2013 Dolomites Research Week on Approximation

Lecture 4: Global and local kernel methods for approximating derivatives on the sphere



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- Methods currently in use for computational geosciences
- Transport equation on the sphere
- Global method RBF collocation method
 Numerical examples and comparisons to other methods
- RBF finite difference method (RBF-FDM)
 o Numerical examples
- RBF partition of unity method (RBF-PUM)
 o Numerical examples and comparison to RBF-FDM
- Concluding remarks

Bottom line for numerical methods: Need numerical methods that provide highresolution and accuracy at low computational costs to resolve the multi-scale features.

Grids, meshes, nodes, used in large scale models

• Grids/meshes/nodes used used in methods for large-scale applications:



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- Methods used:
 - Finite-difference, finite-element, finite-volume, semi-Lagrangian
 - Double Fourier, spherical harmonics, spectral elements, discontinuous Galerikin (DG), and radial basis functions (RBF)

Highlights of some high-order methods in large-scale models

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Double Fourier series:

- Strength: Exponential accuracy Computationally fast because of FFTs
- Weakness: No practical option for local mesh refinement

Spherical harmonics:

- Strength: Exponential accuracy
- Weakness: No practical option for local mesh refinement Relatively high computational cost Poor scalability on parallel computer architectures

Spectral elements:

- Strength: Accuracy approaching exponential Local refinement is feasible but complex
- Weakness: Loss of efficiency due to unphysical element boundaries (Runge phenomenon - oscillations near boundaries → restrictive time-step) High algorithmic complexity High pre-processing cost





(Surface) Div, Grad, Curl, and all that

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	Spherical Coords.	Cartesian Coords.
Point:	$(\lambda, heta, 1)$	(x,y,z)
Unit vectors:	$\hat{\mathbf{i}} = $ longitudinal	$\hat{\mathbf{i}} = x$ -direction
	$\hat{\mathbf{j}} = $ latitudinal $\hat{\mathbf{k}} = $ radial	$\mathbf{\hat{j}} = y$ -direction $\mathbf{\hat{k}} = z$ -direction
Unit tangent vectors:	î, ĵ	$oldsymbol{\zeta} = rac{1}{\sqrt{1-z^2}} egin{bmatrix} -y \ x \ 0 \end{bmatrix}, \ oldsymbol{\mu} = rac{1}{\sqrt{1-z^2}} egin{bmatrix} -zx \ -zy \ 1-z^2 \end{bmatrix}$
Unit normal vector:	ĥ	$\mathbf{x} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}} + z\hat{\mathbf{k}}$
Gradient of scalar g :	$\mathbf{u}_s = abla_s \; g = rac{1}{\cos heta} rac{\partial g}{\partial \lambda} \mathbf{\hat{i}} + rac{\partial g}{\partial heta} \mathbf{\hat{j}}$	$\mathbf{u}_{c}=P(abla_{c}~g)=P\left(rac{\partial g}{\partial x}\mathbf{\hat{i}}+rac{\partial g}{\partial y}\mathbf{\hat{j}}+rac{\partial g}{\partial z}\mathbf{\hat{k}} ight)$
Surface divergence of u :	$ abla_s \cdot \mathbf{u}_s = rac{1}{\cos heta} rac{\partial u_s}{\partial \lambda} + rac{\partial v_s}{\partial heta}$	$(abla_c) \cdot \mathbf{u}_c = abla_c \cdot \mathbf{u}_c - \mathbf{x} \cdot abla (\mathbf{u}_c \cdot \mathbf{x})$
Curl of a scalar f :	$\mathbf{u}_s = \hat{\mathbf{k}} imes (abla_s f) = -rac{\partial f}{\partial heta} \hat{\mathbf{i}} + rac{1}{\cos heta} rac{\partial f}{\partial \lambda} \hat{\mathbf{j}}$	$\mathbf{u}_c = \mathbf{x} imes (P abla_c f) = Q P(abla_c f) = Q(abla_c f)$
Surface curl of a vector u :	$\hat{\mathbf{k}} \cdot (abla_s imes \mathbf{u}_s) = - abla_s \cdot (\hat{\mathbf{k}} imes \mathbf{u}_s)$	$\mathbf{x} \cdot ((P abla_c) imes \mathbf{u}_c) = - abla_c \cdot (Q \mathbf{u}_c)$
Here: P	$= I - \mathbf{x}\mathbf{x}^T = egin{bmatrix} 1 - x^2 & -xy & -xy & -xy & -xy & -xy & 1-y^2 & -xz & -yz & 1 \ -xz & -yz & 1 \ \end{pmatrix}$	$egin{array}{c} -xz \ -yz \ -z^2 \end{array} ight] Q = egin{bmatrix} 0 & -z & y \ z & 0 & -x \ -y & x & 0 \end{array} ight]$

