

## EXISTENCE OF SOLUTIONS FOR AN OLDROYD MODEL OF VISCOELASTIC FLUIDS

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ABSTRACT. In this paper we investigate the unilateral problem for an Oldroyd model of a viscoelastic fluid. Using the penalty method, Faedo-Galerkin's approximation, and basic result from the theory of monotone operators, we establish the existence of weak solutions.

### 1. INTRODUCTION

It is well know that, the motion of incompressible fluids is described by the system of Cauchy equations

$$\begin{aligned} \frac{\partial u}{\partial t} + u_i \frac{\partial u}{\partial x_i} + \nabla p &= \operatorname{div} \sigma + f \\ \operatorname{div} u &= 0, \end{aligned} \tag{1.1}$$

where  $u = (u_1, \dots, u_n)$  is the velocity,  $p$  is the pressure in the fluid,  $f$  is the density of external forces and  $\sigma$  is the deviator of the stress tensor, that is,  $\sigma$  has the purpose of letting us consider reactions arising in the fluid during its motion. The vector  $(u_i \frac{\partial u_j}{\partial x_i})$ ,  $j = 1, 2, \dots, n$ , is denoted by  $(u \cdot \nabla)u$ . The Hooke's Law establishes a relationship between the stress tensor  $\sigma$  and the deformation tensor  $D_{ij}(u) = \frac{1}{2}(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i})$  and their derivatives. Therefore is the Hooke's Law that establishes the type of fluid. Such relationship is also called of *rheological equation* or *equation of state* (see Serrin [10] or Clifford [1]). For example, for an incompressible Stokes fluid the relationship has the form

$$\sigma = \alpha D + \beta D^2 \tag{1.2}$$

where  $\alpha$  and  $\beta$  are scalar functions. If in (1.2)  $\alpha = 2\nu$  positive constants. and  $\beta \equiv 0$  we have the Newton's Law  $\sigma = 2\nu D$ , which substituting in (1.1) we obtain the equations of motion of Newtonian fluid, which is called the Navier-Stokes equations:

$$u' - \nu \Delta u + (u \cdot \nabla)u + \nabla p = f, \quad \operatorname{div} u = 0,$$

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where  $\nu$  is called the kinematic coefficient of viscosity. The Navier-Stokes model was studied from the mathematical point of view by Leray [15] and later by Ladyzhenskaya [9]. We mention other deep contributions by Lions [16], Temam [21], Tartar [19] and many others researchers.

The model studied in this work, introduced by Oldroyd [11, 12], was proposed for viscous incompressible fluids whose defining equations have the form

$$(1 + \lambda \frac{\partial}{\partial t})\sigma = 2\nu(1 + k\nu^{-1} \frac{\partial}{\partial t})D, \quad (1.3)$$

where  $\lambda, \nu, k$  are positive constants with  $\nu - \frac{k}{\lambda} > 0$ . In this fluid the stress after instantaneous cessation of the motion die out like  $e^{-\lambda^{-1}t}$ , while the velocities of the flow after instantaneous removal of the stress die out like  $e^{-k^{-1}t}$ .

Assuming that  $\sigma(0) = 0$  and  $D(0) = 0$ , we write the relationship (1.3) in the form of integral equation

$$\sigma(x, t) = 2k\lambda^{-1}D(x, t) + 2\lambda^{-1}(\nu - k\lambda^{-1}) \int_0^t e^{-\frac{(t-\xi)}{\lambda}} D(x, \xi) d\xi. \quad (1.4)$$

Thus, the equation for the motion of Oldroyd fluid can be written by the system of integro-differential equations

$$\frac{\partial u}{\partial t} + (u \cdot \nabla)u - \mu \Delta u - \int_0^t \beta(t - \xi) \Delta u(x, \xi) d\xi + \nabla p = f, \quad x \in \Omega, t > 0 \quad (1.5)$$

and the incompressible condition

$$\operatorname{div} u = 0, \quad x \in \Omega, t > 0,$$

with initial and boundary conditions

$$u(x, 0) = u_0, \quad x \in \Omega, \quad \text{and} \quad u(x, t) = 0, \quad x \in \Gamma, t \geq 0.$$

Here,  $\mu = k\lambda^{-1} > 0$  and  $\beta(t) = \gamma e^{-\delta t}$ , where  $\gamma = \lambda^{-1}(\nu - k\lambda^{-1})$  with  $\delta = \lambda^{-1}$ . For physical details and mathematical modelling see [2, 5, 11, 22].

The mixed problem above was investigated by Oskolkov [2], where he proves existence of weak solution for all  $n \in \mathbb{N}$  in certain Sobolev class.

In Brézis [6] we find investigation for a unilateral problem for the case of the Navier-Stokes equations.

In the present work we consider a unilateral problem similar to Brézis [6], adding a memory term, that is  $-\int_0^t g(t-\sigma) \Delta u(\sigma) d\sigma$ . More precisely, in this paper we study a unilateral problem or a variational inequality, c.f. Lions [16], for the operator

$$L = \frac{\partial u}{\partial t} + (u \cdot \nabla)u - \mu \Delta u - \int_0^t \beta(t - \xi) \Delta u(x, \xi) d\xi + \nabla p - f$$

under standard hypothesis on  $f$  and  $u_0$ . Making use of the penalty method and Galerkin's approximations, we establish existence and uniqueness of weak solutions.

This work is organized as follows: In Section 2, we introduce the notation and main results. In Section 3, we proof to the results. Finally, in Section 4, we prove an simple result of uniqueness.

## 2. NOTATION AND MAIN RESULTS

Let  $\Omega$  be a bounded domain in  $\mathbb{R}^n$  with the boundary  $\partial\Omega$  of class  $C^2$ . For  $T > 0$ , we denote by  $Q_T$  the cylinder  $(0, T) \times \Omega$ , with lateral boundary  $\Sigma_T = (0, T) \times \partial\Omega$ . By  $\langle \cdot, \cdot \rangle$  we will represent the duality pairing between  $X$  and  $X'$ ,  $X'$  being the topological dual of the space  $X$ , and by  $C$  we denote various positive constants. We propose the variational inequality

$$\begin{aligned} u' - \mu\Delta u + (u \cdot \nabla)u - \int_0^t g(t - \sigma)\Delta u(\sigma)d\sigma + \nabla p &\geq f \quad \text{in } Q_T \\ \operatorname{div} u &= 0 \quad \text{in } Q_T \\ u &= 0 \quad \text{on } \Sigma_T \\ u(x, 0) &= u_0(x) \quad \text{in } \Omega, \end{aligned} \tag{2.1}$$

where  $g : [0, \infty) \rightarrow [0, \infty)$  is a function of  $W^{1,1}(0, \infty)$  satisfying

$$\frac{\mu}{2} - 2 \int_0^\infty g(s) ds > 0; \tag{2.2}$$

$$-C_1 g \leq g' \leq -C_2 g, \tag{2.3}$$

where  $C_1, C_2, C_3$  are positive constants;

$$g(0) > 0, \tag{2.4}$$

As an example,  $g(s) = e^{-\frac{s}{\mu}}$  satisfies the three conditions above.

