

Solid-like response

Linear viscoelasticity

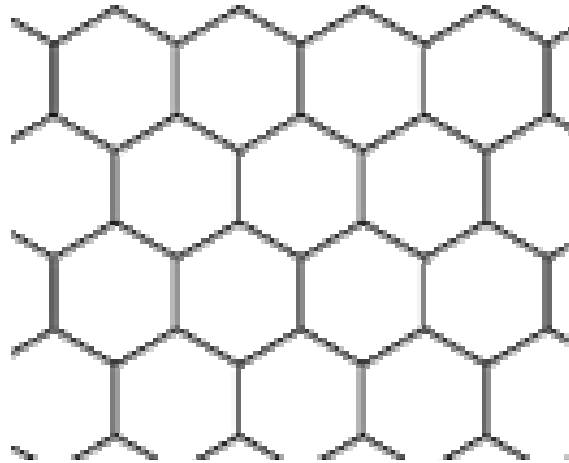
Non linear elasticity



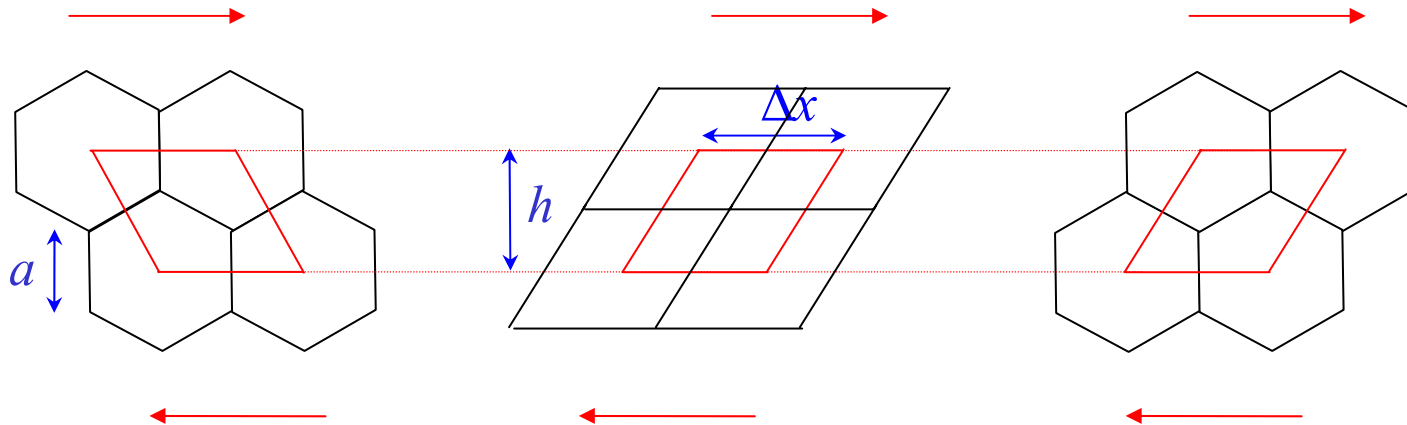
Laboratoire de Physique des Matériaux Divisés et des Interfaces



Princen model

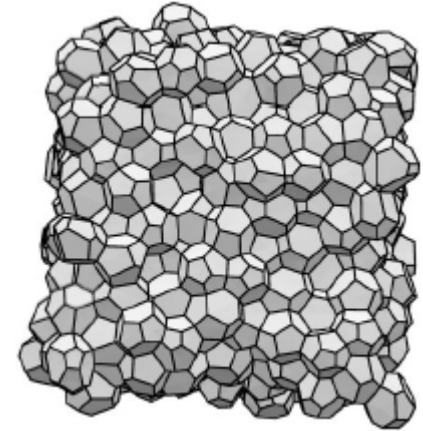
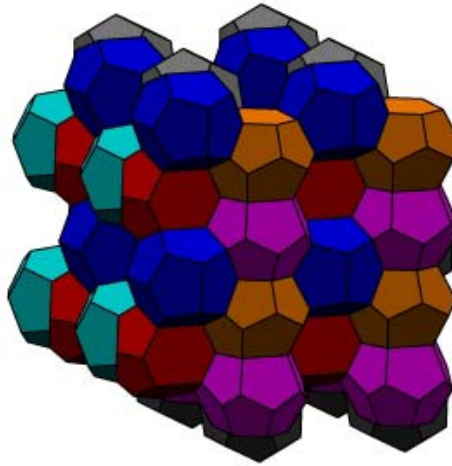
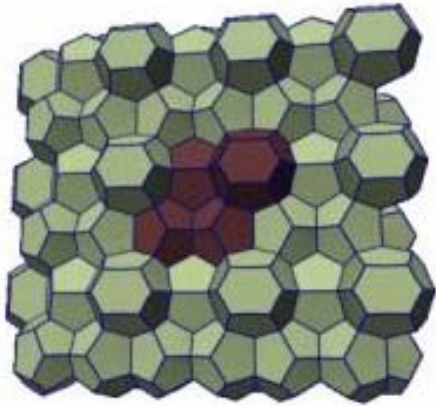


Vincent-Bonnieu



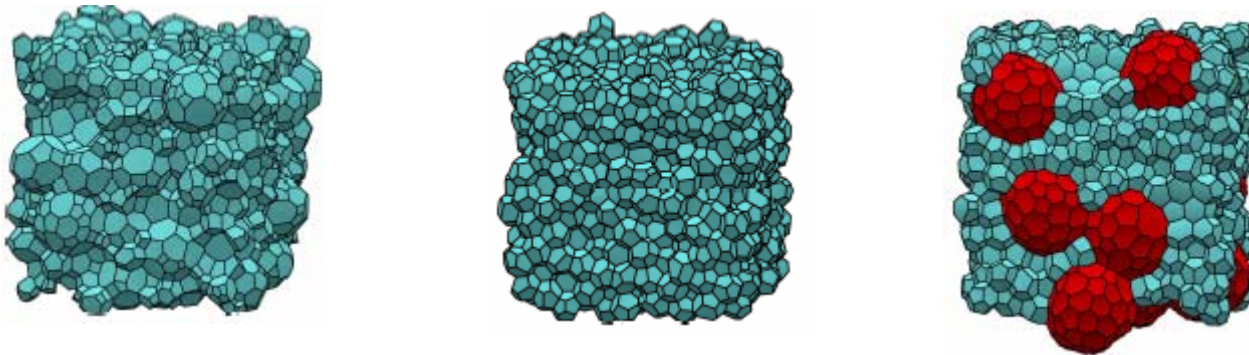
Princen 1983

3D dry foams

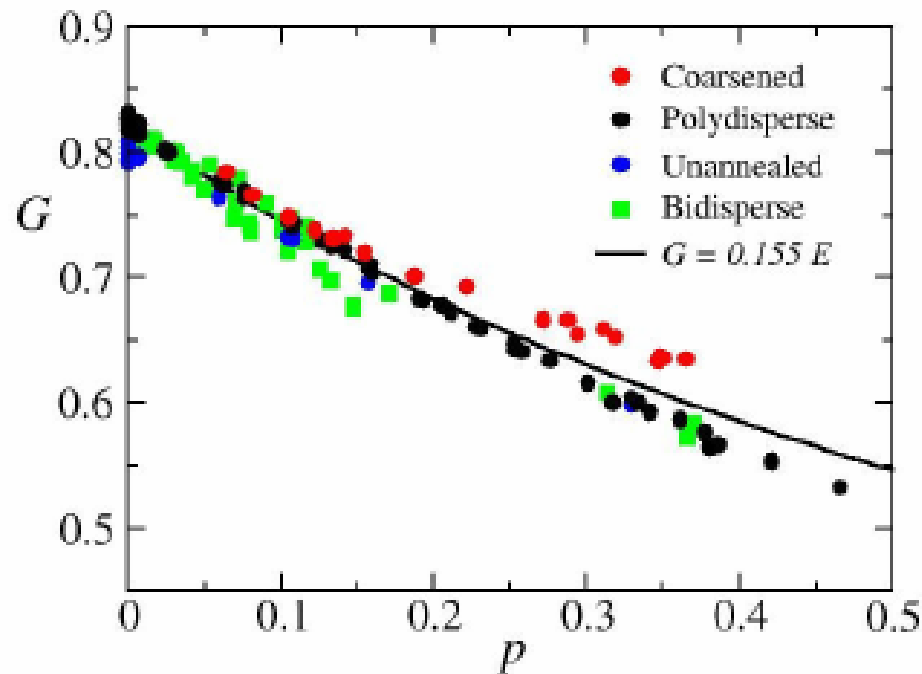


	$R_{G/T}$	$\Delta G/G$
Kelvin	0.50	0.5
Williams	0.49	0.5
Weaire-Phelan	0.54	0.04
Disordered monodisperse	0.51	0
Derjaguin estimation	0.8	0