Shallow water wave equations on a rotating sphere

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- Model for the nonlinear dynamics of a shallow, hydrostatic, homogeneous, and inviscid fluid layer.



• Idealized test-bed for horizontal dynamics of all 3-D global climate models.

Equations	Momentum	Transport		
Spherical coordinates	$\frac{\partial \mathbf{u}_s}{\partial t} + \mathbf{u}_s \cdot \nabla_s \mathbf{u}_s + f \hat{\mathbf{k}} \times \mathbf{u}_s + g \nabla_s h = 0$	$\frac{\partial h^*}{\partial t} + \nabla_s \cdot (h^* \mathbf{u}_s) = 0$		
	Singu	larity at poles!		
Cartesian coordinates	$\frac{\partial \mathbf{u}_c}{\partial t} + P \begin{bmatrix} (\mathbf{u}_c \cdot P\nabla_c)u_c + f(\mathbf{x} \times \mathbf{u}_c) \cdot \hat{\mathbf{i}} + g(\mathbf{x} - \mathbf{u}_c) \cdot \hat{\mathbf{j}} +$	$ \begin{bmatrix} P\hat{\mathbf{i}} \cdot \nabla_c h \\ P\hat{\mathbf{j}} \cdot \nabla_c h \\ P\hat{\mathbf{k}} \cdot \nabla_c h \end{bmatrix} = 0 \frac{\partial h^*}{\partial t} + (P\nabla_c) \cdot (h^* \mathbf{u}_c) = 0 $		
	Smooth	over entire sphere!		

Simpler problem: transport equation on the sphere

- For this tutorial we focus on the transport equation for a scalar valued quantity h on the surface of the unit sphere in an incompressible velocity field **u**.
- The governing PDE can be written in Cartesian coordinates as:

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$$h_t + \mathbf{u} \cdot (P\nabla h) = 0$$

P projects arbitrary three-dimensional vectors onto a plane tangent to the unit sphere at \mathbf{x} .

• Surface gradient operator:

$$P\nabla = (\mathbf{I} - \mathbf{x}\mathbf{x}^T)\nabla = \begin{bmatrix} (1 - x^2) & -xy & -xz \\ -xy & (1 - y^2) & -yz \\ -xz & -yz & (1 - z^2) \end{bmatrix} \begin{bmatrix} \partial_x \\ \partial_y \\ \partial_z \end{bmatrix} = \begin{bmatrix} \mathbf{p}_x \cdot \nabla \\ \mathbf{p}_y \cdot \nabla \\ \mathbf{p}_z \cdot \nabla \end{bmatrix}$$

• Goal: Show how to construct good numerical approximations to

$$\mathcal{D}_x = \mathbf{p}_x \cdot \nabla, \ \mathcal{D}_y = \mathbf{p}_y \cdot \nabla, \ \mathcal{D}_z = \mathbf{p}_z \cdot \nabla$$