To formulate problem (2.1) we need some notation about Sobolev spaces. We use standard notation of  $L^2(\Omega)$ ,  $L^p(\Omega)$ ,  $W^{m,p}(\Omega)$  and  $C^p(\Omega)$  for functions that are defined on  $\Omega$  and range in  $\mathbb{R}$ , and the notation  $L^2(\Omega)^n$ ,  $L^p(\Omega)^n$ ,  $W^{m,p}(\Omega)^n$  and  $C^p(\Omega)^n$  for functions that range in  $\mathbb{R}^n$ . Besides, we work also with the spaces  $L^p(0, T; H^m(\Omega))$  or  $L^p(0, T; H^m(\Omega))^n$ . To complete this recall on functional spaces, see for instance, Lions [16].

Also we define the following spaces

$$\mathcal{V} = \{\varphi \in \mathcal{D}(\Omega)^n : \operatorname{div} \varphi = 0\},$$

$V = V(\Omega)$  is the closure of  $\mathcal{V}$  in the space  $H_0^1(\Omega)^n$  with inner product and norm denoted, respectively by

$$((u, z)) = \sum_{i,j=1}^n \int_{\Omega} \frac{\partial u_i}{\partial x_j}(x) \frac{\partial z_i}{\partial x_j}(x) dx, \quad \|u\|^2 = \sum_{i,j=1}^n \int_{\Omega} \left( \frac{\partial u_i}{\partial x_j}(x) \right)^2 dx,$$

$H = H(\Omega)$  is the closure of  $\mathcal{V}$  in the space  $L^2(\Omega)^n$  with inner product and norm defined, respectively, by

$$(u, v) = \sum_{i=1}^n \int_{\Omega} u_i(x)v_i(x) dx, \quad |u|^2 = \sum_{i=1}^n \int_{\Omega} |u_i(x)|^2 dx$$

and  $V_2$  is the closure of  $\mathcal{V}$  in  $H^2(\Omega)^n$  with inner product and norm denoted, respectively by

$$((u, z))_{V_2} = \sum_{i=1}^n (u_i, v_i)_{H^2(\Omega)}, \quad \|u\|_{V_2}^2 = ((u, u))_{V_2},$$

**Remark 2.1.**  $V$ ,  $H$  and  $V_2$  are Hilbert's spaces,  $V_2 \hookrightarrow V \hookrightarrow H \hookrightarrow V'$  with embedding dense and continuous.

Let  $K$  be a closed and convex subset of  $V \cap V_2$  with  $0 \in K$ . We introduce the following bilinear and the trilinear forms:

$$a(u, v) = \sum_{i,j=1}^n \int_{\Omega} \frac{\partial u_i}{\partial x_j}(x) \frac{\partial v_i}{\partial x_j}(x) dx = ((u, v)),$$

$$b(u, v, w) = \sum_{i,j=1}^n \int_{\Omega} u_i(x) \frac{\partial v_j}{\partial x_i}(x) w_j(x) dx,$$

We also assume that

$$a(v, v) + b(v, \varphi, v) + \int_0^t g(t - \sigma)((v, v)) d\sigma \geq 0 \quad \forall \varphi \in K, \forall v \in V. \quad (2.5)$$

Next we shall state the main results of this paper.

**Theorem 2.2.** *If  $f \in L^2(0, T; H)$  and hypotheses (2.5) holds, then there exists a function  $u$  such that*

$$u \in L^2(0, T; V) \cap L^\infty(0, T; H) \quad (2.6)$$

$$u(t) \in K \quad a.e. \quad (2.7)$$

$$\begin{aligned} & \int_0^T \langle \varphi', \varphi - u \rangle + \mu a(u, \varphi - u) + b(u, u, \varphi - u) \\ & - \left( \int_0^t g(t - \sigma) \Delta u(\sigma) d\sigma, \varphi - u \right) dt \\ & \geq \int_0^T \langle f, \varphi - u \rangle dt, \quad \forall \varphi \in L^2(0, T; V), \varphi' \in L^2(0, T; V'), \\ & \varphi(0) = 0, \quad \varphi(t) \in K \quad a.e. \\ & u(0) = u_0. \end{aligned} \quad (2.8)$$

**Theorem 2.3.** *Assumption (2.5),  $n = 2$ , and*

$$f \in L^2(0, T; V), \quad f' \in L^2(0, T; V') \quad (2.9)$$

$$u_0 \in K. \quad (2.10)$$

Suppose also that

$$(f(0), v) - \mu a(u_0, v) - b(u_0, u_0, v) = (u_1, v) \quad \text{for all } v \in V \text{ some } u_1 \in V. \quad (2.11)$$

Then there exists a unique function  $u$  such that

$$u \in L^2(0, T; V \cap V_2), \quad u' \in L^2(0, T; V) \cap L^\infty(0, T; H) \quad (2.12)$$

$$u(t) \in K, \quad \forall t \in [0, T] \quad (2.13)$$

$$\begin{aligned} & (u'(t), v - u(t)) + \mu a(u(t), v - u(t)) + b(u(t), u(t), v - u(t)) \\ & + \int_0^T \int_0^t g(t - \sigma)((u(\sigma), v - u(t))) d\sigma dt \\ & \geq (f(t), v - u(t)) \quad \forall v \in K, \quad a.e. \text{ in } t, \end{aligned} \quad (2.14)$$

$$u(0) = u_0. \quad (2.15)$$

The proof of Theorems 2.2 and 2.3 will be given in Section 3 by the penalty method. It consists in considering a perturbation of the operator  $L$  adding a singular term called penalty, depending on a parameter  $\epsilon > 0$ . We solve the mixed problem

in  $Q$  for the penalized operator and the estimates obtained for the local solution of the penalized equation, allow to pass to limits, when  $\epsilon$  goes to zero, in order to obtain a function  $u$  which is the solution of our problem.

First of all, let us consider the penalty operator  $\beta : V \rightarrow V'$  associated to the closed convex set  $K$ , c.f. Lions [16, p. 370]. The operator  $\beta$  is monotonous, hemicontinuous, takes bounded sets of  $V$  into bounded sets of  $V'$ , its kernel is  $K$  and  $\beta : L^2(0, T; V) \rightarrow L^2(0, T; V')$  is equally monotone and hemicontinuous. The penalized problem associated with the variational inequalities (2.8) and (2.14) consists in, given  $0 < \epsilon < 1$ , find  $u_\epsilon$  satisfying

$$\begin{aligned} (u'_\epsilon, v) + \mu a(u_\epsilon, v) + b(u_\epsilon, u_\epsilon, v) - \int_0^t g(t - \sigma)(\Delta u_\epsilon(\sigma), v) d\sigma + \frac{1}{\epsilon}(\beta(u_\epsilon), v) &= (f, v), \\ \forall v \in V, \quad u_\epsilon \in L^2(0, T; V), \quad u'_\epsilon \in L^2(0, T; V') \\ u_\epsilon(x, 0) &= u_{\epsilon_0}(x). \end{aligned} \tag{2.16}$$

We suppose  $n = 2$ . The solution of this problem is given by the followings theorems.

