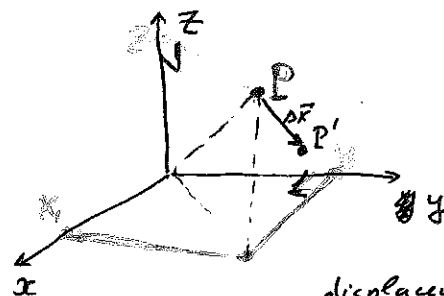


Inertial frame: a coordinate system in which $\vec{v} = \text{const}$ in time, if no forces acting on the particle.

In physics you need means to describe position objects



$$(t; x, y, z)$$

$$(t'; x', y', z')$$

$$\frac{\Delta \vec{r}}{\Delta t} = \vec{v} \leftarrow \text{velocity of the particle}$$

time to make displacement

- All inertial frames move with constant speeds $\vec{v}$ with respect to another

$$\vec{x}' = \vec{x} - \vec{v}t$$

- Galilean principle of relativity:

Physical laws are the same in any inertial frame

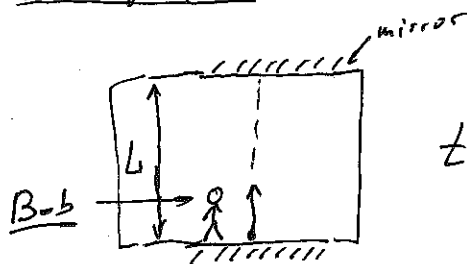
- Einstein's principle of relativity:

Physical laws are the same in any ~~inertial~~ inertial frame.

• speed of light is constant. $\leftarrow$ physical law! here is the difference.

$$c = 299792458 \text{ m/s} \approx 3 \cdot 10^8 \text{ m/s}$$

Measuring of time

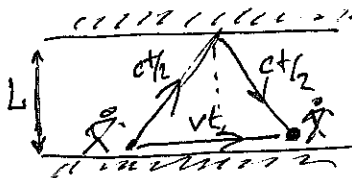


$$t = \frac{2L}{c}; \quad t_2 = \frac{L}{c}$$

Moving frame



Frame of external observer



$$(vt_2)^2 + L^2 = (ct_2)^2$$

$$t_2 = \frac{L}{\sqrt{c^2 - v^2}} = \frac{1}{\sqrt{1 - v^2/c^2}} \frac{L}{c}$$

$$= \frac{L}{c} \gamma$$

$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}}; \quad v/c = \beta$$

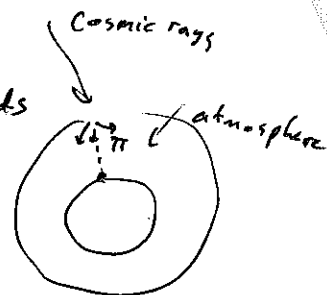
$$\gamma = \frac{1}{\sqrt{1 - \beta^2}}$$

Condition: $\frac{t_B}{2} = \frac{L}{c}$

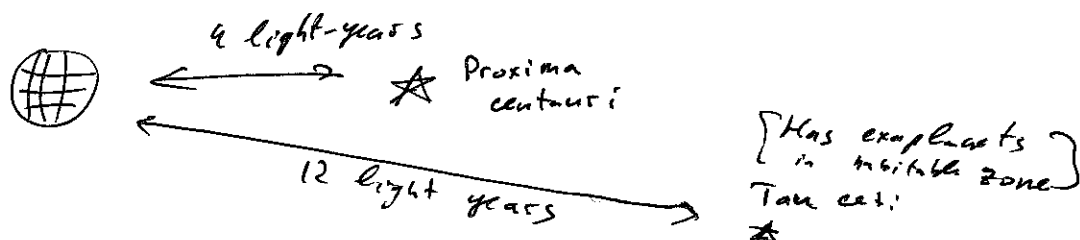
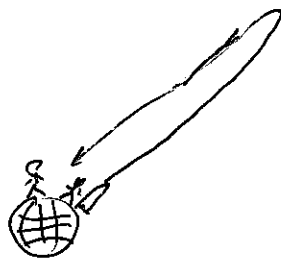
$$\frac{t_A}{2} = \gamma \cdot \frac{L}{c}, \quad \gamma \geq 1$$

