Proper homothetic vector fields in Bianchi Type-I Space-Time

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Abstract. A study of proper homothetic vector field in Bianchi type-1 space-times is given by using direct integration technique. Using the above mentioned technique we show that the above space-times admit proper homothetic vector field in a very special choice of f, k and h. The dimensions of homothetic algebra are 4, 5, 7 and 11.

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Key words: proper homothetic vector fields, direct integration technique.

1 Introduction

Throughout M is representing the four dimensional, connected, hausdorff space-time manifold with Lorentz metric g of signature (-, +, +, +). The curvature tensor associated with g_{ab} , through Levi-Civita connection, is denoted in component form by $R^a{}_{bcd}$, and the Ricci tensor components are $R_{ab} = R^c{}_{acb}$. The usual covariant, partial and Lie derivatives are denoted by a semicolon, a comma and the symbol L, respectively. Round and square brackets denote the usual symmetrization and skew-symmetrization, respectively.

Any vector field X on M can be decomposed as

(1.1)
$$X_{a;b} = \frac{1}{2}h_{ab} + F_{ab}$$

where $h_{ab} = h_{ba}$ and $F_{ab} = -F_{ba}$ are symmetric and skew-symmetric tensor on M, respectively. If

$$h_{ab} = \alpha g_{ab}, \quad \alpha \in R$$

equivalent (1.2)

1.2)
$$g_{ab,c}X^{c} + g_{bc}X^{c}_{,a} + g_{ac}X^{c}_{,b} = \alpha g_{ab}$$

then X is called a homothetic vector field on M. If X is homothetic and $\alpha \neq 0$ then it is called proper homothetic while $\alpha = 0$ it is Killing [3, 4]. Further consequences and geometrical interpretations of (1.2) are explored in [2, 5]. It also follows from (1.2) that [2, 4]

$$L_X R^a{}_{bcd} = 0 \qquad \qquad L_X R_{ab} = 0$$

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2 Main results

A Bianchi type-1 space-time is a spatially homogeneous space-time which admits an abelian Lie algebra of isometries G_3 , acting on spacelike hypersurfaces, generated by the spacelike Killing vector fields which are:

(2.1)
$$\frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z}.$$

The line element in usual coordinate system is [7]

(2.2)
$$ds^{2} = -dt^{2} + f(t)dx^{2} + k(t)dy^{2} + h(t)dz^{2},$$

where f, k and h are some nowhere zero functions of t only. The possible Segre type of the above space-time is $\{1, 111\}$ or one of its degeneracies. A vector field X is said to be a homothetic vector field if it satisfy equation (1.2). One can write (1.2) explicitly using (2.2) we have

(2.3)
$$X_{,0}^0 = \frac{\alpha}{2}$$

(2.4)
$$fX_{,0}^1 - X_{,1}^0 = 0,$$

(2.5)
$$kX_{,0}^2 - X_{,2}^0 = 0$$

(2.6)
$$hX_{,0}^3 - X_{,3}^0 = 0,$$

(2.7)
$$\dot{f} X^0 + 2f X^1_{,1} = \alpha f,$$

(2.8)
$$kX_{,1}^2 + fX_{,2}^1 = 0,$$

(2.9)
$$hX_{,1}^3 + fX_{,3}^1 = 0.$$

(2.10)
$$\dot{k} X^0 + 2kX_{,2}^2 = \alpha \dot{k}$$

(2.11)
$$hX_{,2}^3 + kX_{,3}^2 = 0,$$

(2.12)
$$\dot{h} X^0 + 2hX^3_{,3} = \alpha h.$$

Equations (2.3), (2.4), (2.5) and (2.6) give

(2.13)
$$\begin{cases} X^{0} = \frac{\alpha}{2}t + A^{1}(x, y, z) \\ X^{1} = A_{x}^{1}(x, y, z) \int \frac{1}{t}dt + A^{2}(x, y, z) \\ X^{2} = A_{y}^{1}(x, y, z) \int \frac{1}{k}dt + A^{3}(x, y, z) \\ X^{3} = A_{z}^{1}(x, y, z) \int \frac{1}{h}dt + A^{4}(x, y, z) \end{cases} ,$$

where $A^1(x, y, z)$, $A^2(x, y, z)$, $A^3(x, y, z)$ and $A^4(x, y, z)$ are functions of integration. In order to determine $A^1(x, y, z)$, $A^2(x, y, z)$, $A^3(x, y, z)$ and $A^4(x, y, z)$ we need to integrate the remaining six equations. To avoid lengthy calculations here we will only present results for full details see [1].