**Theorem 2.4.** *If  $f \in L^2(0, T; H)$  and hypotheses (2.2) holds, then, for each  $0 < \epsilon < 1$  and  $u_{\epsilon_0} \in H$ , there exists a function  $u_\epsilon$  with  $u_\epsilon \in L^2(0, T; V) \cap L^\infty(0, T; H)$ ,  $u'_\epsilon \in L^2(0, T; V')$  solution of (2.16).*

**Theorem 2.5.** *If  $f \in L^2(0, T; V)$  and  $f' \in L^2(0, T; V')$  and hypotheses (2.2) holds, then for each  $0 < \epsilon < 1$  and  $u_{\epsilon_0} \in V$ , there exists a function  $u_\epsilon$  with  $u_\epsilon \in L^\infty(0, T; V \cap V_2)$ ,  $u'_\epsilon \in L^2(0, T; V) \cap L^\infty(0, T; H)$  satisfying (2.16).*

### 3. PROOF OF THE RESULTS

**Proof of Theorem 2.2.** We first prove Theorem 2.4 for the penalized problem. We employ the Faedo-Galerkin method. We note that the embedding  $V \hookrightarrow V \xrightarrow{\text{comp}} H \hookrightarrow V'$  are continuous and dense and that  $V$  is compactly and densely embedded in  $H$ . Let  $\{w_\nu, \lambda_\nu\}$ ,  $\nu \in \mathbb{N}$ , be solutions of the spectral problem

$$((w, v)) = \lambda(w, v), \quad \forall v \in V. \tag{3.1}$$

We consider  $(w_\nu)_{\nu \in \mathbb{N}}$  a Hilbertian basis for Faedo-Galerkin method. We represent by  $V_m = [w_1, w_2, \dots, w_m]$  the  $V$  subspace generated by the vectors  $w_1, w_2, \dots, w_m$  and let us consider

$$u_{\epsilon_m}(t) = \sum_{j=1}^m g_{j_m}(t) w_j$$

solution of approximate problem

$$\begin{aligned} (u'_{\epsilon_m}, w_j) + \mu a(u_{\epsilon_m}, w_j) + b(u_{\epsilon_m}, u_{\epsilon_m}, w_j) \\ - \int_0^t g(t - \sigma)(\Delta u_{\epsilon_m}(\sigma), w_j) d\sigma + \frac{1}{\epsilon} \langle \beta u_{\epsilon_m}, w_j \rangle \\ = \langle f(t), w_j \rangle, \quad j = 1, 2, \dots, m \\ u_{\epsilon_m}(x, 0) \rightarrow u_\epsilon(x, 0) \quad \text{strongly in } V. \end{aligned} \tag{3.2}$$

This system of ordinary differential equations has a solution on a interval  $[0, t_m[$ ,  $0 < t_m < T$ . The first estimate permits us to extend this solution to the whole interval  $[0, T]$ .

**Remark 3.1.** To obtain a better notation, we omit the parameter  $\epsilon$  in the approximate solutions.

**First estimate.** Multiplying both sides of (3.2) by  $g_j$  and adding from  $j = 1$  to  $j = m$ , we obtain

$$\frac{1}{2} \frac{d}{dt} |u_m(t)|^2 + \mu \|u_m(t)\|^2 + \int_0^t g(t-\sigma) (\nabla u_m(\sigma), \nabla u_m(t)) d\sigma = (f(t), u_m(t)),$$

since  $b(u_m, u_m, u_m) = 0$  (see Lions [16]) and  $(\beta u_m(t), u_m(t)) \geq 0$  because  $\beta$  is monotone and  $0 \in K$ . It follows that

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u_m(t)|^2 + \mu \|u_m(t)\|^2 \\ & \leq \left| \int_{\Omega} \nabla u_m(x, t) \left( \int_0^t g(s-\sigma) \nabla u_m(x, \sigma) \right) d\sigma dx \right|_{\mathbb{R}} \\ & \quad + |f(t)| |u_m(t)| + \int_{\Omega} |\nabla u_m(x, t)|_{\mathbb{R}} |g * \nabla u_m(x, t)|_{\mathbb{R}} dx, \end{aligned} \quad (3.3)$$

where  $*$  denotes the convolution in  $t$ . It follows from (3.3) that

$$\begin{aligned} & \frac{d}{dt} |u_m(t)|^2 + 2\mu \|u_m(t)\|^2 \\ & \leq 2 \int_{\Omega} |\nabla u_m(x, t)|_{\mathbb{R}} |g * \nabla u_m(x, t)|_{\mathbb{R}} dx + 2|f(t)| C \|u_m(t)\| \\ & = 2 \int_{\Omega} |\nabla u_m(x, t)|_{\mathbb{R}} |g * \nabla u_m(x, t)|_{\mathbb{R}} dx + 2\sqrt{\frac{2}{3\mu}} C |f(t)| \sqrt{\frac{3\mu}{2}} \|u_m(t)\| \\ & = 2 \int_{\Omega} |\nabla u_m(x, t)|_{\mathbb{R}} |g * \nabla u_m(x, t)|_{\mathbb{R}} dx + \frac{3\mu}{2} \|u_m(t)\|^2 + \frac{2}{3\mu} C^2 |f(t)|^2. \end{aligned}$$

**Remark 3.2.** We note that from Cauchy-Schwarz inequality and Fubini's theorem we have

$$\|g * \nabla u_m\|_{L^2(Q)} \leq \|g\|_{L^1(0; \infty)} \|\nabla u_m\|_{L^2(Q)}.$$

Thus, integrating 0 to  $t$  the inequality above, using the Remark 3.2 and using Gronwall's inequality we obtain

$$\|u_m\|_{L^\infty(0, T; H)}^2 + \left( \frac{\mu}{2} - 2\|g\|_{L^1(0, \infty)} \right) \|u_m\|_{L^2(0, T; V)}^2 \leq \frac{2}{3\mu} C + C^2 |f|_{L^2(0, T; H)}^2.$$

Integrating these last inequality in  $t \in [0, T]$  and using (2.2), we have

$$u_m \text{ is bounded in } L^\infty(0, T; H) \quad (3.4)$$

$$u_m \text{ is bounded in } L^2(0, T; V). \quad (3.5)$$

From (3.5), we obtain

$$\beta(u_m) \text{ is bounded in } L^2(0, T; V') \quad (3.6)$$

**Second estimate.** By Remark 3.2, we observe that,

$$\text{if } \xi \in L^2(0, T; H) \text{ then } \int_0^t g(t-\sigma) \xi(\sigma) d\sigma \in L^2(0, T; H). \quad (3.7)$$

