

CLASSICAL FIELD THEORY
ASSIGNMENT

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1a,b We have a Lagrangian density

$$\mathcal{L}(\phi, \partial_\mu \phi) = \frac{1}{2} \partial_\lambda \phi \partial^\lambda \phi + \frac{1}{3} \sigma \phi^3.$$

The Euler-Lagrange equations of motion of a scalar field $\phi(x)$ are

$$\partial_\mu \left[\frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} \right] - \frac{\partial \mathcal{L}}{\partial \phi} = 0.$$

Therefore, for the above Lagrangian density we have

$$\frac{\partial}{\partial(\partial_\mu \phi)} \mathcal{L} = \frac{\partial}{\partial(\partial_\mu \phi)} \left[\frac{1}{2} \partial_\lambda \phi \partial^\lambda \phi \right]$$

$$\text{as } \frac{\partial}{\partial(\partial_\mu \phi)} \left[\frac{1}{3} \sigma \phi^3 \right] = 0. \text{ So,}$$

$$\begin{aligned} \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} &= \frac{1}{2} \frac{\partial(\partial_\lambda \phi)}{\partial(\partial_\mu \phi)} \partial^\lambda \phi + \frac{1}{2} \frac{\partial(\partial^\lambda \phi)}{\partial(\partial_\mu \phi)} \partial_\lambda \phi \\ &= \frac{1}{2} \frac{\partial(\partial_\lambda \phi)}{\partial(\partial_\mu \phi)} \partial^\lambda \phi + \frac{1}{2} \frac{\partial(\partial_\lambda \phi)}{\partial(\partial_\mu \phi)} \partial^\lambda \phi \\ &= \frac{1}{2} \delta^\mu_\lambda \partial^\lambda \phi + \frac{1}{2} \delta^\mu_\lambda \partial^\lambda \phi = \partial^\mu \phi \end{aligned}$$

and therefore,

$$\partial_\mu \left[\frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} \right] = \partial_\mu \partial^\mu \phi = \square \phi,$$

$$\frac{\partial \mathcal{L}}{\partial \phi} = \sigma \phi^2 \text{ so the E-L gives}$$

$$\square \phi - \sigma \phi^2 = 0. \quad \checkmark$$

$$T^{\mu\nu} = \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} \partial^\nu \phi - g^{\mu\nu} \mathcal{L}$$

$$= \partial^\mu \phi \partial^\nu \phi - g^{\mu\nu} \left[\frac{1}{2} \partial_\lambda \phi \partial^\lambda \phi + \frac{1}{3} \sigma \phi^3 \right].$$

d The 4-divergence of this stress tensor is

$$\partial_\mu T^{\mu\nu} = \partial_\mu (\partial^\mu \phi \partial^\nu \phi) - \partial_\mu g^{\mu\nu} \left[\frac{1}{2} \partial_\lambda \phi \partial^\lambda \phi + \frac{1}{3} \sigma \phi^3 \right]$$

$$= (\partial_\mu \partial^\mu \phi) \partial^\nu \phi + (\partial_\mu \partial^\nu \phi) \partial^\mu \phi - g^{\mu\nu} \left[\frac{1}{2} \partial_\mu (\partial_\lambda \phi \partial^\lambda \phi) + \frac{1}{3} \sigma \partial_\mu \phi^3 \right]$$

$$= (\partial_\mu \partial^\mu \phi) \partial^\nu \phi + (\partial_\mu \partial^\nu \phi) \partial^\mu \phi - g^{\mu\nu} \left[\frac{1}{2} (\partial_\mu \partial_\lambda \phi) \partial^\lambda \phi + \frac{1}{2} (\partial_\mu \partial^\lambda \phi) \partial_\lambda \phi + \frac{1}{3} \sigma \partial_\mu \phi^3 \right]$$

$$= (\partial_\mu \partial^\mu \phi) \partial^\nu \phi + (\partial_\mu \partial^\nu \phi) \partial^\mu \phi - \left[\frac{1}{2} (\partial^\nu \partial_\lambda \phi) \partial^\lambda \phi + \frac{1}{2} (\partial^\nu \partial^\lambda \phi) \partial_\lambda \phi + \frac{1}{3} \sigma \partial^\nu \phi^3 \right]$$

$$= (\partial_\mu \partial^\mu \phi) \partial^\nu \phi + (\partial_\mu \partial^\nu \phi) \partial^\mu \phi - \frac{1}{2} (\partial^\nu \partial_\lambda \phi) \partial^\lambda \phi - \frac{1}{2} (\partial^\nu \partial^\lambda \phi) \partial_\lambda \phi - \frac{1}{3} \sigma \partial^\nu \phi^3.$$

e Assuming our scalar field $\phi(x^\mu)$ obeys the equation of motion, i.e.

