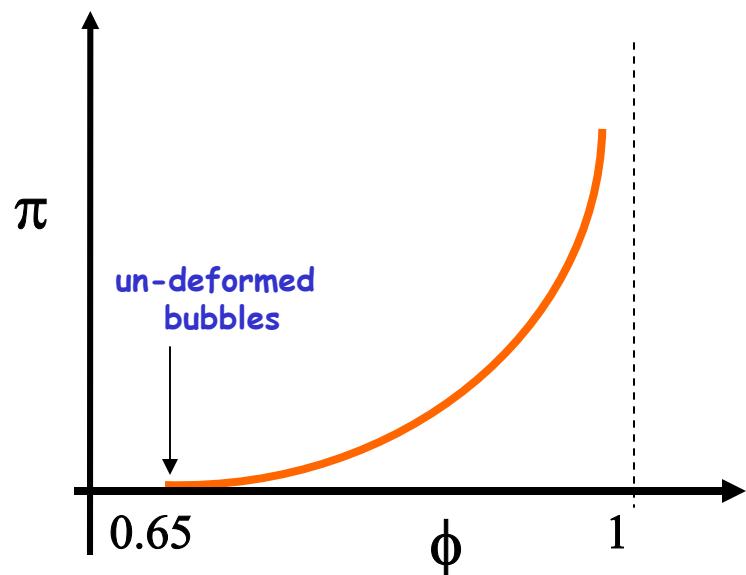
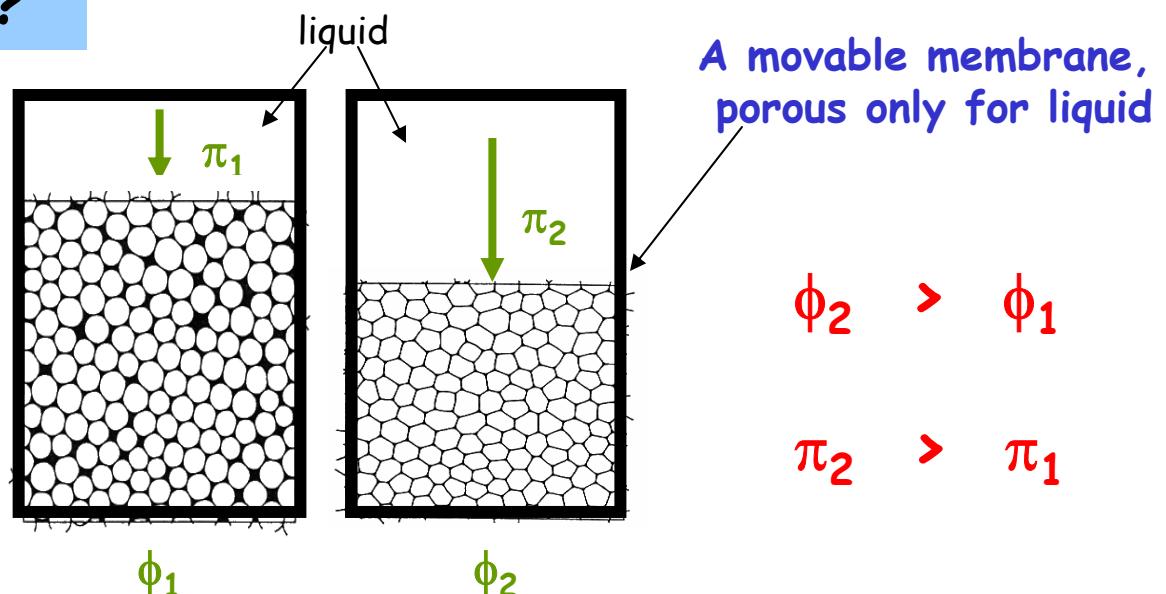
stability →

# osmotic pressure, $\pi$ ?

Princen (1979)

$A(\phi)$  ?  $E(\phi)$  ?

*in analogy  
with solutions...*



$$\pi(\varphi) = \frac{\gamma}{r} \phi^2 \frac{\partial}{\partial \phi} \left[ \frac{A}{A_0} \right]$$