3D dry disordered polydisperse foams

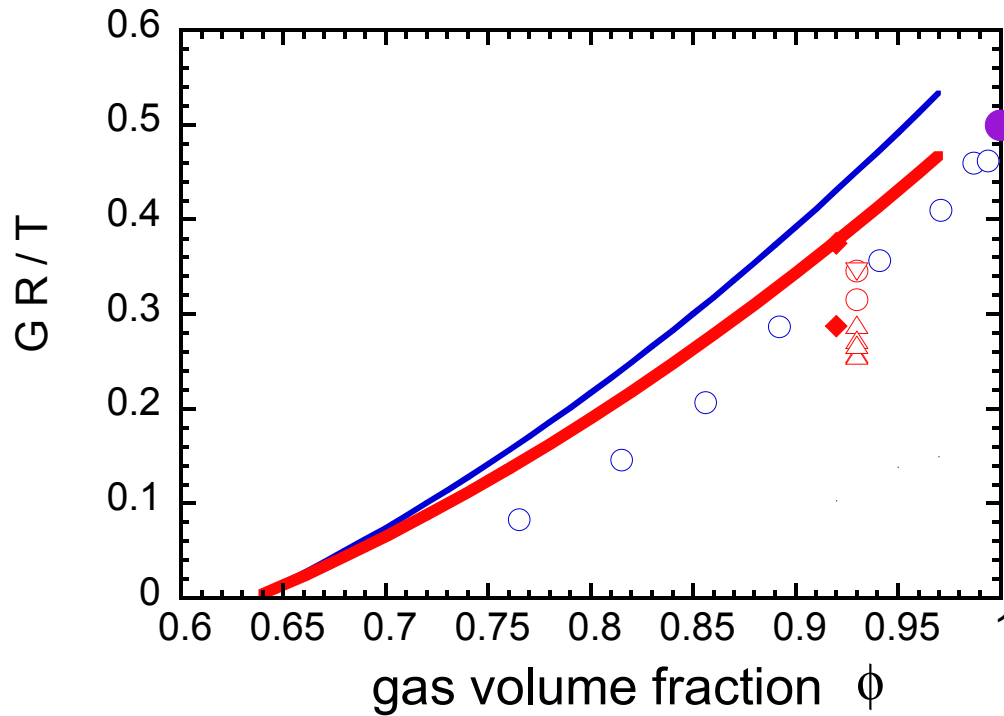


$$p = \frac{R_{32}}{\langle R^3 \rangle^{1/3}} - 1$$



$$G \cong 0.5 \frac{T}{R_{32}}$$

Effect of liquid content



- Emulsions
- Monodisperse, R_1
 - Polydisperse, R_{32}
- Polydisperse foams
- R_1
 - Δ ∇ R_{32}
- Surface Evolver simulation

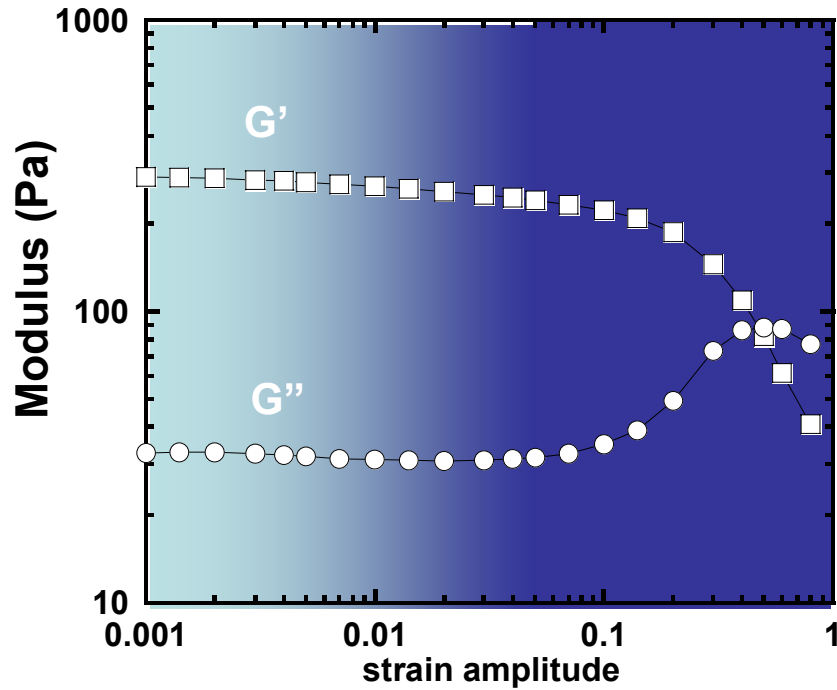
$$G = \alpha \phi (\phi - \phi_c) \frac{T}{R_{32}} \quad \alpha \cong 1.4$$

*Princen, Kiss 1986; Mason et al 1995; Saint-Jalmes, Durian 1999;
Gopal, Durian 1999; Höhler, CA 2005; Kraynik, Reinelt 2004
Bolton, Weaire 1990*

Does a linear regime exist ?