Similarly we obtain

$$\int_0^t g(t - \sigma)\xi(\sigma)d\sigma \in V \quad \text{if } \xi(t) \in V, \tag{3.8}$$

$$\int_0^t g(t - \sigma)\xi(\sigma)d\sigma \in V' \quad \text{if } \xi(t) \in V'. \tag{3.9}$$

We consider  $\tilde{u}_m = u_m, \tilde{w} = w$  in  $[0, T]$  and  $\tilde{u}_m = 0, \tilde{w} = 0$  out of  $[0, T], \tilde{g}(\xi) = g(\xi)$  if  $\xi \geq 0$  and zero if  $\xi < 0$ . Therefore,  $\nabla\tilde{u}_m \in L^2(\mathbb{R}; H), \tilde{w} \in L^2(\mathbb{R}; V)$  and  $\tilde{g} \in L^1(\mathbb{R})$ . This implies

$$\begin{aligned} & \int_0^T \int_0^t g(t - \sigma) ((u_m(\sigma), w(t))) \, d\sigma \, dt \\ &= \int_{\mathbb{R}} \int_{\mathbb{R}} \tilde{g}(t - \sigma) \int_{\Omega} \nabla\tilde{u}_m(x, \sigma) \nabla\tilde{w}(x, t) \, dx \, d\sigma \, dt \\ &= \int_{\mathbb{R}} \int_{\Omega} \tilde{g} * \nabla\tilde{u}_m(x, t) \nabla\tilde{w}(x, t) \, dx \, dt \\ &= \int_{\mathbb{R}} \int_{\Omega} \nabla\tilde{u}_m(x, \sigma) \tilde{g} * \nabla\tilde{w}(x, \sigma) \, dx \, d\sigma, \end{aligned}$$

where  $\tilde{g}(x) = \tilde{g}(-x)$ . We observe that (3.5) implies that

$$\int_0^T ((u_m(t), w))dt \rightarrow \int_0^T ((u(t), w))dt, \quad \forall w \in L^2(0, T; V). \tag{3.10}$$

From (3.7), we have that  $\tilde{g} * \tilde{w}(t) \in V, \forall w \in L^2(0, T; V)$ , therefore (3.10) yield

$$\int_{\mathbb{R}} (\nabla\tilde{u}_m(\sigma), \tilde{g} * \nabla\tilde{w}(\sigma)) \, dt \rightarrow \int_{\mathbb{R}} (\nabla\tilde{u}(\sigma), \tilde{g} * \nabla\tilde{w}(\sigma)) \, dt.$$

We note that

$$\begin{aligned} \int_{\mathbb{R}} (\nabla\tilde{u}(\sigma), \tilde{g} * \nabla\tilde{w}(\sigma)) \, dt &= \int_{\mathbb{R}} (\tilde{g} * \nabla\tilde{u}(\sigma), \nabla\tilde{w}(\sigma)) \, dt \\ &= \int_{\Omega} \int_{\mathbb{R}} \int_{\mathbb{R}} \tilde{g}(t - \sigma) \nabla\tilde{u}(x, \sigma) d\sigma \nabla\tilde{w}(x, t) d\sigma \, dt \, dx \\ &= \int_0^T \int_0^t g(t - \sigma) ((u(\sigma), w(t))) \, d\sigma \, dt. \end{aligned}$$

Then

$$\int_0^T \int_0^t g(t - \sigma) ((u_m(\sigma), w(t))) \, d\sigma \, dt \rightarrow \int_0^T \int_0^t g(t - \sigma) ((u(\sigma), w(t))) \, d\sigma \, dt, \tag{3.11}$$

for all  $w \in L^2(0, T; V)$ .

Let  $P_m$  be the orthogonal projection  $H \mapsto V_m$ ; that is,

$$P_m \varphi = \sum_{j=1}^m (\varphi, w_j) w_j, \quad \varphi \in H.$$

By the choice of  $(w_\nu)_{\nu \in \mathbb{N}}$  we have

$$\|P_m\|_{\mathcal{L}(V, V)} \leq 1 \quad \text{and} \quad \|P_m^*\|_{\mathcal{L}(V', V')} \leq 1.$$

We note that  $P_m u'_m = u'_m$ . Multiplying both sides of the approximate equation (3.2) by the vector  $w_j$  and adding from  $j = 1$  to  $j = m$ , we obtain using the notations and ideas of Lions [16, pages 75-76] and (3.7), that

$$(u'_m) \text{ is bounded in } L^2(0, T; V'). \quad (3.12)$$

The boundedness in (3.5), (3.12) and the Aubin-Lions compactness Theorem imply that there exists a subsequence from  $(u_m)$ , still denoted by  $(u_m)$ , such that

$$u_m \rightarrow u \text{ strongly in } L^2(0, T; H) \text{ and a. e. in } Q. \quad (3.13)$$

Returning to the notation  $u_{\epsilon_m}$ , using (3.4), (3.5) and (3.13) (see Lions [16, pages 76-77]), (3.6) and (3.11) we obtain

$$\begin{aligned} (u'_\epsilon, v) + a(u_\epsilon, v) + b(u_\epsilon, u_\epsilon, v) - \int_0^t g(t - \sigma)(\Delta u(\sigma), v) d\sigma + \frac{1}{\epsilon}(\zeta, v) &= (f, v), \\ \forall v \in V, \quad u_\epsilon \in L^2(0, T; V), \quad u'_\epsilon \in L^2(0, T; V') \\ u_\epsilon(x, 0) &= u_{\epsilon_0}(x). \end{aligned} \quad (3.14)$$

It is necessary to prove that  $\zeta = \beta(u_\epsilon)$ . We make this using the monotony of the operator  $\beta$  (see Lions [16, Chap. 2]). Therefore, we have proved the Theorem 2.4.

**Proof of Theorem 2.2.** From (3.4), (3.5), (3.13) and Banach-Steinhaus theorem, it follows that there exists a subnet  $(u_\epsilon)_{0 < \epsilon < 1}$ , such that it converges to  $u$  as  $\epsilon \rightarrow 0$ , in the weak sense. This function satisfies (2.6). On the other hand, we have from (3.14) that

$$\beta u_\epsilon = \epsilon [f - u'_\epsilon - Au_\epsilon - Bu_\epsilon - \int_0^t g(t - \sigma) \Delta u(\sigma) d\sigma]. \quad (3.15)$$

Where  $\langle Au_\epsilon, v \rangle = a(u, v)$  and  $\langle Bu_\epsilon, v \rangle = b(u_\epsilon, u_\epsilon, v)$ .