Examples from real life ②

- π -meson - decays in ~ 30 nano-seconds $2.8 \cdot 10^{-8} \text{ s}$



- Twin paradox:



$$t_{\text{Earth}} = 12 \text{ light years} / \beta ; \beta = v/c = 0.99$$

$$t_{\text{rocket}} = \frac{12 \text{ light years}}{\sqrt{1-\beta^2}} = \frac{12 \text{ l.y.}}{\beta} (1 + \frac{1}{2})$$

$$t_{\text{rocket}} = t_{\text{Earth}} \sqrt{1-\beta^2} \approx t_{\text{Earth}} \sqrt{0.02} = 0.447 t_{\text{Earth}}$$

β	$\sqrt{1-\beta^2}$	time [years] to τ -eti / Earth	rocket
0.1	0.995	120	119
0.5	0.87	24	21
0.9	0.44	13	6
0.99	0.14	12	2
0.999	0.04	12	0.5

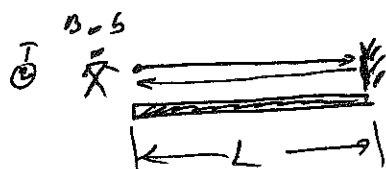
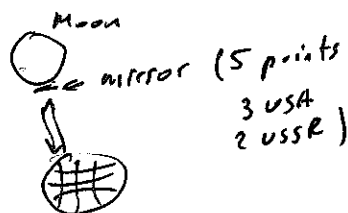
current rocket speed
speed of voyager

$$622 \text{ km/h} \quad 1.7 \cdot 10^4 \text{ m/s} \quad (62000 \frac{\text{km}}{\text{h}})$$

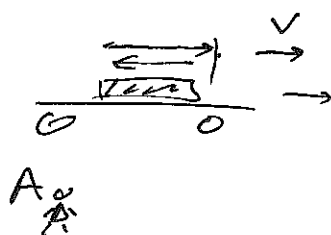
$$\beta \approx 0.5 \cdot 10^{-9}$$

Measuring of distance

③



$$L_B = \frac{1}{2} c \cdot t_B$$



$$t_{\rightarrow} = \frac{L}{c-v}$$

$$t_{\leftarrow} = \frac{L}{c+v} \quad (1)$$

$$t_{\rightarrow} + t_{\leftarrow} = t_A$$

$$\frac{t_{\rightarrow}}{t_{\leftarrow}} = \frac{c+v}{c-v}$$

$$\left(\frac{c+v}{c-v} + 1 \right) t_{\leftarrow} = t_A$$

$$t_A t_{\leftarrow} = \cancel{t_A} t_A \cdot \frac{c-v}{2c} \quad (2)$$

$$(1) \rightarrow L = \frac{c^2 - v^2}{2c} t_A = \left(1 - \frac{v^2}{c^2} \right) \frac{1}{2} c \cdot t_A$$

(1)

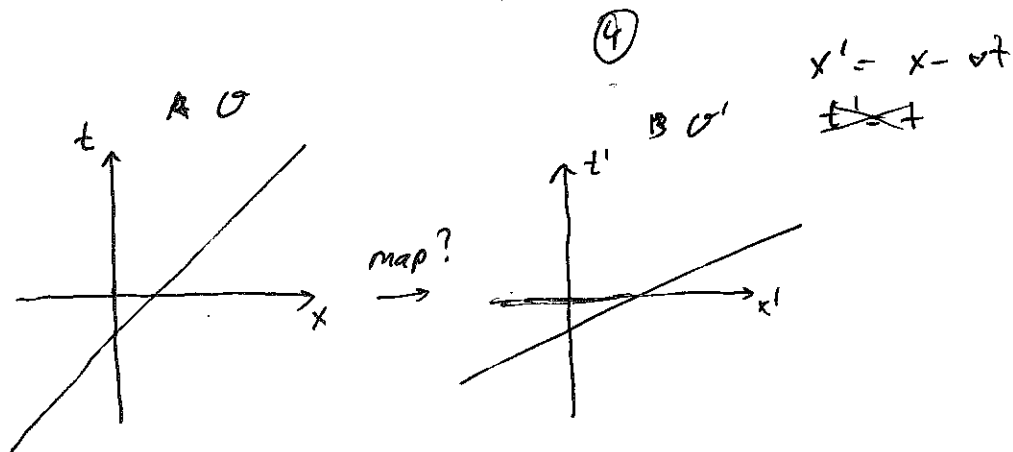
$$L_A = \frac{1}{\gamma^2} \frac{1}{2} c t_A = \frac{1}{\gamma^2} \left(\frac{1}{2} c \cdot t_B \right) \cdot \gamma = \frac{L_B}{\gamma} \quad \gamma > 1$$