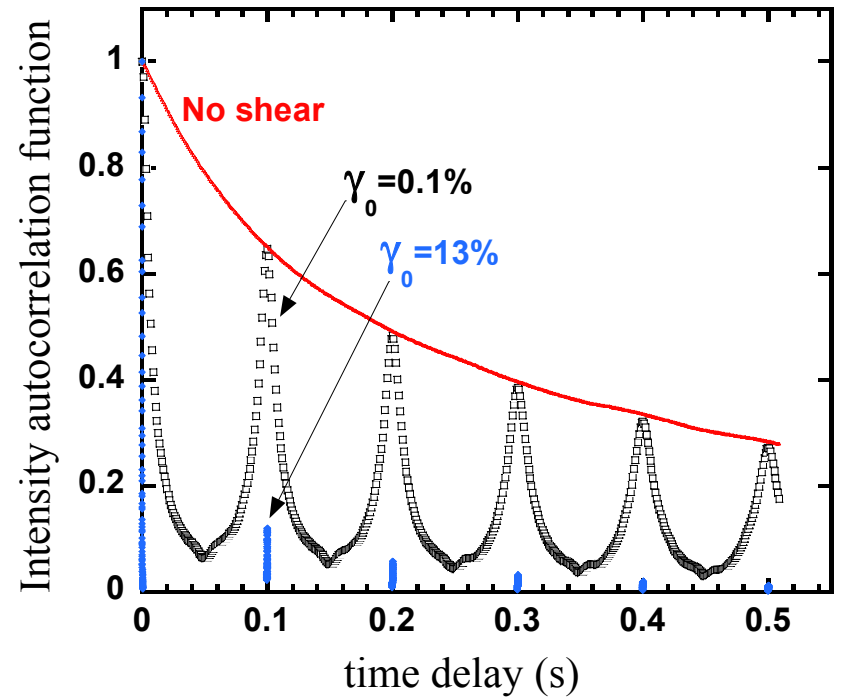
Weaire, Fortes 1994

Rheometry

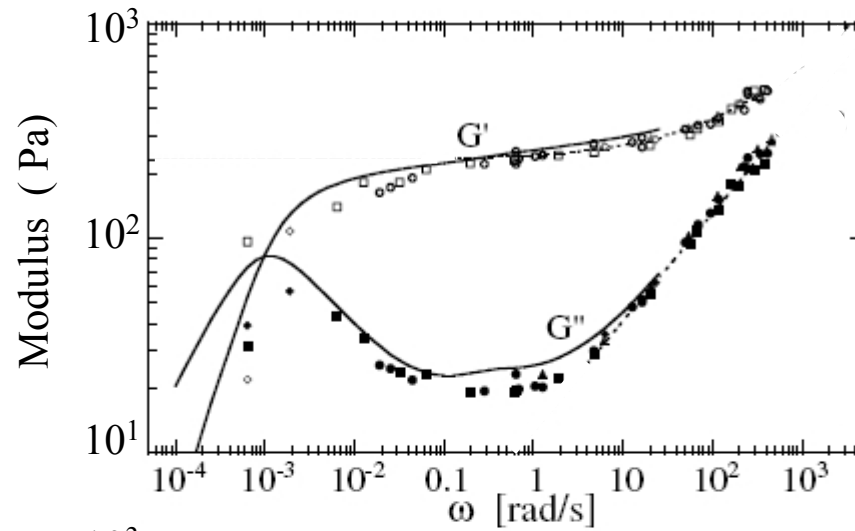


Gillette foam
Age = 60 min
Frequency = 10 Hz

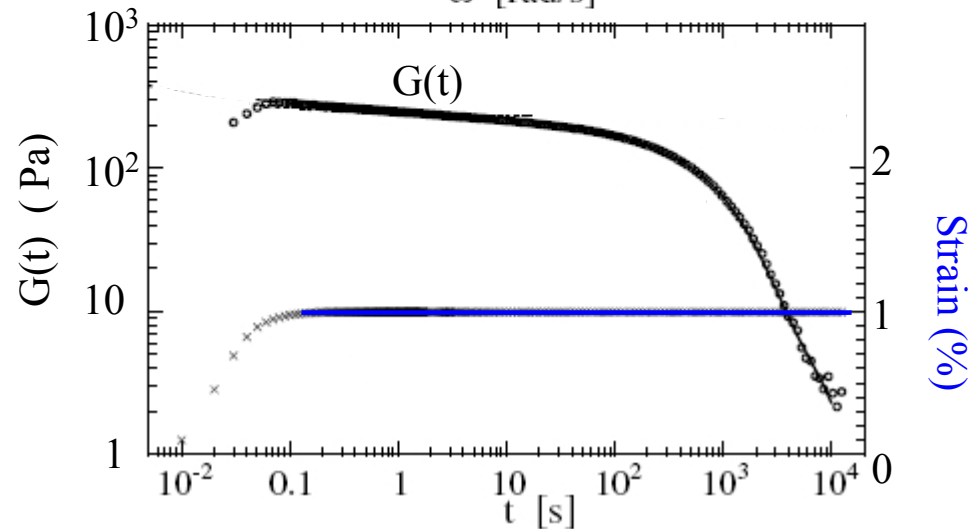
DWS echoes



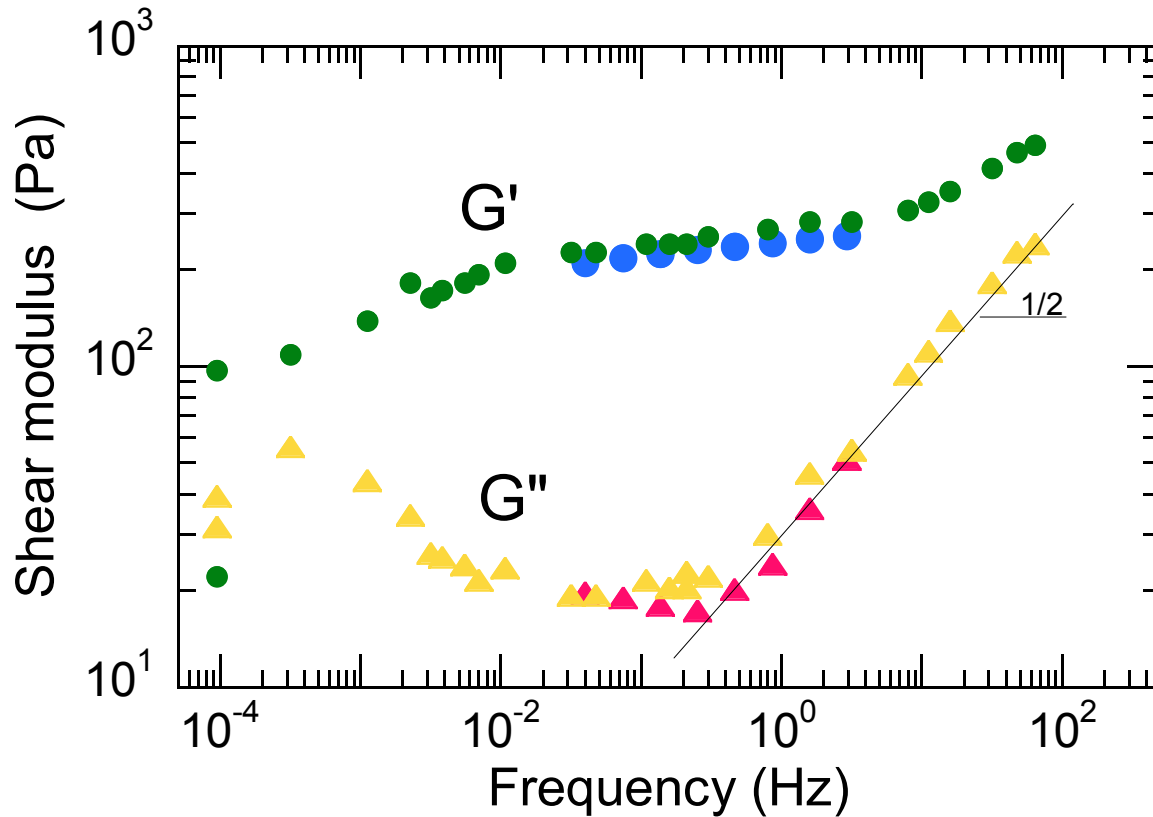
Equivalence between temporal and frequency responses



$$G^*(\omega) = i\omega \mathcal{L}[G(t)][i\omega]$$



Viscoelastic response



Gillette foam
Age $\cong$ 100 min

● ▲ C-A, Hoballah, Höhler 1998
● ▲ Gopal, Durian 2003

Dissipation at the film scale

- Fluid transport in films and borders
- Intrinsic surface viscosity of the interfaces

Buzza, Lu, Cates, *J Phys II* 1995

Edwards, Wasan 1996

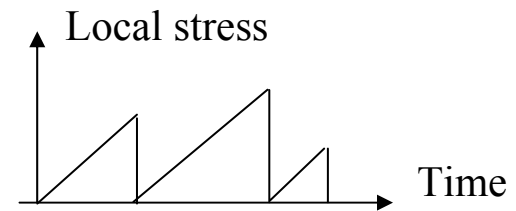
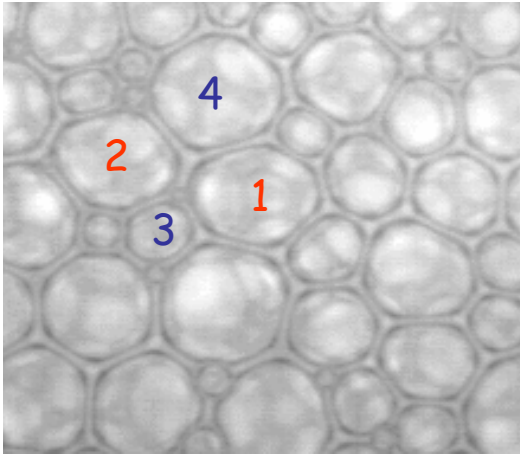
Dissipation at a mesoscopic scale “Soft Glassy Materials”