Surface gradient approximation: Global RBF method

- Setup: $X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$ and $f\Big|_X$ samples of a target function.
- ϕ is some differentiable PD or CPD(1) kernel on \mathbb{R}^3 .
- RBF interpolant of $f|_X$ is given by

$$s(\mathbf{x}) = \sum_{j=1}^{N} c_j \phi(\|\mathbf{x} - \mathbf{x}_j\|)$$



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• The coefficients c_j are determined from:

$$\underbrace{\begin{bmatrix} \phi(\|\mathbf{x}_{1} - \mathbf{x}_{1}\|) & \phi(\|\mathbf{x}_{1} - \mathbf{x}_{2}\|) \cdots \phi(\|\mathbf{x}_{1} - \mathbf{x}_{N}\|) \\ \phi(\|\mathbf{x}_{2} - \mathbf{x}_{1}\|) & \phi(\|\mathbf{x}_{2} - \mathbf{x}_{2}\|) \cdots \phi(\|\mathbf{x}_{2} - \mathbf{x}_{N}\|) \\ \vdots & \vdots & \ddots & \vdots \\ \phi(\|\mathbf{x}_{N} - \mathbf{x}_{1}\|) & \phi(\|\mathbf{x}_{N} - \mathbf{x}_{2}\|) \cdots \phi(\|\mathbf{x}_{N} - \mathbf{x}_{N}\|) \end{bmatrix}} \underbrace{\begin{bmatrix} c_{1} \\ c_{2} \\ \vdots \\ c_{N} \end{bmatrix}}_{\underline{C}} = \underbrace{\begin{bmatrix} f_{1} \\ f_{2} \\ \vdots \\ f_{N} \end{bmatrix}}_{\underline{f}}$$

• Discretization of the projected gradient closely follows Flyer & W (2007,2009).

Surface gradient approximation: Global RBF method DRWA 2013 Lecture 4

• Approximate the *x*-component of the surface gradient using collocation:

$$\mathbf{p}_{x} \cdot \nabla s(\mathbf{x}) \left| \left| \mathbf{x} = \mathbf{x}_{j} \right|_{\mathbf{x} = \mathbf{x}_{j}} = \sum_{k=1}^{N} c_{k} \left[\mathbf{p}_{x} \cdot \nabla \phi_{k} (\|\mathbf{x} - \mathbf{x}_{k}\|) \right] \right|_{\mathbf{x} = \mathbf{x}_{j}}, \quad j = 1, \dots, N$$

$$= \sum_{k=1}^{N} c_{k} \left[x_{j} \mathbf{x}_{j}^{T} \mathbf{x}_{k} - x_{k} \right] \left(\frac{\phi_{k}^{'}(\|\mathbf{x} - \mathbf{x}_{k}\|)}{\|\mathbf{x} - \mathbf{x}_{k}\|} \right) \right|_{\mathbf{x}_{j}}, \quad j = 1, \dots, N$$

$$= B^{x} \underline{c}$$

$$= \left(B^{x} A^{-1} \right) \underline{f}$$

$$= D_{N}^{x} \underline{f}$$

- D_N^x is an N-by-N differentiation matrix (DM).
- It represents the discrete RBF approximation to $\mathbf{p}_x \cdot \nabla$ at nodes X.
- DMs D_N^y and D_N^z can similarly be constructed fo $(\mathbf{p}_y \cdot \nabla)$ and $(\mathbf{p}_z \cdot \nabla)$.

Global RBF collocation for transport equation

• Continuous transport equation for some $\mathbf{u} = (u, v, w) \in T_X \mathbb{S}^2$:

$$h_t + \mathbf{u} \cdot (P\nabla h) = 0$$

- Let h and $\mathbf{u} = (u, v, w)$ be sampled at X.
- Semi-discrete formulation (method-of-lines) of transport equation:

$$\underline{h}_t = -\left(\operatorname{diag}(\underline{u})D_N^x + \operatorname{diag}(\underline{v})D_N^y + \operatorname{diag}(\underline{w})D_N^z\right)\underline{h} = -D_N\underline{h}.$$

- Advance the system in time using some standard ODE solver.
- This is a purely hyperbolic problem and temporal stability can be an issue.
 - We stabilize the method by including some high-order diffusion operator L_N (hyperviscosity):

$$\underline{h}_t = -D_N \underline{h} + \mu L_N \underline{h}$$

 $-L_N$ is a discrete approximation to a high power of the Laplacian: Δ^{2k} .