$$\square \phi - \sigma \phi^2 = 0,$$

$$(\partial^\nu \phi^3 = 3\phi^2 \partial^\nu \phi)$$

the stress tensor is conserved, that is $\partial_\mu T^{\mu\nu} = 0$, so

$$\begin{aligned} \partial_\mu T^{\mu\nu} &= (\partial_\mu \partial^\mu \phi) \partial^\nu \phi + (\partial_\mu \partial^\nu \phi) \partial^\mu \phi - \frac{1}{2} \delta_\lambda^\nu \sigma \phi^2 \partial^\lambda \phi - \frac{1}{2} \delta^{\nu\lambda} \sigma \phi^2 \partial_\lambda \phi - \sigma \phi^2 \partial^\nu \phi \\ &= 0. \quad \checkmark \end{aligned}$$

- a We have that the electromagnetic field tensor $F^{\mu\nu}$ transforms as

$$F'^{\mu\nu} = \Lambda^\mu_\rho \Lambda^\nu_\sigma F^{\rho\sigma}$$

with

$$F^{\mu\nu} = \begin{bmatrix} 0 & -E^1 & -E^2 & -E^3 \\ E^1 & 0 & -B^3 & B^2 \\ E^2 & B^3 & 0 & -B^1 \\ E^3 & -B^2 & B^1 & 0 \end{bmatrix}, \quad \Lambda^\mu_\rho = \begin{bmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}.$$

$$\begin{bmatrix} \gamma & -\beta\gamma & 0 & 0 \\ -\beta\gamma & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} 0 & -E^1 & -E^2 & -E^3 \\ E^1 & 0 & -B^3 & B^2 \\ E^2 & B^3 & 0 & -B^1 \\ E^3 & -B^2 & B^1 & 0 \end{bmatrix} = \begin{bmatrix} -\beta\gamma E_1 & -\gamma E_1 & -\gamma(E_2 + \beta B_3) & -\gamma(E^3 - \beta B^2) \\ E_1\gamma & \beta\gamma E_1 & \gamma(\beta E^2 - B^3) & \gamma(\beta E^3 + B^2) \\ E_2 & B_3 & 0 & -B^1 \\ E_3 & B_2 & B_1 & 0 \end{bmatrix}.$$

Multiplying again by Λ gives

$$\begin{bmatrix} 0 & -E_1 & -\gamma(E_2^2 + \beta B^3) & -\gamma(E_3^2 + \beta B^2) \\ E_1 & 0 & \gamma(E^2\beta - B^3) & -\gamma(E^3\beta + B^2) \\ \gamma(E_2^2 - \beta B^3) & \gamma(\beta^3 - \beta E^2) & 0 & -B_1 \\ \gamma(E_3^2 - \beta B^2) & \gamma(B^2 + \beta E^3) & B^1 & 0 \end{bmatrix} = F'^{\mu\nu}.$$

Thus, the electric and magnetic induction fields transform according to

$$E'_1 = E_1$$

$$B'_1 = B_1$$

$$E'_2 = \gamma E_2 - \gamma\beta B_3$$

$$B'_2 = \gamma B_2 + \gamma\beta E_3$$

$$E'_3 = \gamma E_3 + \gamma\beta B_2$$

$$B'_3 = \gamma B_3 - \gamma\beta E_2.$$

b From Jackson,

$$E'_1 = -\frac{qv't'}{r'^3}, \quad E'_2 = \frac{4b}{r'^3}, \quad E'_3 = 0$$

$$B'_1 = 0, \quad B'_2 = 0, \quad B'_3 = 0$$

with

$$r' = \sqrt{b^2 + (vt')^2}$$

where b is the distance of closest approach. Thus,

$$E'_1 = -\frac{qv't}{(b^2 + (v't)^2)^{3/2}}, \quad E'_2 = \frac{\gamma q b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}, \quad E'_3 = 0$$

as required.

$$B_3 = \gamma(B'_3 + \beta E'_2),$$

but $B'_3 = 0$ (from Jackson), so

$$B_3 = \beta E'_2 \\ = \frac{\beta \gamma q b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}.$$

We know $B_1 = B'_1$ and $B_2 = \gamma(B'_2 - \beta E'_3)$, so $B_1 = B_2 = 0$. This gives us

$$E^1 = -\frac{\gamma q v t}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}$$

$$B^1 = 0$$



$$E^2 = \frac{\gamma q b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}$$

$$B^2 = 0$$



$$E^3 = 0$$

$$B^3 = \frac{\beta \gamma q b}{(b^2 + \gamma^2 v^2 t^2)^{3/2}}.$$