- Irreversible loss of elastic energy upon bubble rearrangements

Sollich et al *PRL* 1997

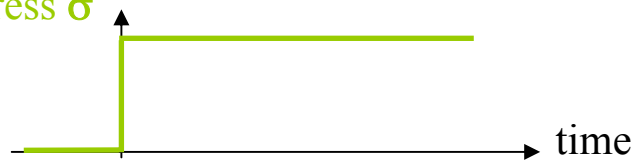
Hébraud, Lequeux *PRL* 1998

Coarsening is an intrinsic source of dynamics in foams

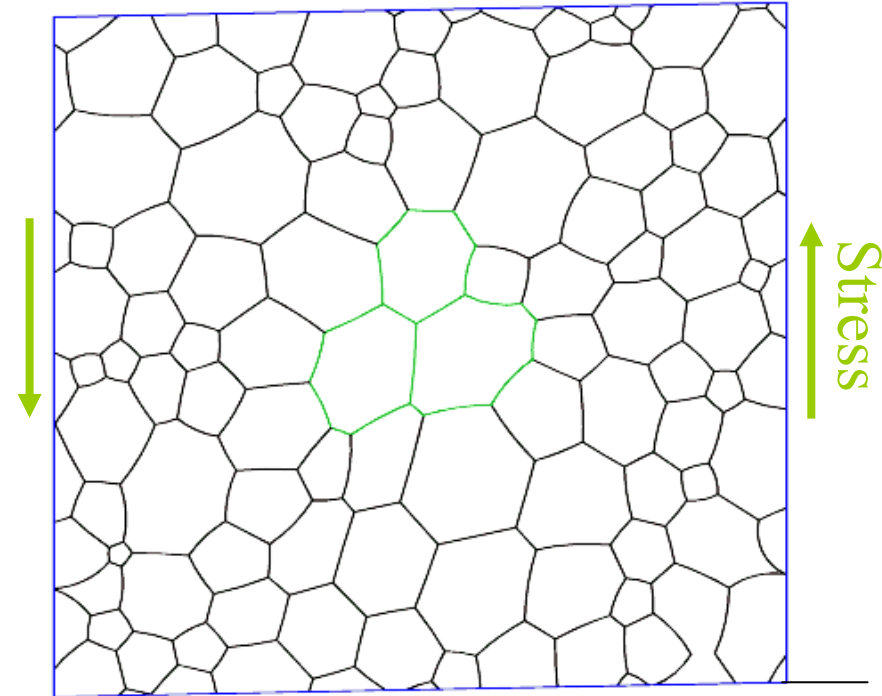
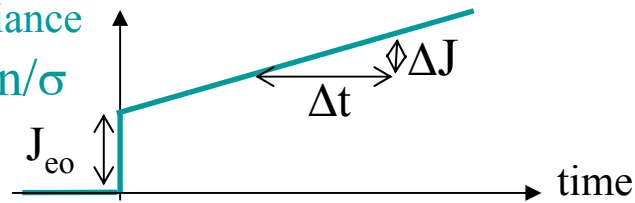


A simple mesoscopic model of relaxation due to coarsening-induced rearrangements

Applied shear stress σ

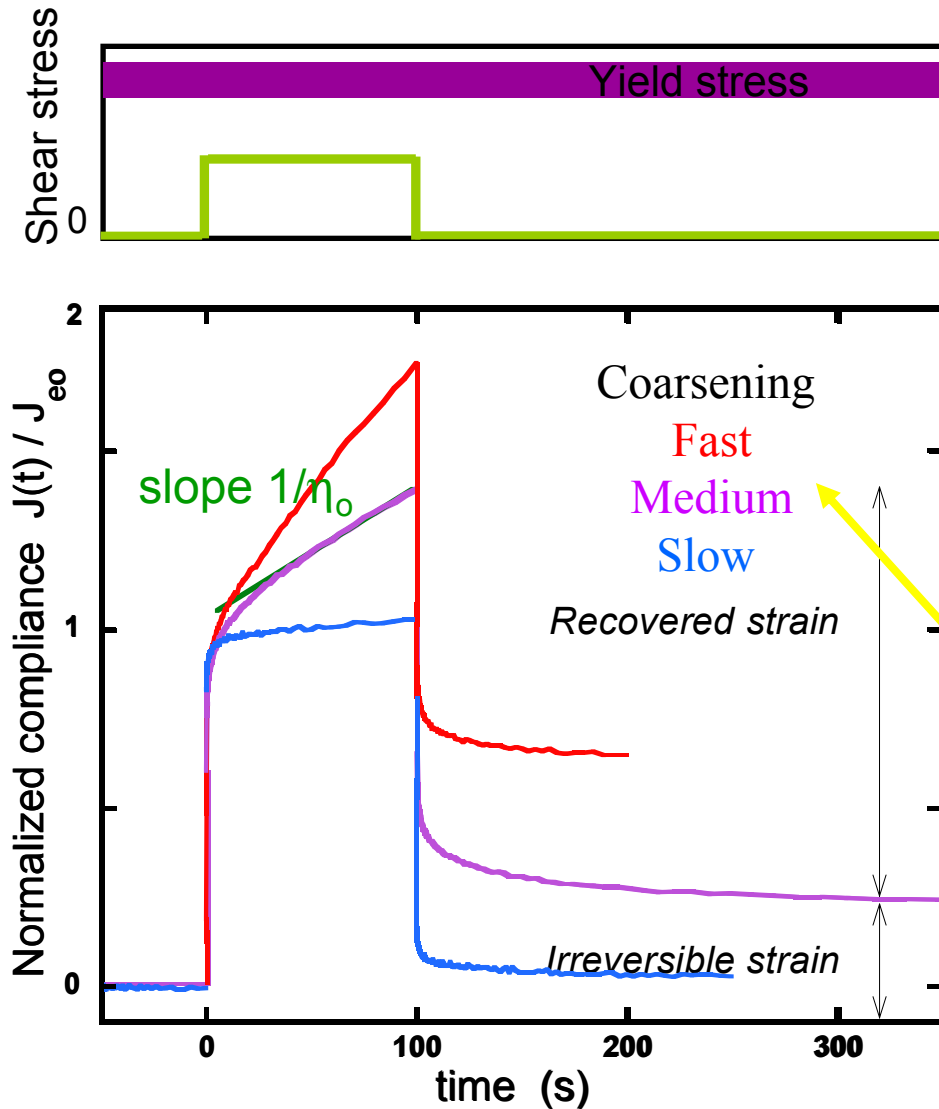


Shear compliance
 $J(t) = \text{strain}/\sigma$



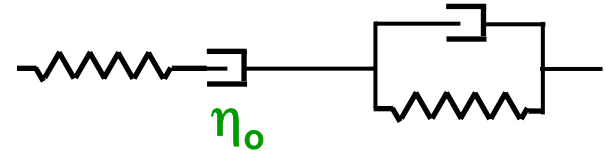
Surface Evolver simulations
S. Vincent-Bonnieu

Creep, DWS and DTS experiment



Only two relaxation processes !

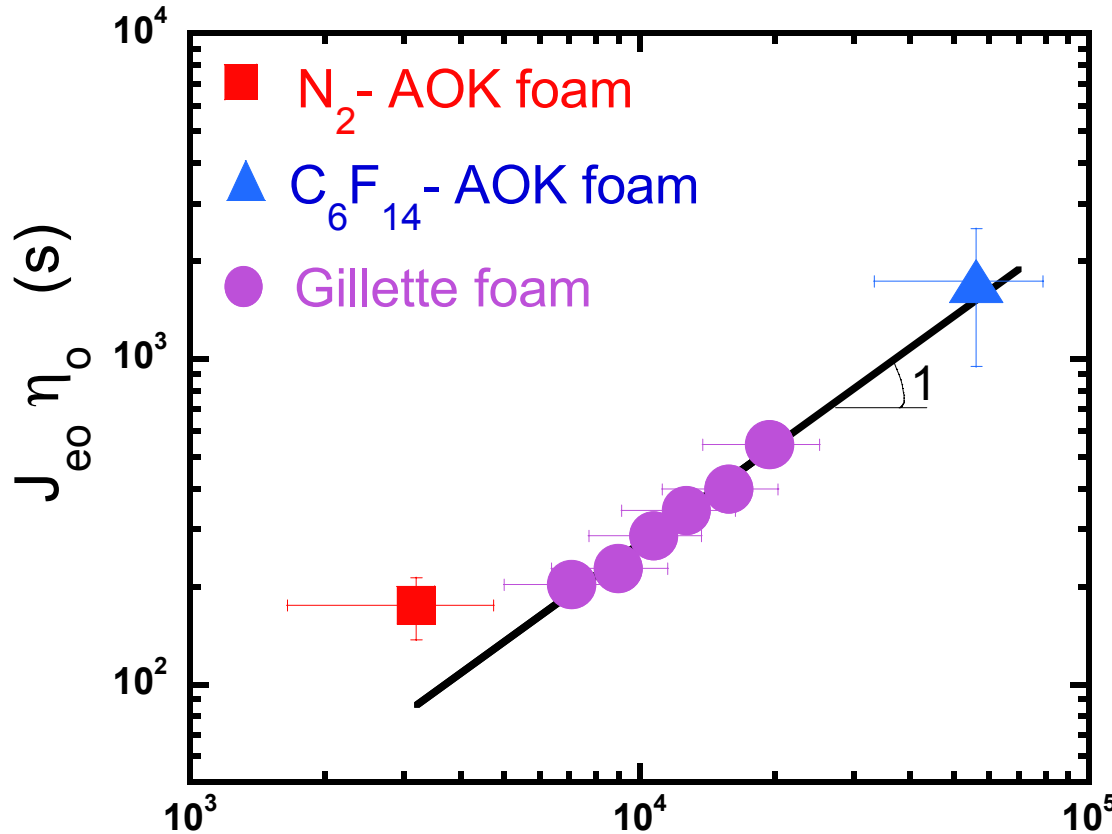
$$J(t) = J_o + \frac{t}{\eta_o} + J_1 \left(1 - \exp\left(-\frac{t}{\eta_1 J_1}\right) \right)$$