Since  $\int_0^t g(t - \sigma) \Delta u(\sigma) d\sigma \in V'$  and  $[f - u'_\epsilon - Au_\epsilon - Bu_\epsilon]$  is bounded, we have

$$\beta u_\epsilon \rightarrow 0 \text{ in } \mathcal{D}'(0, T; V'). \quad (3.16)$$

Since  $\beta u_\epsilon$  is bounded in  $L^2(0, T; V')$ , we have

$$\beta u_\epsilon \rightarrow 0 \text{ weak in } L^2(0, T; V'). \quad (3.17)$$

On the other hand we deduce from (3.14) that

$$0 \leq \int_0^T \langle \beta u_\epsilon, u_\epsilon \rangle dt \leq \epsilon C. \quad (3.18)$$

Thus  $\int_0^T \langle \beta u_\epsilon, u_\epsilon \rangle dt \rightarrow 0$ . We have that

$$\int_0^T \langle \beta u_\epsilon - \beta \varphi, u_\epsilon - \varphi \rangle dt \geq 0, \quad \forall \varphi \text{ in } L^2(0, T; V),$$

because  $\beta$  is a monotonous operator. Thus,

$$\int_0^T \langle \beta u_\epsilon, u_\epsilon \rangle dt - \int_0^T \langle \beta u_\epsilon, \varphi \rangle dt - \int_0^T \langle \beta \varphi, u_\epsilon - \varphi \rangle dt \geq 0. \quad (3.19)$$

We have from (3.17) and (3.19) that

$$\int_0^T \langle \beta \varphi, u(t) - \varphi \rangle dt \leq 0. \quad (3.20)$$

Taking  $\varphi = u - \lambda v$ , with  $v \in L^2(0, T; V)$  and  $\lambda > 0$ , we deduce using the hemicontinuity of  $\beta$  that

$$\beta(u(t)) = 0, \quad (3.21)$$

and this implies that  $u(t) \in K$  a. e.

Next, we prove that  $u$  is a solution of inequality (2.8). Let us consider  $\mathbf{X}_\epsilon$  defined by

$$\begin{aligned} \mathbf{X}_\epsilon &= \int_0^T \langle \varphi', \varphi - u_\epsilon \rangle dt + \int_0^T a(u_\epsilon, \varphi - u_\epsilon) dt + \int_0^T b(u_\epsilon, u_\epsilon, \varphi - u_\epsilon) dt \\ &\quad + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon(\sigma), \varphi - u_\epsilon)) d\sigma dt, - \int_0^T \langle f, \varphi - u_\epsilon \rangle dt, \end{aligned} \quad (3.22)$$

with  $\varphi \in L^2(0, T; V)$ ,  $\varphi' \in L^2(0, T; V')$ ,  $\varphi(0) = 0$ ,  $\varphi(t) \in K$  a.e. It follows from (3.22) that

$$\begin{aligned} \mathbf{X}_\epsilon &= \int_0^T \langle \varphi', \varphi \rangle dt - \int_0^T \langle \varphi', u_\epsilon \rangle dt + \int_0^T a(u_\epsilon, \varphi) dt - \int_0^T a(u_\epsilon, u_\epsilon) dt \\ &\quad + \int_0^T b(u_\epsilon, u_\epsilon, \varphi) dt - \int_0^T b(u_\epsilon, u_\epsilon, u_\epsilon) dt + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon(\sigma), \varphi)) d\sigma dt \\ &\quad - \int_0^T \int_0^t g(t - \sigma)((u_\epsilon(\sigma), u_\epsilon)) d\sigma dt - \int_0^T \langle f, \varphi \rangle dt + \int_0^T \langle f, u_\epsilon \rangle dt. \end{aligned} \quad (3.23)$$

On the other hand, taking  $v = \varphi - u_\epsilon$  in (2.16) and integrating in  $Q_T$ , we obtain that

$$\begin{aligned} & - \int_0^T \langle u'_\epsilon, \varphi \rangle dt + \int_0^T \langle u'_\epsilon, u_\epsilon \rangle dt - \int_0^T a(u_\epsilon, \varphi) dt + \int_0^T a(u_\epsilon, u_\epsilon) dt \\ & - \int_0^T b(u_\epsilon, u_\epsilon, \varphi) dt + \int_0^T b(u_\epsilon, u_\epsilon, u_\epsilon) dt - \int_0^T \int_0^t g(t - \sigma)((u_\epsilon(\sigma), \varphi)) d\sigma dt \\ & + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon(\sigma), u_\epsilon)) d\sigma dt - \frac{1}{\epsilon} \int_0^T \langle \beta u_\epsilon - \beta \varphi, \varphi - u_\epsilon \rangle dt \\ & + \int_0^T \langle f, \varphi \rangle dt - \int_0^T \langle f, u_\epsilon \rangle dt = 0, \end{aligned} \quad (3.24)$$

because  $\beta \varphi = 0$ . Adding member to member (3.23) and (3.24), we obtain

$$\begin{aligned} \mathbf{X}_\epsilon &= \int_0^T \langle \varphi', \varphi \rangle dt - \int_0^T \langle \varphi', u_\epsilon \rangle dt - \int_0^T \langle u'_\epsilon, \varphi \rangle dt \\ &\quad + \int_0^T \langle u'_\epsilon, u_\epsilon \rangle dt + \frac{1}{\epsilon} \int_0^T \langle \beta \varphi - \beta u_\epsilon, \varphi - u_\epsilon \rangle dt \geq 0, \end{aligned} \quad (3.25)$$

because

$$\begin{aligned} & \int_0^T \langle \varphi', \varphi \rangle dt - \int_0^T \langle \varphi', u_\epsilon \rangle dt - \int_0^T \langle u'_\epsilon, \varphi \rangle dt + \int_0^T \langle u'_\epsilon, u_\epsilon \rangle dt \\ & = \int_0^T \langle \varphi' - u'_\epsilon, \varphi - u_\epsilon \rangle dt \geq 0. \end{aligned}$$

On the other hand,  $b(u_\epsilon, u_\epsilon, u_\epsilon) = 0$ . From (3.22)-(3.23) it follows that

$$\begin{aligned} \mathbf{X}_\epsilon &= \int_0^T \langle \varphi', \varphi - u_\epsilon \rangle dt + \int_0^T a(u_\epsilon, \varphi) dt \\ &\quad + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, \varphi)) d\sigma dt - \int_0^T \langle f, \varphi - u_\epsilon \rangle dt \\ &\geq \int_0^T a(u_\epsilon, u_\epsilon) dt + \int_0^T b(u_\epsilon, \varphi, u_\epsilon) dt + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, u_\epsilon)) d\sigma dt. \end{aligned} \quad (3.26)$$

Consider

$$\mathbf{Y}_\epsilon = \int_0^T a(u_\epsilon, u_\epsilon) dt + \int_0^T b(u_\epsilon, \varphi, u_\epsilon) dt + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, u_\epsilon)) d\sigma dt. \quad (3.27)$$

It follows from (2.5) with  $v = u - u_\epsilon$  that

$$a(u - u_\epsilon, u - u_\epsilon) + b(u - u_\epsilon, \varphi, u - u_\epsilon) + \int_0^t g(t - \sigma)((u - u_\epsilon, u - u_\epsilon)) d\sigma \geq 0.$$