$$L_A = \frac{L_B}{\gamma}$$

← Lorentz contraction.

GPS

$$t_A = \gamma t_B$$



we want that lines go to lines [principle that $\vec{v} = \text{const}$ if no forces applied]

Hence it should be a linear transformation $x \mapsto \begin{pmatrix} t' \\ x' \end{pmatrix} = \hat{A} \cdot \begin{pmatrix} t \\ x \end{pmatrix} + \vec{b}$

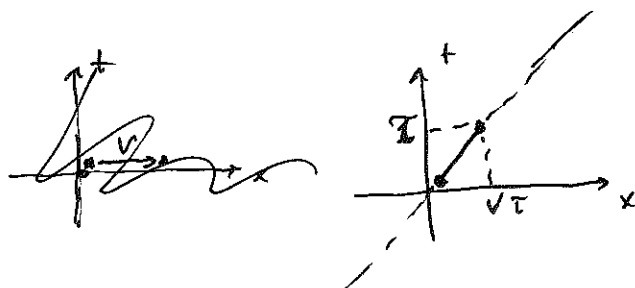
Let us choose $x \mapsto x$, $t=0 \leftrightarrow t'=0$
 $x=0 \leftrightarrow x'=0$ ← just the question is where our origin

$$\begin{pmatrix} t' \\ x' \end{pmatrix} = \hat{A} \begin{pmatrix} t \\ x \end{pmatrix} = \begin{pmatrix} A t + B x \\ C t + D x \end{pmatrix} \quad \hat{A} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

Consider first $x' = \gamma t + \delta x$

moving with velocity v means that $x = vt$

$$x' = C(-vt + x)$$



What is C ?

put $t=0$ $x' = \gamma x$ $C = \gamma$?

$\uparrow$ distance to origin $\uparrow$ distance to origin

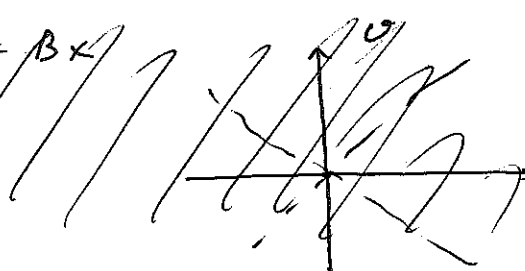
What is γ ?

$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}}$$

If O' moves with v with respect

(5)

Consider $t' = \gamma (t + \beta x)$
 put $x = 0$



Now let us prove the following statement

$x' = \gamma (x - vt)$ can. $O \rightarrow O'$

Consider opposite transform $O' \rightarrow O$

$x = \gamma (x' + vt')$

$\rightarrow t' = \gamma (t - vx)$

$$t' = \frac{1}{\gamma} \left(\frac{x}{v} - x' \right) = \frac{1}{\gamma} \left(\frac{x}{v} - \gamma x + \gamma v t \right) =$$

$$= \gamma \left(t + \frac{x}{v} \left(\frac{1}{\gamma^2} - 1 \right) \right) =$$

$$= \gamma \left(t + \frac{x}{v} \left(\frac{v^2}{c^2} \left(1 - \frac{v^2}{c^2} \right) - 1 \right) \right) =$$

$$= \gamma \left(t + \frac{vx}{c^2} \right)$$

Lorentz transformation:

$$\begin{cases} x' = \gamma (x - vt) \\ t' = \gamma \left(t - \frac{vx}{c^2} \right) \end{cases}$$

$$\begin{cases} x = \gamma (x' + \beta ct') \\ ct' = \gamma (ct - \beta x) \end{cases}$$