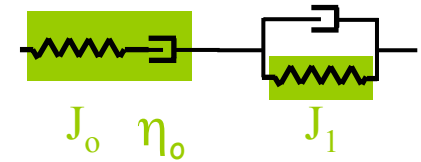
*In agreement with
Marze, Saint-Jalmes, Langevin 2005*

Tunable by the physico-chemistry

Coarsening induced rearrangements are at the origin of the steady creep



Mean time interval between rearrangements $1/(R d^3)$ (s)



$$J_{eo} = J_o + J_1$$

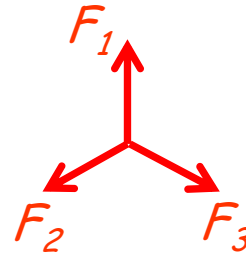
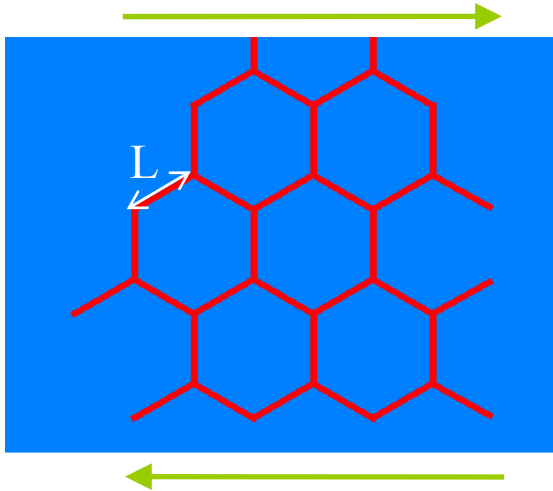
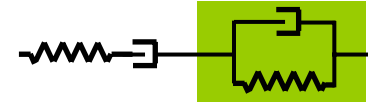
$$J_{eo} \eta_o \cong \frac{1}{R V}$$

$$V \cong (3 d)^3$$

Foams and other solids can flow on long time scales



Interfacial relaxation



Forces acting at the film junctions:
Edwards, Brenner, Wasan 1991

$$F_i = 2 \left(T + \kappa \frac{1}{L} \frac{\partial L}{\partial t} \right)$$

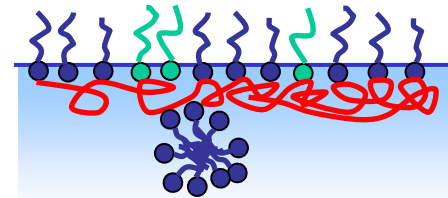
Surface
tension

Dilatational
surface viscosity

Scaling law for the relaxation time, independent of bubble size : $\tau_1 \approx \kappa / T$

$$\tau_1 \cong 3 - 5 \text{ s} \rightarrow \kappa \approx 0.05 - 0.15 \text{ kg s}^{-1}$$

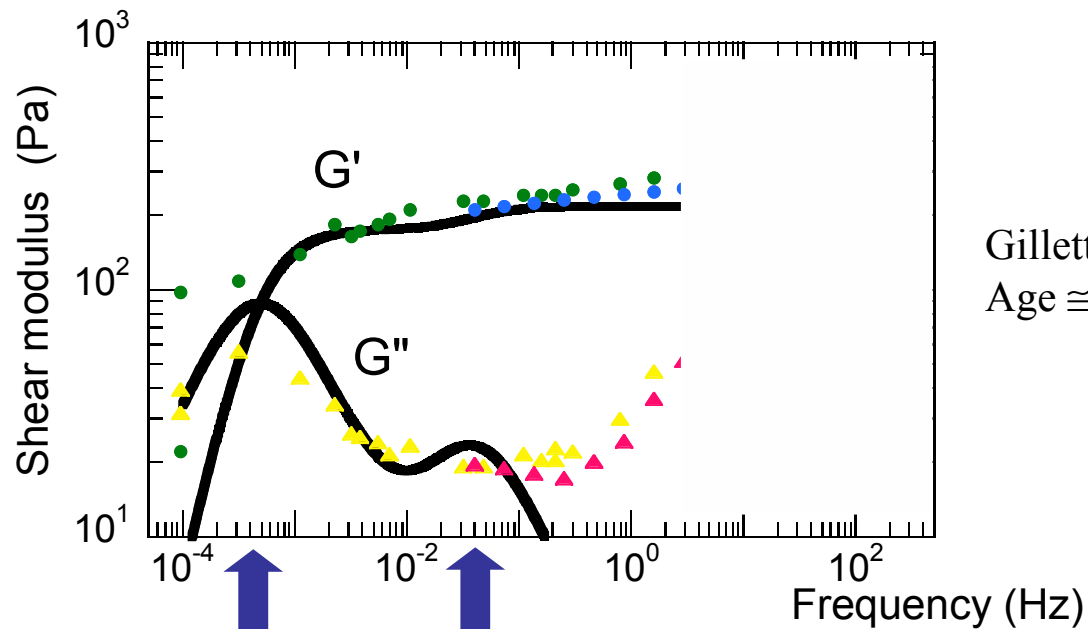
Surfactants + Dodecanol + Polymer



Back to the frequency response

Complex shear modulus :
$$G^*(\omega) = \frac{1}{i\omega \mathcal{L}[J(t)][i\omega]}$$

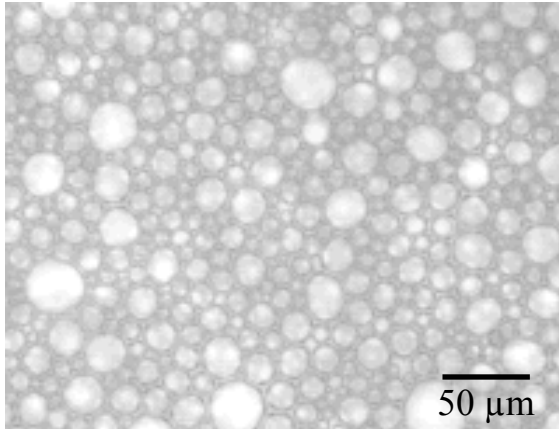
$J_0 \quad \eta_0 \quad J_1 \quad \eta_1$