On the other hand, we can write

$$\begin{aligned} \mathbf{Y}_\epsilon &= \int_0^T a(u_\epsilon - u, u_\epsilon - u) dt + \int_0^T b(u_\epsilon - u, \varphi, u_\epsilon - u) dt \\ &\quad + \int_0^T a(u, u_\epsilon - u) dt + \int_0^T a(u_\epsilon, u) dt + \int_0^T b(u, \varphi, u_\epsilon - u) dt \\ &\quad + \int_0^T b(u_\epsilon, \varphi, u) dt + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon - u, u_\epsilon - u)) d\sigma dt \\ &\quad + \int_0^T \int_0^t g(t - \sigma)((u, u_\epsilon - u)) d\sigma dt + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, u)) d\sigma dt. \end{aligned}$$

This implies

$$\begin{aligned} \mathbf{Y}_\epsilon &\geq \int_0^T a(u_\epsilon, u) dt + \int_0^T a(u, u_\epsilon - u) dt + \int_0^T b(u, \varphi, u_\epsilon - u) dt \\ &\quad + \int_0^T b(u_\epsilon, \varphi, u) dt + \int_0^T \int_0^t g(t - \sigma)((u, u_\epsilon - u)) d\sigma dt \\ &\quad + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, u)) d\sigma dt. \end{aligned} \quad (3.28)$$

Taking lim sup in (3.28) we obtain

$$\limsup \mathbf{Y}_\epsilon \geq \int_0^T a(u, u) dt + \int_0^T b(u, \varphi, u) dt + \int_0^T \int_0^t g(t - \sigma)((u, u)) d\sigma dt. \quad (3.29)$$

It follows from (3.26) and (3.29) that

$$\begin{aligned} &\limsup \left\{ \int_0^T \langle \varphi', \varphi - u_\epsilon \rangle dt + \int_0^T a(u_\epsilon, \varphi) dt \right. \\ &\quad \left. + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, u)) d\sigma dt - \int_0^T \langle f, \varphi - u_\epsilon \rangle dt \right\} \\ &\geq \int_0^T a(u, u) dt + \int_0^T b(u, \varphi, u) dt + \int_0^T \int_0^t g(t - \sigma)((u, u)) d\sigma dt. \end{aligned} \quad (3.30)$$

It follows from (3.30) that

$$\begin{aligned} & \int_0^T \langle \varphi', \varphi - u \rangle dt + \int_0^T a(u, \varphi - u) dt + \int_0^T b(u, u, \varphi - u) dt \\ & + \int_0^T \int_0^t g(t - \sigma) \langle (u, \varphi - u) \rangle d\sigma dt \\ & \geq \int_0^T \langle f, \varphi - u \rangle dt \end{aligned}$$

for all  $\varphi \in L^2(0, T; V)$ ,  $\varphi' \in L^2(0, T; V')$ ,  $\varphi(0) = 0$ ,  $\varphi(t) \in K$  a.e.

**Proof of Theorem 2.3.** We first prove Theorem 2.5 for the penalized problem. As in the proof of Theorem 2.2, we employ the Faedo-Galerkin Method. Let  $(w_\nu)_{\nu \in \mathbb{N}}$  be a Hilbertian basis of  $V$ . By  $V_m = [w_1, w_2, \dots, w_m]$  we represent the subspace generated by the  $m$  first vectors of  $(w_\nu)$ . Consider

$$u_{\epsilon_m} = \sum_{j=1}^m g_{jm} w_j$$

solution of approximate penalized problem

$$\begin{aligned} & (u'_{\epsilon_m}, w_j) + \mu a(u_{\epsilon_m}, w_j) + b(u_{\epsilon_m}, u_{\epsilon_m}, w_j) \\ & - \int_0^t g(t - \sigma) \langle \Delta u_{\epsilon_m}(\sigma), v \rangle d\sigma + \frac{1}{\epsilon} \langle \beta u_{\epsilon_m}, w_j \rangle \\ & = \langle f(t), w_j \rangle, \quad j = 1, 2, \dots, m \\ & u_{\epsilon_m}(x, 0) \rightarrow u_\epsilon(x, 0) \quad \text{strongly in } V. \end{aligned} \tag{3.31}$$

**First estimate.** As in the proof of Theorem 2.4, omitting the parameter  $\epsilon$  and taking  $v = u_m$  in the approximate equation (3.31) we obtain

$$(u_m) \quad \text{is bounded in } L^\infty(0, T; H), \tag{3.32}$$

$$(u_m) \quad \text{is bounded in } L^2(0, T; V), \tag{3.33}$$

**Second estimate.** In both sides of (3.31) we take the derivatives with respect  $t$  and consider  $v = u'_m(t)$ . We obtain

$$\begin{aligned} & (u''_m(t), u'_m(t)) + \mu a(u'_m(t), u'_m(t)) \\ & + b(u'_m(t), u_m(t), u'_m(t)) + b(u_m(t), u'_m(t), u'_m(t)) \\ & + \frac{1}{\epsilon} ((\beta u_m(t))', u'_m(t)) + \int_0^t g'(t - \sigma) \langle (u_m(t), u'_m(t)) \rangle d\sigma \\ & + g(0) \langle (u_m(t), u'_m(t)) \rangle + \frac{1}{\epsilon} ((\beta u_m)'(t), u'_m(t)) \\ & = (f'(t), u'_m(t)), \end{aligned} \tag{3.34}$$

because

$$\frac{d}{dt} \left( \int_0^t g(t - \sigma) \Delta u_m(\sigma) d\sigma \right) = g(0) \Delta u_m(t) + \int_0^t g'(t - \sigma) \Delta u_m(\sigma) d\sigma.$$

We note that

$$\begin{aligned} u'_m(0) & \rightarrow u_1 \quad \text{strongly in } H, \\ u_m(0) & \rightarrow u_0 \quad \text{strongly in } V. \end{aligned} \tag{3.35}$$

Indeed, (3.35)<sub>1</sub> is obtained using (3.31) with  $t = 0$  and (2.11). Note that  $\beta(u_0) = 0$ . Then

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u'_m(t)|^2 + \mu \|u'_m(t)\|^2 + b(u'_m(t), u_m(t), u'_m(t)) \\ & + \int_0^t g'(t - \sigma) ((u_m(t), u'_m(t))) d\sigma + g(0) \frac{1}{2} \frac{d}{dt} \|u_m(t)\|^2 \\ & = (f'(t), u'_m(t)), \end{aligned} \quad (3.36)$$

because  $b(u_m(t), u'_m(t), u'_m(t)) = 0$  and  $((\beta u_m)'(t), u'_m(t)) \geq 0$  (see Lions [16, page 399]).

**Remark 3.3.** The derivative with respect to  $t$  of  $(\beta(v(t)), w)$  is only formal. The correct method is to consider the difference equation in  $t+h$  and  $t$ , divided by  $h$  and take the limits when  $h \rightarrow 0$ . Here is fundamental the operator  $\beta$  to be monotonous. This justify the formal procedure of taking the derivative with respect to  $t$ , on both sides of (3.31) and take  $v = u'_m(t)$ . See Brezis [6], Browder [8] or Lions [17] for details.