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• Solid body rotation of a non-smooth cosine bell (Williamson et. al. JCP (1992))

Stream Function for flow

$$\psi(\mathbf{x}) = \cos(\alpha)z + \sin(\alpha)y \qquad \alpha = \pi/2$$
 (flow over the poles)

Initial condition (non-smooth: jump in second derivative)





Details:

- Gaussian RBF
- $\Delta t = 30$ minutes
- No stabilization required.
- Minimum energy node sets used.

- Convergence results as number of nodes N increases (Flyer & W, 2007)
- Error results are for one complete revolution of the cosine bell.



log-log scale

Straight line indicates <u>algebraic accuracy</u>

Straight line indicates **spectral accuracy**

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• Comparison to other high order methods (Flyer & W, 2007)

Method	Cost per	ℓ_2 error	Time-step	Number of	Code length	Local mesh
	time-step			grid points	(# of lines)	refinement
RBF	$O(N^2)$	0.006	1/2 hour	4096	< 40	yes
SH	$O(M^{3/2})$	0.005	90 seconds	32768	> 500	no
DF	$O(N \log N)$	0.005	90 seconds	32768	> 100	no
DG	$O(kN_e)$	0.005	6 minutes	7776	> 1000	yes

RBF=radial basis functions, SH=spherical harmonics, DF=double Fourier, DG=discontinuous Galerkin spectral elements

Comments:

- For RBF and DF N = the number of grid points.
- For SH $M = \text{total number of spherical harmonics: } (85+1)^2 = 7396.$
- For DG $N_e = \text{total number of nodes per element, and } k=\text{number of elements.}$

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RBF=radial basis functions, SH=spherical harmonics, DF=double Fourier, DG=discontinuous Galerkin spectral elements

- Need ways to reduce this cost.
- Next two methods we discussed are focused on this

RBF generated finite differences (RBF-FD)

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- Generalization of finite-difference (FD) method to scattered nodes using RBFs to compute the FD weights.
- References:
 - \circ W & Fornberg (2006)
 - Fornberg & Lehto (2011)
 - Flyer, Lehto, Blaise, W & St-Cyr (2012)
 - $\circ~$ Bollig, Flyer & Erlebacher (2012)

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Key Steps:

- 1. For each node \mathbf{x}_i , choose n-1 of it's nearest neighbors: $X_i = \{\mathbf{x}_i\}_{i=1}^n$, with $\mathbf{x}_1 = \mathbf{x}_i$.
- 2. Approximate $\mathbf{p}_x \cdot \nabla f \Big|_{\mathbf{x}_i}$ using linear a combination of the values of f sampled at X_i :

$$\mathbf{p}_x \cdot \nabla f \bigg|_{\mathbf{x}_j} = \sum_{i=1}^n c_i f(\mathbf{x}_i)$$

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Key Steps:

3. Choose the weights c_i such the approximation is exact for $\{\phi(\|\mathbf{x} - \mathbf{x}_k\|)\}_{k=1}^n$:

$$\underbrace{\left[\mathbf{p}_{x}\cdot\nabla\phi(\|\mathbf{x}-\mathbf{x}_{k}\|)\right]}_{\mathcal{D}_{x}} \equiv \sum_{i=1}^{n}c_{i}\phi(\|\mathbf{x}_{k}-\mathbf{x}_{i}\|), \ k=1,\ldots,n$$

Similar to standard FD formulas that use polynomials.

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3. The weights $\{c_i\}_{i=1}^n$ can be computed by solving:

 $\begin{bmatrix} \phi(\|\mathbf{x}_1 - \mathbf{x}_1\|) & \phi(\|\mathbf{x}_1 - \mathbf{x}_2\|) \cdots \phi(\|\mathbf{x}_1 - \mathbf{x}_N\|) \\ \phi(\|\mathbf{x}_2 - \mathbf{x}_1\|) & \phi(\|\mathbf{x}_2 - \mathbf{x}_2\|) \cdots \phi(\|\mathbf{x}_2 - \mathbf{x}_N\|) \\ \vdots & \vdots & \ddots & \vdots \\ \phi(\|\mathbf{x}_N - \mathbf{x}_1\|) & \phi(\|\mathbf{x}_N - \mathbf{x}_2\|) \cdots \phi(\|\mathbf{x}_N - \mathbf{x}_N\|) \end{bmatrix} \begin{bmatrix} c_1 \\ c_2 \\ \vdots \\ c_N \end{bmatrix} = \begin{bmatrix} \mathcal{D}_x \phi(\|\mathbf{x}_j - \mathbf{x}_1\|) \\ \mathcal{D}_x \phi(\|\mathbf{x}_j - \mathbf{x}_2\|) \\ \vdots \\ \mathcal{D}_x \phi(\|\mathbf{x}_j - \mathbf{x}_n\|) \end{bmatrix}$

4. Combine all the weights into a differentiation matrix.

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- Example differentiation matrix (DM) for N=16384, n=101:



• Compare to the global RBF method, which results in a dense differentiation matrix.

RBF-FD method for transport equation

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• Continuous transport equation for some $\mathbf{u} = (u, v, w) \in T_X \mathbb{S}^2$:

$$h_t + \mathbf{u} \cdot (P\nabla h) = 0$$

- Let h and $\mathbf{u} = (u, v, w)$ be sampled at X.
- Semi-discrete formulation (method-of-lines) of transport equation:

 $\underline{h}_t = -\left(\operatorname{diag}(\underline{u})D_N^x + \operatorname{diag}(\underline{v})D_N^y + \operatorname{diag}(\underline{w})D_N^z\right)\underline{h} = -D_N\underline{h}.$

- Advance the system in time using some standard ODE solver.
- This is a purely hyperbolic problem and temporal stability is an issue.
 - We stabilize the method by including some high-order diffusion operator L_N (hyperviscosity):

$$\underline{h}_t = -D_N \underline{h} + \mu L_N \underline{h}$$

- L_N is a discrete approximation to a high power of the Laplacian: Δ^{2k} .

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• Solid body rotation of a non-smooth cosine bell (Williamson et. al. JCP (1992))

Stream Function for flow

$$\psi(\mathbf{x}) = \cos(\alpha)z + \sin(\alpha)y \qquad \alpha = \pi/2$$
 (flow over the poles)

Initial condition (non-smooth: jump in second derivative)





Details:

- Gaussian RBF
- Stabilization required.
- Minimum energy node sets used.

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- Convergence results as number of nodes N increases (Fornberg & Lehto, 2011)
- Error results are for 10 complete revolution of the cosine bell.



- Errors compare favorably with the global RBF method.
- RBF-FD method much more computationally efficient than global method.

RBFs and partition-of-unity (RBF-PUM) on the sphere

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- Recent work with my graduate student Kevin Aiton. G.B. Wright & K. Aiton. A radial basis function partition of unity method for transport on the sphere. In preparation.



Background references for interpolation with RBF-PUM

- R. Cavoretto & A. DeRossi, Fast and accurate interpolation of large scattered data sets on the sphere. J. Comput. Appl. Math. 234 (2010), 1505–1521.
- R. Cavoretto & A. DeRossi, Spherical interpolation using the partition of unity method: An efficient and flexible algorithm. *Appl. Math. Lett.* 25 (2012), 1251-1256.

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Key Steps:

1. Generate a set of overlapping patches (spherical caps) $\Omega = \{\Omega_k\}_{k=1}^M$ that creates a uniform partition of the sphere with respect to the density of the nodes in X, and each patch contains roughly n nodes of X.