Coarsening induced
rearrangements

Surface viscosity

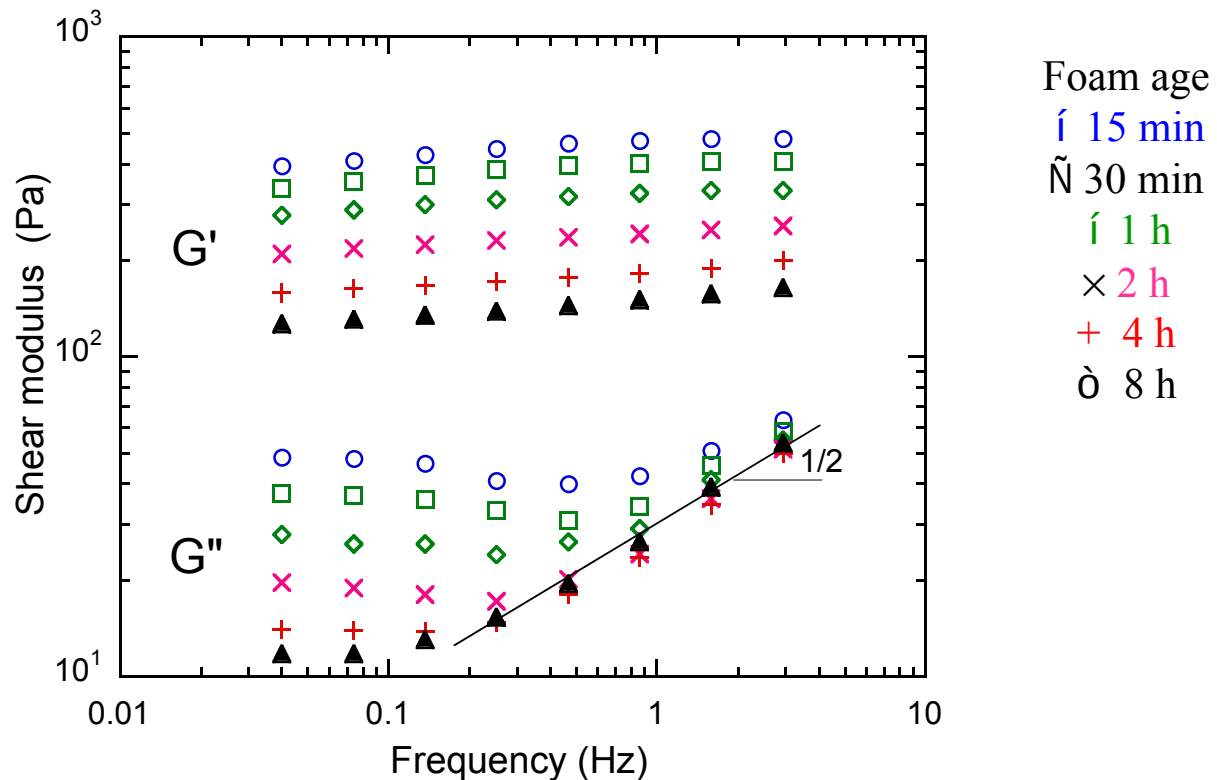
Evolution upon coarsening



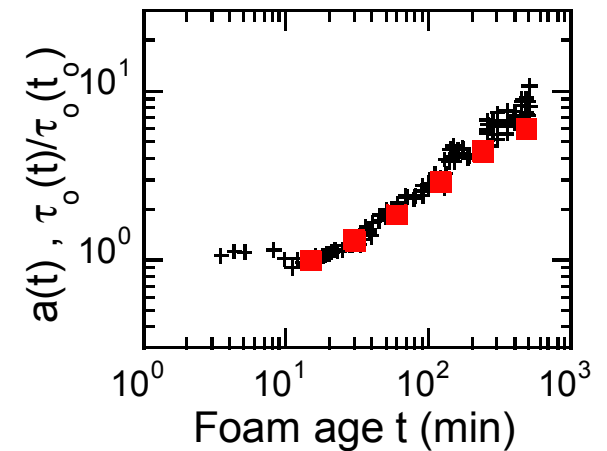
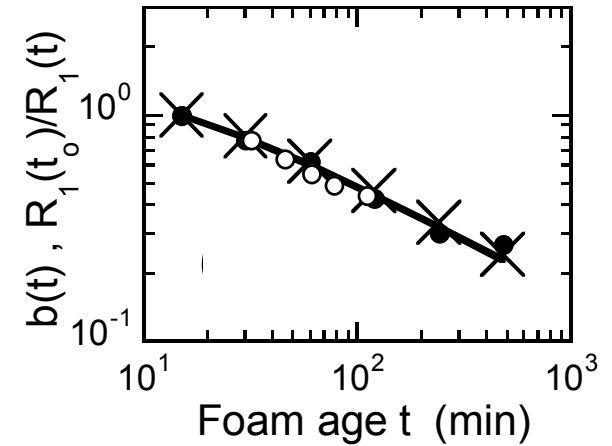
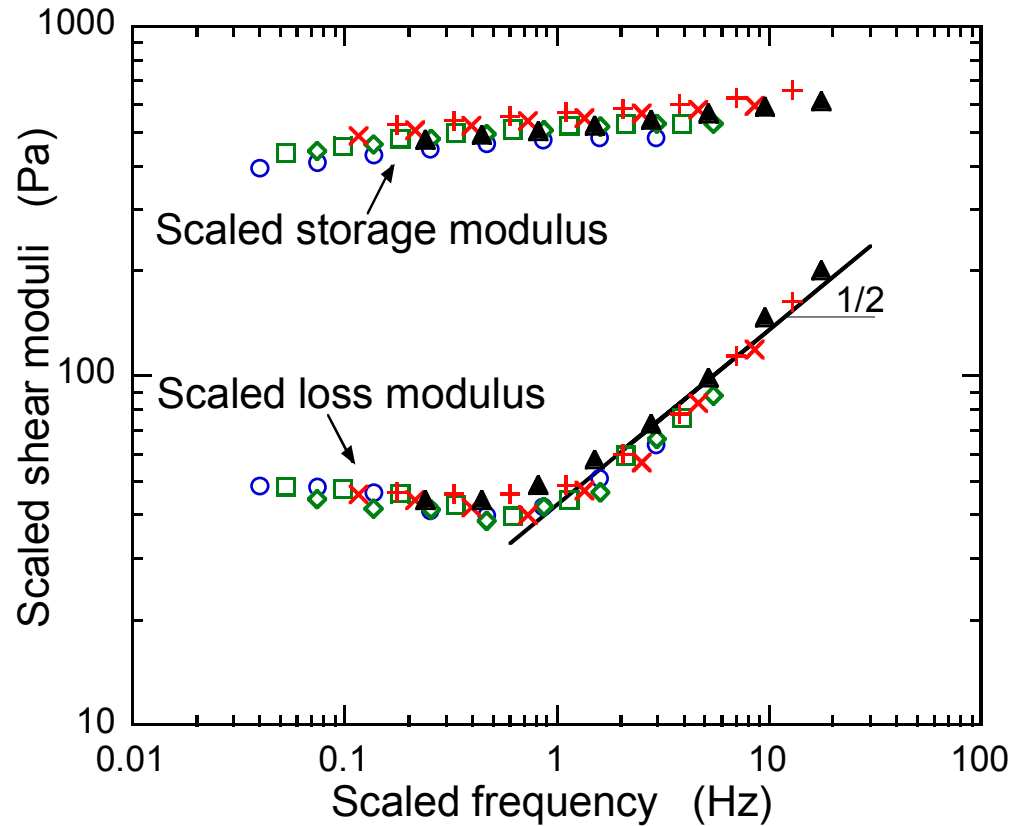
Growth law of mean bubble diameter due to Laplace pressure differences

$$d^2 = d_o^2 + K (t-t_o)$$

Mullins 1986



Scaling law for $G^(\omega, \text{age})$*



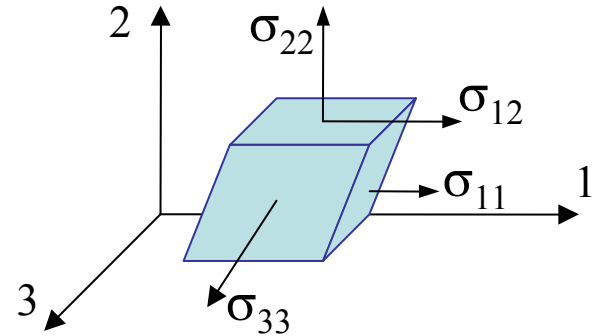
$$G^*(\omega, t) = b(t) G^*(\omega a(t), t_0)$$

C-A, Hoballah, Höhler 1998

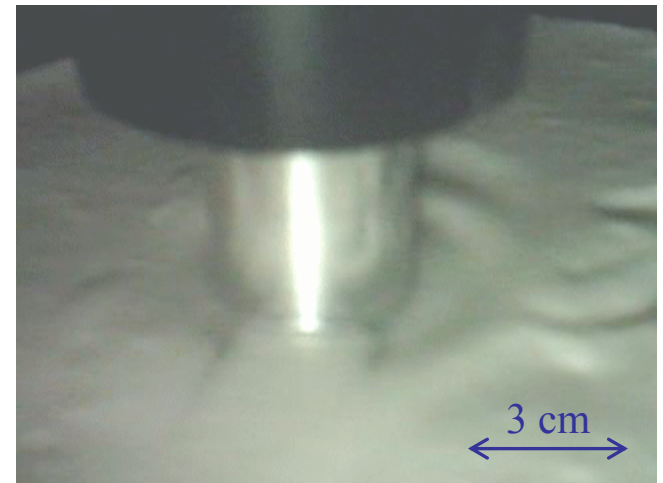
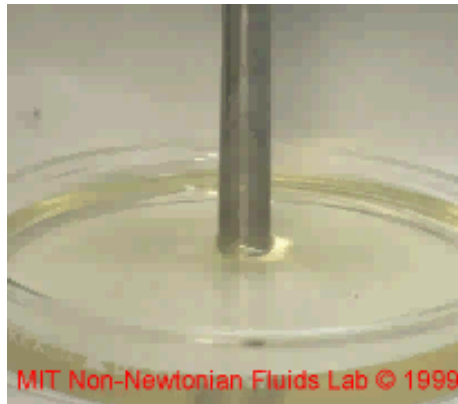
Shear induced normal stresses