As  $n = 2$ , we have (see Lions [16, page 70])

$$\|u\|_{L^4(\Omega)}^2 \leq C \|u\| \|u\|, \quad \forall u \in H_0^1(\Omega). \quad (3.37)$$

Moreover,  $H_0^1(\Omega) \hookrightarrow L^2(\Omega)$ ; therefore,

$$|b(u, v, u)| \leq \sum_{i,j=1}^2 \int_{\Omega} |u_i(x)| \left| \frac{\partial v_j}{\partial x_i}(x) \right| |u_j(x)| \leq \|u\|_{(L^4(\Omega))^2} \|v\|.$$

This and (3.37) imply

$$\begin{aligned} |b(u'_m(t), u_m(t), u'_m(t))| & \leq \sqrt{\frac{\mu}{2}} \|u'_m(t)\| C \sqrt{\frac{2}{\mu}} \|u'_m(t)\| \|u_m(t)\| \\ & \leq \frac{\mu}{4} \|u'_m(t)\|^2 + \frac{C^2}{\mu} \|u_m(t)\|^2 \|u'_m(t)\|^2. \end{aligned} \quad (3.38)$$

Therefore,

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u'_m(t)|^2 + \mu \|u'_m(t)\|^2 + g(0) \frac{1}{2} \frac{d}{dt} \|u_m(t)\|^2 \\ & \leq \left| \int_0^t g'(t - \sigma) ((u_m(t), u'_m(t))) d\sigma \right|_{\mathbb{R}} + |b(u'_m(t), u_m(t), u'_m(t))|_{\mathbb{R}} \\ & + \sqrt{\frac{2}{\mu}} \|f'(t)\|_{V'} \sqrt{\frac{\mu}{2}} \|u'_m(t)\|, \end{aligned} \quad (3.39)$$

Therefore, from (3.39), (3.37) and Remark 3.2 we obtain

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} |u'_m(t)|^2 + g(0) \frac{1}{2} \frac{d}{dt} \|u_m(t)\|^2 + \mu \|u'_m(t)\|^2 \\ & \leq \frac{\mu}{4} \|u'_m(t)\|^2 + \frac{C^2}{2\mu} \|u_m(t)\| \|u'_m(t)\|^2 \\ & + \|g'\|_{L^1(0,\infty)} \|u'_m(t)\| \|u_m(t)\| + \frac{C_1}{\mu} \|f'(t)\|^2 + \frac{\mu}{4} \|u'_m(t)\|^2, \end{aligned} \quad (3.40)$$

Integrating (3.40) from 0 to  $t$  and using the hypothesis (2.3), (2.4) we obtain

$$\begin{aligned} & |u'_m(t)|^2 + \left(\frac{\mu}{2} - 2\|g\|_{L^1(0,\infty)}\right) \int_0^t \|u'_m(s)\|^2 ds \\ & \leq C_2^2 \|g\|_{L^1(0,\infty)} \int_0^T \|u_m(t)\|^2 dt + C \int_0^t \|u_m(s)\|^2 |u'_m(s)|^2 ds + C \int_0^T \|f'(t)\|^2 dt. \end{aligned} \tag{3.41}$$

From (3.4) and hypothesis on  $f$  we obtain

$$|u'_m(t)|^2 + \left(\frac{\mu}{2} - 2\|g\|_{L^1(0,\infty)}\right) \int_0^t \|u'_m(s)\|^2 ds \leq C + C \int_0^t \|u_m(s)\|^2 |u'_m(s)|^2 ds. \tag{3.42}$$

Being  $(u_m)$  is bounded in  $L^2(0, T; V)$  we have, using Gronwall's inequality in (3.41) and hypothesis  $H1$ , that

$$(u'_m) \text{ is bounded in } L^2(0, T; V) \tag{3.43}$$

$$(u'_m) \text{ is bounded in } L^\infty(0, T; H). \tag{3.44}$$

**Third estimate.** Let  $(w_\nu)$  be the orthonormal system of  $V \cap V_2$  formed by the eigenfunctions of the Laplace operator.

As in the proof of Theorem 2.4, omitting the parameter  $\epsilon$  and taking  $w_j = -\Delta u_m$  in the approximate equation (3.31) we obtain

$$\begin{aligned} & \frac{1}{2} \frac{d}{dt} \|u_m(t)\|^2 + \mu |\Delta u_m(t)|^2 \\ & \leq |b(u_m(t), u_m(t), -\Delta u_m(t))|_{\mathbb{R}} \\ & \quad + \left| \int_{\Omega} \Delta u_m(x, t) \left( \int_0^t g(t - \sigma) \Delta u_m(x, \sigma) d\sigma \right) \right|_{\mathbb{R}} dx \\ & \quad + \frac{1}{\mu} |f(t)|^2 + \frac{\mu}{4} |\Delta u_m(t)|^2 \end{aligned} \tag{3.45}$$

because  $\langle \beta u_m, -\Delta u_m \rangle \geq 0$  (see Haraux [4, page 58]). We note that

$$\begin{aligned} |b(u_m(t), u_m(t), -\Delta u_m(t))| & \leq \sum_{i,j=1}^2 \int |u_{m_j}(t)| \left| \frac{\partial u_{m_j}}{\partial x_i}(t) \right| |\Delta u_{m_j}(t)| \\ & \leq \|u_m(t)\|_{(L^3(\Omega))^2}^2 \left\| \frac{\partial u_m}{\partial x_i}(t) \right\| |\Delta u_m(t)|, \end{aligned} \tag{3.46}$$

because  $H_0^1(\Omega) \hookrightarrow L^3(\Omega)$ ,  $H_0^1(\Omega) \hookrightarrow L^6(\Omega)$ , with  $\frac{1}{2} + \frac{1}{3} + \frac{1}{6} = 1$ .

Substituting (3.46) in (3.45) and using the Remark 3.2, we obtain

$$\begin{aligned} & \frac{d}{dt} \|u_m(t)\|^2 + \left(\frac{\mu}{2} - 2\|g\|_{L^1(0,\infty)}\right) |\Delta u_m(t)|^2 \\ & \leq C \|f(t)\|^2 + \left(C \|u_m(t)\|_{(L^3(\Omega))^2}^2 - C\right) \|u_m(t)\|^2. \end{aligned}$$

Integrating the above inequality from 0 to  $t$ , observing that  $u_m \in L^2(0, T; V) \subset L^2(0, T; (L^3(\Omega))^2)$  and using the Gronwall's Lemma, it follows that

$$u_m \text{ is bounded in } L^\infty(0, T; V) \tag{3.47}$$

$$u_m \text{ is bounded in } L^2(0, T; V_2). \tag{3.48}$$

To complete the proof of Theorem 2.5, we use the same argument used in the proof of Theorem 2.4.