• Consider $X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$, where $\mathbf{x}_j = (x_j, y_j, z_j)$:



M total patches n nodes per patch $\xi_k = ext{center of patch } \Omega_k$

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M total patches n nodes per patch $\xi_k = ext{center of patches } \Omega_k$

Key Steps:

2. Letting X_k denote the set of nodes in patch Ω_k , construct RBF interpolants s_k , for $k = 1, \ldots, M$:

$$s_k(\mathbf{x}) = \sum_{j=1}^n c_j^k \phi(\|\mathbf{x} - \mathbf{x}_j^k\|)$$

• Consider $X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$, where $\mathbf{x}_j = (x_j, y_j, z_j)$:



M total patches n nodes per patch $\xi_k = ext{center of patches } \Omega_k$

Key Steps:

3. Define weight functions $w_k : \mathbb{S}^2 \to \mathbb{R}, k = 1 \dots, M$, with the properties that each w_k is compactly supported over Ω_k and the set of all w_k form a partition-of-unity over Ω :

$$\sum_{k=1}^{M} w_k(\mathbf{x}) \equiv 1, \, \mathbf{x} \in \Omega$$

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• Consider
$$X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$$
, where $\mathbf{x}_j = (x_j, y_j, z_j)$:



M total patches n nodes per patch $\xi_k = ext{center of patches } \Omega_k$

Key Steps:

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• Consider $X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$, where $\mathbf{x}_j = (x_j, y_j, z_j)$:



M total patches n nodes per patch $\xi_k = ext{center of patches } \Omega_k$

Key Steps:

Weight function details:

$$\psi_k(\mathbf{x}) = \psi\left(\frac{\|\mathbf{x} - \boldsymbol{\xi}_k\|}{\rho_k}\right)$$

 $\rho_k = \text{radius of patch } \Omega_k$ ψ has compact support over [0, 1]

$$w_k(\mathbf{x}) = \frac{\psi_k(\mathbf{x})}{\sum_{i=1}^{M} \psi_i(\mathbf{x})}$$

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• Consider $X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$, where $\mathbf{x}_j = (x_j, y_j, z_j)$:



M total patches n nodes per patch $\xi_k = ext{center of patches } \Omega_k$

Key Steps:

4. Create a global interpolant for X as M

$$s(\mathbf{x}) = \sum_{k=1}^{n} w_k(\mathbf{x}) s_k(\mathbf{x})$$

• Consider $X = {\mathbf{x}_j}_{j=1}^N \subset \mathbb{S}^2$, where $\mathbf{x}_j = (x_j, y_j, z_j)$:



M total patches n nodes per patch $\xi_k = ext{center of patches } \Omega_k$

Key Steps:

5. Apply projected gradient operator $\mathcal{D}_x := \mathbf{p}_x \cdot \nabla$ to interpolant and evaluate at each node \mathbf{x}_j :

$$\mathcal{D}_x s(\mathbf{x}) \Big|_{\mathbf{x} = \mathbf{x}_j} = \sum_{k=1}^M \mathcal{D}_x \left(w_k(\mathbf{x}) s_k(\mathbf{x}) \right) \Big|_{\mathbf{x} = \mathbf{x}_j}$$

Weights can be generated and stored in a differentiation matrix.

Choosing the nodes and patches

<u>Nodes</u>: We use the *maximal determinant* (MD) node sets, which are quasi-uniformly distributed over the sphere. R.S. Womersley & I. Sloan (2001)

<u>Patches</u>: We use *minimum energy* (ME) points, which are also quasiuniformly distributed over the sphere. D.P. Hardin & E.B. Saff (2004)



<u>Parameters</u>: Given N nodes, there are 2 parameters to choose for determining the total number of patches M:

- \circ n = approx. number of nodes in each patch;
- \circ q = measure of the amount the patches overlap.

Choosing the nodes and patches



• Using the quasi-uniformity of the nodes and patches, we compute the **radii of the patches** using the approximation:

$$\rho \approx 2\sqrt{\frac{n}{N}}$$

• The overlap parameter q determines the **average number of patches a node belongs to**, and satisfies the relationship:

$$\frac{4\pi}{M} \approx \frac{\pi \rho^2}{q} \implies M = \left\lceil q \frac{N}{n} \right\rceil$$

Choosing the nodes and patches

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- Illustration of the patches for N=4096, n=100, and different q:



The computational cost for evaluating a derivative grows at most linearly with q and n.

Comparison: RBF-FD and RBF-PUM

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- Example differentiation matrix (DM) for N=16384, n=101:



• Recall global RBF-type methods result in dense matrices.