$$N_1 = \sigma_{11} - \sigma_{22}$$

$$N_2 = \sigma_{22} - \sigma_{33}$$



- Stationary flow : Weissenberg effect

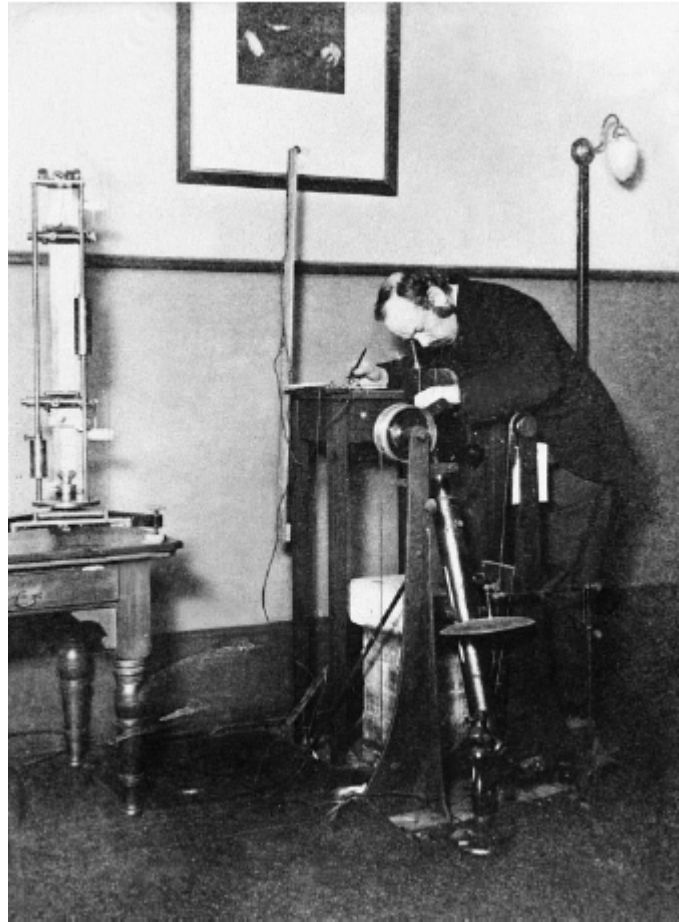


Labiausse

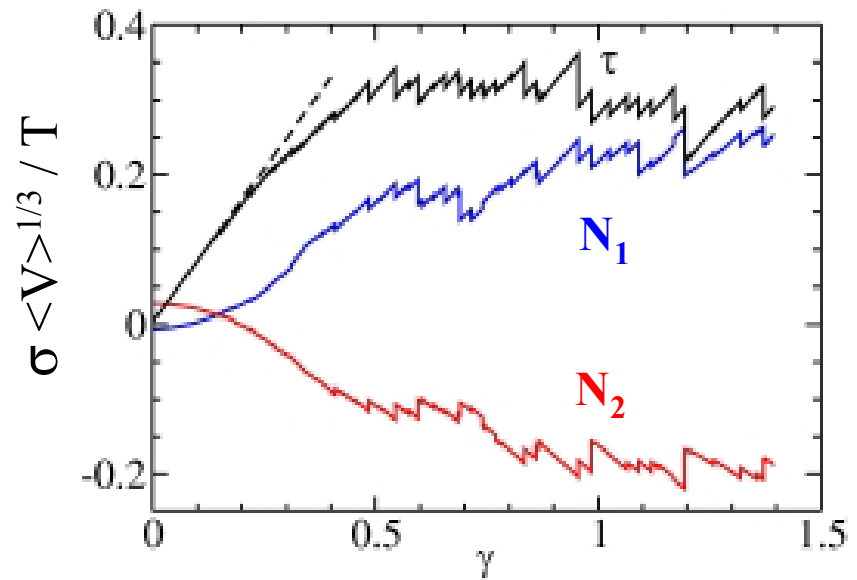
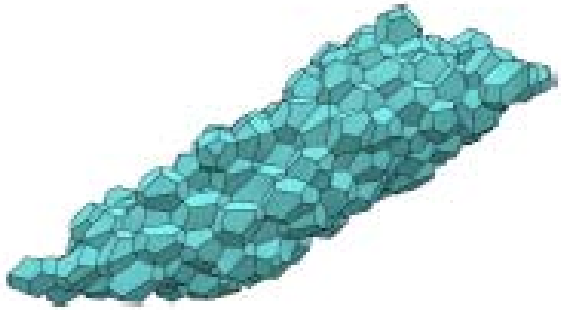
- Elastic regime : Poynting effect 1909

$$N_1 = \sigma_{12} \gamma = G \gamma^2$$

Poynting 1852-1914



Surface evolver simulations

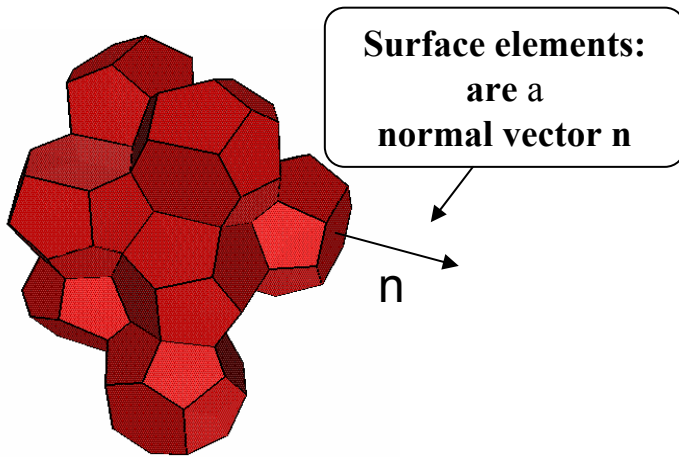


$$N_1 \cong G \gamma^2$$

$$N_2 \cong -0.85 N_1$$

Kraynik, Reinelt 2004

A nonlinear analytical constitutive law for foam



Application of a homogeneous affine strain :

Modified area a' : $a'^2 = a^2 \mathbf{n}^T \mathbf{C}^{-1} \mathbf{n}$

Energy density : $W(\mathbf{C}) = \frac{Ta}{V} \sum_j \sqrt{\mathbf{n}_j^T \mathbf{C}^{-1} \mathbf{n}_j}$

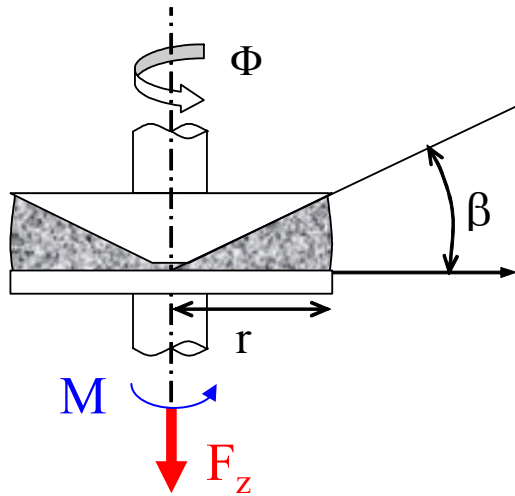
Surface tension T , sample volume V
Right Cauchy Green tensor $\mathbf{C}$

Express W in terms of the invariants of $\mathbf{C}$, denoted I_c and II_c : $W = \frac{G}{14} ((I_c - 3) + 6(II_c - 3))$

Finger tensor $\mathbf{B}$

$$\boldsymbol{\sigma} = -p \mathbf{I} + \frac{G}{7} (\mathbf{B} - 6 \mathbf{B}^{-1}) \rightarrow N_1 = \sigma_{12} \gamma \quad N_2 = -\frac{6}{7} N_1$$