$$\begin{pmatrix} ct' \\ x' \end{pmatrix} = \begin{pmatrix} \gamma & -\beta\gamma \\ -\beta\gamma & \gamma \end{pmatrix} \begin{pmatrix} ct \\ x \end{pmatrix}$$

$$\vec{x}' = \gamma (\vec{x} - \vec{\beta} ct)$$

$$ct' = \gamma (ct - \vec{\beta} \cdot \vec{x})$$

⑥

Interval

Analogy: 3-dimensional space

Consider all possible ^{nongenerate} linear transformations

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = \underset{\substack{\uparrow \\ \text{3x3 matrix}}}{\hat{G}} \begin{pmatrix} x \\ y \\ z \end{pmatrix} \quad \det G \neq 0$$

Consider a subspace of transformations $\mathcal{O}$ that preserve distances:

$$\vec{V}^2 = \sum_{i=1}^3 V_i^2 = \sum_{i,j} \delta_{ij} V_i V_j$$

$$\vec{V} \cdot \vec{W} = \sum_{i,j} \delta_{ij} V_i W_j$$

$$V_i' = \mathcal{O}_{i,j}^i V_j$$

$$\delta_{i,j'} V_i' W_j' = \delta_{ij} V_i W_j$$

$$W_j' = \mathcal{O}_{j,i}^j W_i$$

$$(\delta_{i,j'} \mathcal{O}_{i,i}^i \mathcal{O}_{j,j}^j) V_i W_j = \delta_{ij} V_i W_j$$

$$\boxed{\delta_{i,j'} \mathcal{O}_{i,i}^i \mathcal{O}_{j,j}^j = \delta_{ij}}$$

Orthogonal transformations
 $\mathcal{O}(3)$

[All possible rotations $\leftarrow \det \mathcal{O} = 1$
or reflections $\leftarrow \det \mathcal{O} = -1$
+ rotations]

$$\mathcal{O} \mathcal{O}^T = \mathbb{1}$$

$$\Downarrow$$

$$1 = \det(\mathcal{O} \mathcal{O}^T) = \{\det(\mathcal{O})\}^2 = 1$$

$$\det \mathcal{O} = \pm 1$$

(7)

Minkowski space

{space-time}

1+3

4 dimensional vector space

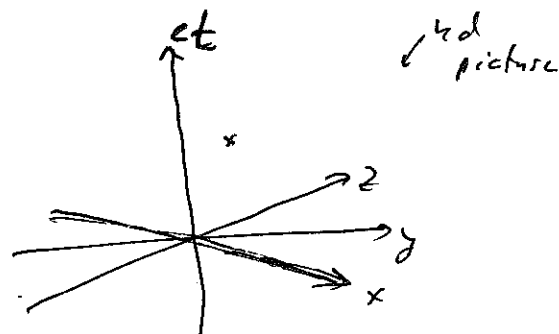
with metric

$$\eta_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

"Length"

Points in this space are
called space-time events

$$\bar{X} = \{ct, x, y, z\}$$



1+1 dimension space

"Length" between events $\{ct_2, x_2, y_2, z_2\}$

and $\{ct_1, x_1, y_1, z_1\}$

is defined by

~~$\bar{X}$~~

$$\eta_{\mu\nu}(\Delta X)^\mu (\Delta X)^\nu = \boxed{c^2(\Delta t)^2 - (\Delta \bar{X})^2} = (\Delta s)^2$$

← interval

Δs^2 can be ≥ 0 or < 0 so it is called

Consider all possible transformations that preserve interval

~~$\eta_{\mu\nu}$~~ ~~$L_\mu^{\mu'}$~~ ~~$L_{\mu'}^{\nu'}$~~

$$\boxed{O(1,3) \det L = \pm 1}$$

$$SO(1,3) \det L = 1$$

$$\eta_{\mu'\nu'} L_\mu^{\mu'} L_\nu^{\nu'} = \eta_{\mu\nu}$$

$$L \eta L^T = \eta$$

$$\det(L \eta L^T) = \det \eta = 1$$

$$\det L \cdot \det \eta \cdot \det L^T = \det \eta$$

$$\boxed{\det L = \pm 1}$$

$$\begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & & & \\ 0 & & & \\ 0 & & & \end{pmatrix}$$

← normal 3d rotations