RBF-PUM method for transport equation

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• Continuous transport equation for some $\mathbf{u} = (u, v, w) \in T_X \mathbb{S}^2$:

$$h_t + \mathbf{u} \cdot (P\nabla h) = 0$$

- Let h and $\mathbf{u} = (u, v, w)$ be sampled at X.
- Semi-discrete formulation (method-of-lines) of transport equation:

$$\underline{h}_t = -\left(\operatorname{diag}(\underline{u})D_N^x + \operatorname{diag}(\underline{v})D_N^y + \operatorname{diag}(\underline{w})D_N^z\right)\underline{h} = -D_N\underline{h}.$$

- Advance the system in time using some standard ODE solver.
- This is a purely hyperbolic problem and temporal stability will be an issue.
 - We stabilize the method by including some high-order diffusion-type operator L_N (hyperviscosity):

$$\underline{h}_t = -D_N \underline{h} + \mu L_N \underline{h}$$

 $-L_N$ is a like a discrete approximation to a high power of the Laplacian.

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Details for all numerical results:

- We use the Gaussian RBF
- Time-step Δt is not optimized

- Overlap is set to q=4
- We set $\epsilon = a_n \sqrt{N} + b_n$
- Solid body rotation of a non-smooth cosine bell (Williamson et. al. JCP (1992))

Stream Function for flow

$$\psi(\mathbf{x}) = \cos(\alpha)z + \sin(\alpha)y \qquad \alpha = \pi/2$$
 (flow over the poles)

Initial condition (non-smooth: jump in second derivative)

$$h(\mathbf{x}) = \begin{cases} \frac{1}{2} \left(1 + \cos(3\pi r(\mathbf{x})) \right) & r(\mathbf{x}) < 1/3 \\ 0 & r(\mathbf{x}) \ge 1/3 \end{cases}$$

 $r(\mathbf{x}) = \arccos(x)$

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Plots of the RBF-PUM solution for $N=12544, n=100, \Delta t=2\pi/1600$:





Plots of the RBF-PUM error for $N=12544, n=100, \Delta t=2\pi/1600$:



Comparison of the errors for RBF-FD and RBF-PUM:

Relative ℓ_2 error vs. \sqrt{N} (logscale)



Convergence rates are as expected given smoothness of the initial condition.

Numerical results: deformational flow

Deformational/rotational flow (R.D. Nair and P.H. Lauritzen, JCP (2010))



Smooth initial condition:

Non-smooth

Simulation for non-smooth IC, N=20736, n=100, $\Delta t=5/2400$



Time = 0.000

Convergence plots for increasing N and n, $\Delta t=5/2400$

• Non-smooth initial condition:



Convergence rates are as expected given smoothness of the initial condition.

Numerical results: deformational flow

DRWA 2013 Lecture 4

Convergence plots for increasing N and n, $\Delta t=5/2400$

• Smooth initial condition:



• Test: Deformational/Rotational flow smooth initial condition

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• MATLAB R2013a, Intel Xeon 3.1GHz processor



Derivative approx.	Global RBF	RBF-FD*	RBF-PUM*
Construction:	$O(N^3)$	$O(n^3N)$	$O(n^3M)$
Evaluation:	$O(N^2)$	O(nN)	O(nN)

*Constants for the RBF-PUM are higher than for RBF-FD.

- The Global RBF collocation method is competitive in terms of accuracy per degree of freedom.
- It are not competitive in terms of computational complexity.
- The RBF generated finite difference (RBF-FD) method shows great promise in terms of accuracy and computational cost.
 - More comparisons with other state-of-the art methods in the next lecture.
 - Parallelization on multi-GPU has already been implemented (Bollig, Flyer, & Erlebacher, 2012).
- The RBF Partition of unity method (RBF-PUM) also shows great promise.
 - More comparisons are needed between RBF-PUM and RBF-FD in terms of computational cost to achieve a certain accuracy.
 - Parallel implementations also needed.
- More work is needed in developing stable algorithms for "flat" RBFs when working on patches of the sphere.