We shall now prove Theorem 2.3. From the previous convergence, and Banach-Steinhaus theorem, it follows that there exists a subnet  $(u_\epsilon)_{0 < \epsilon < 1}$ , such that it converges to  $u$  as  $\epsilon \rightarrow 0$ , in the sense of previous convergence.

This function satisfies (2.12) and (2.13). Using the same arguments used in Theorem 2.2 we obtain that  $\beta u = 0$ . Therefore,  $u$  satisfy (2.15) of Theorem 2.3.

We need to show only that  $u$  is a solution of inequality (2.14) a.e. in  $t$ . In fact, we have that  $u_\epsilon$  satisfies

$$\begin{aligned} (u'_\epsilon, \tilde{v}) + \mu a(u_\epsilon, \tilde{v}) + b(u_\epsilon, u_\epsilon, \tilde{v}) + \int_0^t g(t - \sigma)((u_\epsilon, \tilde{v}))d\sigma + \frac{1}{\epsilon}(\beta u_\epsilon, \tilde{v}) &= (f, \tilde{v}), \\ u_\epsilon(0) &= u_0. \end{aligned} \quad (3.49)$$

for all  $\tilde{v} \in V$ . Then from (3.49), with  $\tilde{v} = v - u_\epsilon$ ,  $v \in K$ , we have

$$\begin{aligned} (u'_\epsilon, v - u_\epsilon) + \mu a(u_\epsilon, v) + b(u_\epsilon, u_\epsilon, v) + \int_0^t g(t - \sigma)((u_\epsilon, v))d\sigma - (f, v - u_\epsilon) \\ \geq \mu a(u_\epsilon, u_\epsilon) + \int_0^t g(t - \sigma)((u_\epsilon, u_\epsilon))d\sigma, \quad \forall v \in K, \end{aligned} \quad (3.50)$$

because  $(\beta u_\epsilon - \beta v, u_\epsilon - v) \geq 0$ . Let us denote

$$X_\epsilon^v = (u'_\epsilon, v - u_\epsilon) + \mu a(u_\epsilon, v) + b(u_\epsilon, u_\epsilon, v) + \int_0^t g(t - \sigma)((u_\epsilon, v))d\sigma - (f, v - u_\epsilon).$$

We obtain

$$X_\epsilon^v \geq \mu a(u_\epsilon, u_\epsilon) + \int_0^t g(t - \sigma)((u_\epsilon, u_\epsilon))d\sigma, \quad \forall v \in V. \quad (3.51)$$

Let  $\psi \in C^0([0, T])$  with  $\psi(t) \geq 0$ . Then  $v\psi \in C^0([0, T]; V)$  for all  $v \in V$ .

$$u_{\epsilon i} u_{\epsilon j} \rightarrow u_i u_j \text{ weakly in } L^2(0, T, L^2(\Omega))$$

It follows from (3.51) that

$$\begin{aligned} \int_0^T \psi(u'_\epsilon, v - u_\epsilon) dt + \mu \int_0^T \psi a(u_\epsilon, v) dt + \int_0^T \psi b(u_\epsilon, u_\epsilon, v) dt \\ + \psi \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, v))d\sigma dt - \int_0^T \psi(f, v - u_\epsilon) dt \\ \geq \mu \int_0^T \psi a(u_\epsilon, u_\epsilon) dt + \int_0^T \int_0^t g(t - \sigma)((u_\epsilon, u_\epsilon))d\sigma dt. \end{aligned} \quad (3.52)$$

Taking lim sup in both side of inequality (3.52) we obtain

$$\begin{aligned} \int_0^T \psi(u', v - u) dt + \mu \int_0^T \psi a(u, v) dt - \int_0^T \psi b(u, u, v) dt \\ + \int_0^T \int_0^t g(t - \sigma)((u, v))d\sigma dt - \int_0^T \psi(f, v - u) dt \\ \geq \mu \int_0^T \psi a(u, u) dt + \int_0^T \int_0^t g(t - \sigma)((u, u))d\sigma dt, \end{aligned} \quad (3.53)$$

because

$$\limsup \mu \int_0^T \psi a(u_\epsilon, u_\epsilon) dt \geq \liminf \mu \int_0^T \psi a(u_\epsilon, u_\epsilon) dt \geq \mu \int_0^T \psi a(u, u) dt$$

and

$$\begin{aligned} \limsup \int_0^T \int_0^t g(t-\sigma)((u_\epsilon, u_\epsilon)) d\sigma dt &= \limsup \int_0^T \int_0^t g(t-\sigma)(-\Delta u_\epsilon, u_\epsilon) d\sigma dt \\ &= \int_0^T \int_0^t g(t-\sigma)(-\Delta u, u) d\sigma dt \\ &= \int_0^T \int_0^t g(t-\sigma)((u, u)) d\sigma dt \end{aligned}$$

From (3.53) we obtain finally

$$\begin{aligned} (u', v-u) + \mu a(u, v-u) + b(u, u, v-u) + \int_0^t g(t-\sigma)((u, v-u)) d\sigma \\ \geq (f, v-u) \quad \forall v \in K, \text{ a.e. in } t. \end{aligned} \quad (3.54)$$

#### 4. UNIQUENESS

We now prove that when  $n = 2$  we have uniqueness in Theorem 2.3. Indeed, suppose that  $u_1, u_2$  are two solutions of (2.14) and set  $w = u_2 - u_1$  and  $t \in (0, T)$ . Taking  $v = u_1$  (resp.  $u_2$ ) in the inequality (2.14) relative to  $v_2$  (resp.  $v_1$ ) and adding up the results we obtain

$$\begin{aligned} - \int_0^t (w', w) dt - \mu \int_0^t a(w, w) dt + \int_0^t b(u_1, u_1, w) dt \\ - \int_0^t b(u_2, u_2, w) dt - \int_0^t \int_0^t g(t-\sigma)((w, w)) \geq 0. \end{aligned}$$

Therefore,

$$\frac{1}{2} \int_0^t \frac{d}{dt} |w(t)|^2 dt + \mu \int_0^t \|w(t)\|^2 dt \leq \int_0^t |b(w, u_2, w)| dt, \quad (4.1)$$

because  $\int_0^t \int_0^t g(t-\sigma)((w, w)) \geq 0$  and  $b(u_2, u_2, w) - b(u_1, u_1, w) = b(w, u_2, w)$ . On the other hand, if  $n = 2$ , we have (see Lions [16, page 70])

$$|b(w(t), u_2(t), w(t))| \leq C \|w(t)\| \|w(t)\| \|u_2(t)\|. \quad (4.2)$$

It follows from (4.1) and (4.2) that

$$|w(t)|^2 + \frac{\mu}{2} \int_0^t \|w(t)\|^2 dt \leq C \int_0^t |w(t)|^2 \|u_2(t)\|^2 dt.$$

This implies, using Gronwall's inequality that  $w = 0$ , because  $u_2 \in L^2(0, T; V)$ , therefore  $u_1(t) = u_2(t)$ , for all  $t \in [0, T]$ .

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