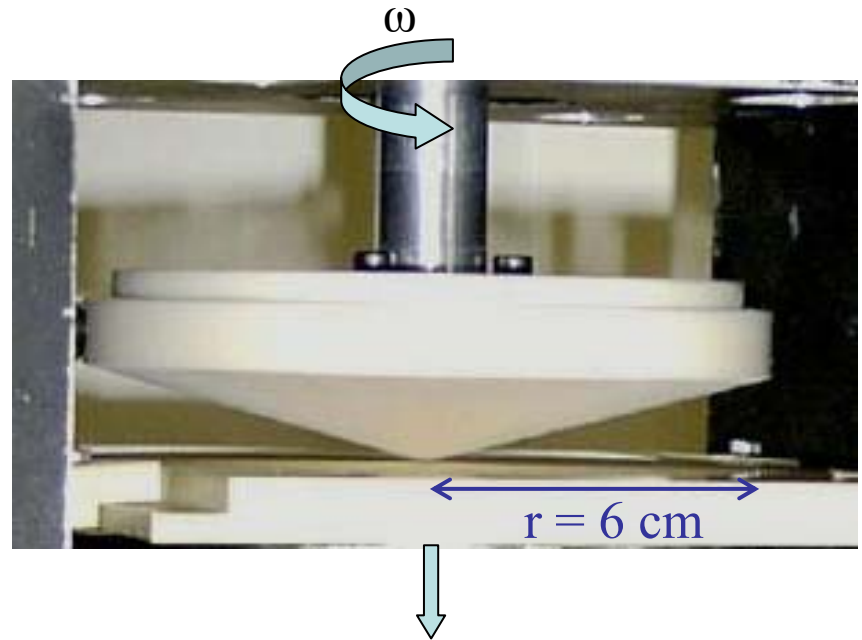
Measuring N_1



$$\gamma = \frac{\phi}{\beta}$$

$$\sigma_{\phi\theta} = \frac{3M}{2\pi r^3}$$

Macosko 1994

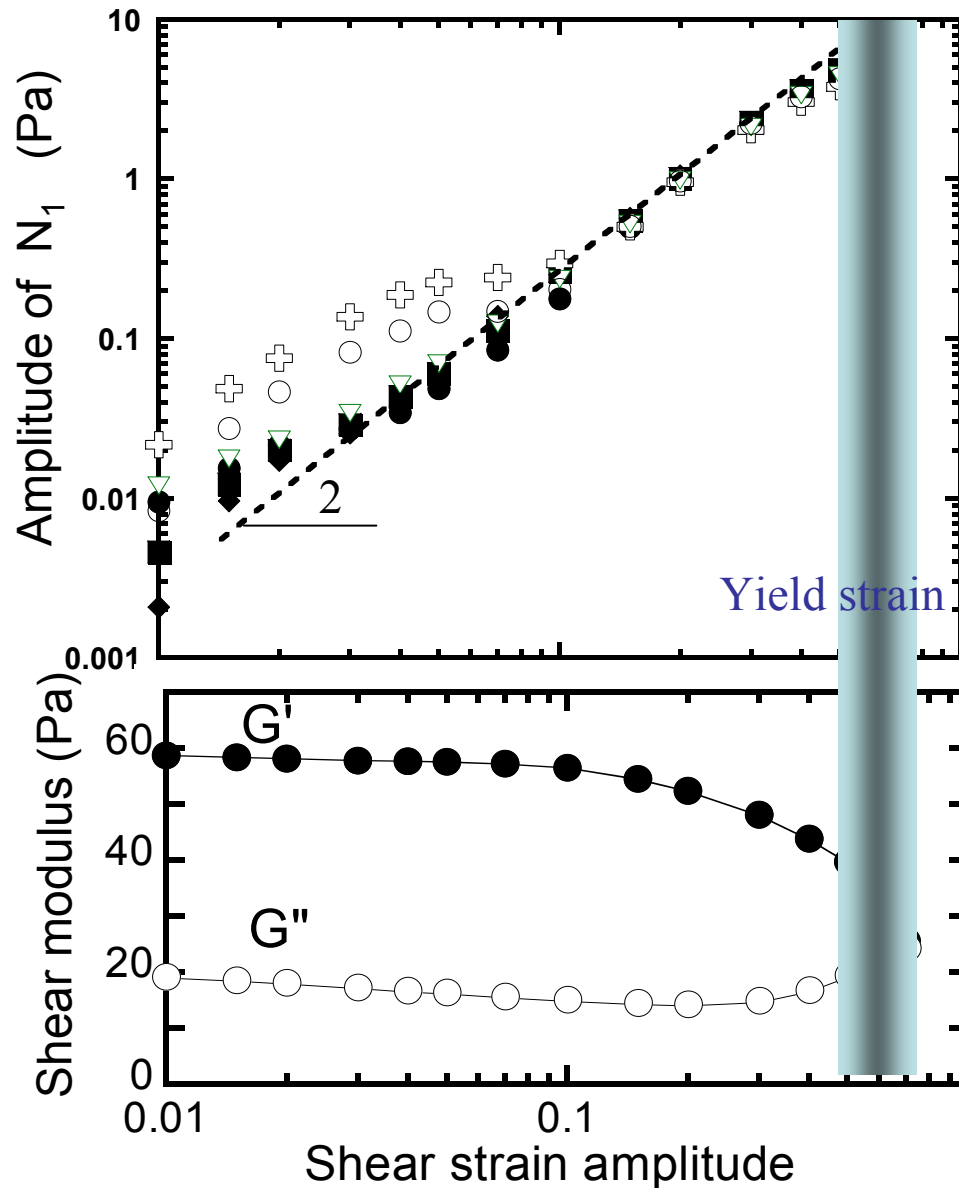


Lock-in detection at 2ω

$$N_1 = \frac{2F_z}{\pi r^2}$$

Good agreement with Poynting's law for dry foams

AOK-N₂ $\phi = 97\%$

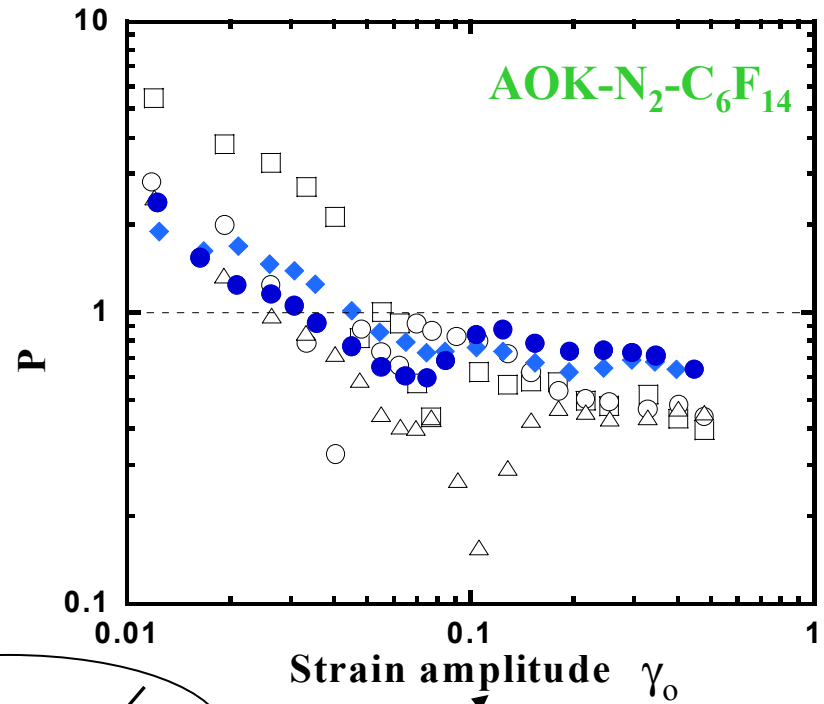
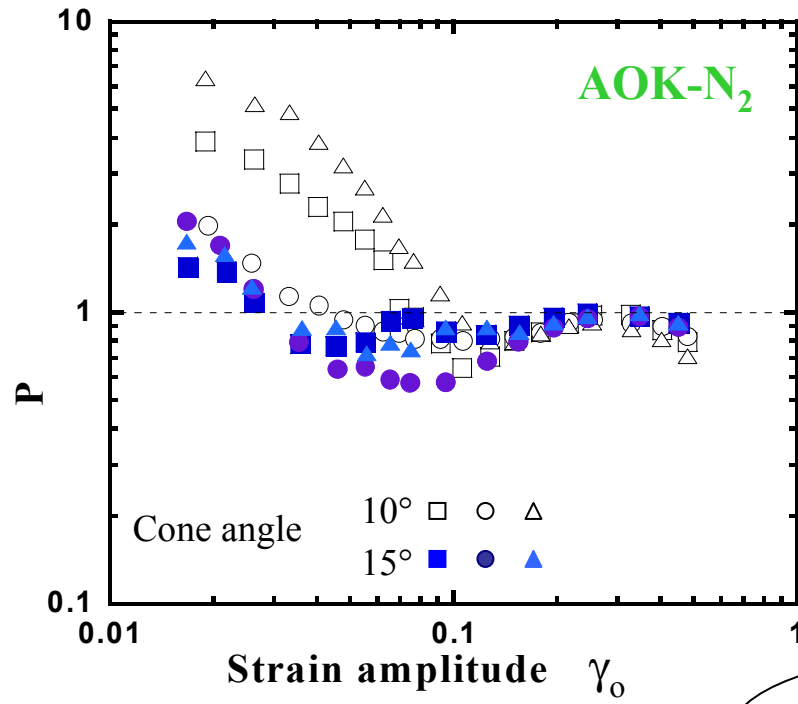


	Cone angle
Empty symbols	10°
Filled symbols	15°

frequency 1 Hz

Strain history effect

$$P \equiv \frac{2N_1}{G\gamma_o^2}$$



Coarsening rate
 / 10

Coarsening releases part of the stresses trapped due to the strain history.
 => more isotropic structure