Case (1) Four independent homothetic vector fields:

In this case the space-time (2.2) takes the form

$$(2.14) \ ds^2 = -dt^2 + (\alpha t + 2d_{10})^{2(1-\frac{2}{\alpha}d_5)}dx^2 + dy^2 + (\alpha t + 2d_{10})^{2(1-\frac{2}{\alpha}d_{14})}dz^2,$$

and homothetic vector fields in this case are

(2.15)
$$\begin{array}{c} X^0 = t\frac{\alpha}{2} + d_{10} \\ X^1 = xd_5 + d_6 \\ X^2 = y\frac{\alpha}{2} + d_{12} \\ X^3 = zd_{14} + d_{15} \end{array} \right\},$$

where $d_5, d_6, d_{10}, d_{12}, d_{14}, d_{15} \in \mathbb{R}$. The above space-time (2.14) admits four independent homothetic vector fields in which three are Killing vector fields which are given in (2.1) and one is proper homothetic vector field which is

(2.16)
$$Z^1 = (t, 0, y, 0).$$

Case (2) Four independent homothetic vector fields:

In this case the space-time (2.2) takes the form

$$ds^{2} = -dt^{2} + (\alpha t + 2d_{10})^{2(1 - \frac{2}{\alpha}d_{5})} dx^{2} + (\alpha t + 2d_{10})^{2(1 - \frac{2}{\alpha}d_{11})} dy^{2} + (\alpha t + 2d_{10})^{2(1 - \frac{2}{\alpha}d_{14})} dz^{2}.$$
(2.17)

Homothetic vector fields in this case are

(2.18)
$$\begin{cases} X^0 = t\frac{\alpha}{2} + d_{10} \\ X^1 = xd_5 + d_6 \\ X^2 = yd_{11} + d_{12} \\ X^3 = zd_{14} + d_{15} \end{cases} ,$$

where $d_5, d_6, d_{10}, d_{11}, d_{12}, d_{14}, d_{15} \in R$. The above space-time (2.17) admits four independent homothetic vector fields in which three are Killing vector fields which are given in (2.1) and one is proper homothetic vector field which is

(2.19)
$$Z^2 = (t, 0, 0, 0).$$

Case (3) Five independent homothetic vector fields:

In this case the space-time (2.2) takes the form

(2.20)
$$ds^{2} = -dt^{2} + (\alpha t + 2d_{4})^{2(1-\frac{2}{\alpha}d_{7})}dx^{2} + (dy^{2} + dz^{2}),$$

and homothetic vector fields in this case are

(2.21)
$$\begin{cases} X^0 = t\frac{\alpha}{2} + d_4 \\ X^1 = xd_5 + d_8 \\ X^2 = y\frac{\alpha}{2} + zd_5 + d_6 \\ X^3 = z\frac{\alpha}{2} - yd_5 + d_3 \end{cases} ,$$

where $d_3, d_4, d_5, d_6, d_7, d_8 \in R$. The above space-time (2.20) admits five independent homothetic vector fields in which four are Killing vector fields and one is proper homothetic vector field which is

(2.22)
$$Z^3 = (t, 0, y, z).$$

Case (4) Five independent homothetic vector fields:

In this case the space-time (2.2) takes the form

$$(2.23) \ ds^2 = -dt^2 + (\alpha t + 2d_4)^{2(1-\frac{2}{\alpha}d_7)}dx^2 + (\alpha t + 2d_4)^{2(1-\frac{2}{\alpha}d_{12})}(dy^2 + dz^2).$$

Homothetic vector fields in this case are

(2.24)
$$\begin{cases} X^0 = t\frac{\alpha}{2} + d_4 \\ X^1 = xd_7 + d_8 \\ X^2 = y\frac{\alpha}{2} + zd_5 + d_{11} \\ X^3 = zd_{12} - yd_5 + d_{13} \end{cases} ,$$

where $d_4, d_5, d_7, d_8, d_{11}, d_{12}, d_{13} \in R$. The above space-time (2.24) admits five independent homothetic vector fields in which four are Killing vector fields and one is proper homothetic vector field which is

(2.25)
$$Z^4 = (t, 0, y, 0).$$

Case (5) Seven independent homothetic vector fields:

In this case the space-time (2.2) takes the form

(2.26)
$$ds^{2} = -dt^{2} + (\alpha t + 2d_{9})^{2(1-\frac{2}{\alpha}d_{8})}(dx^{2}dy^{2} + dz^{2}).$$

It follows from [6] homothetic vector fields in this case are

(2.27)
$$\begin{cases} X^0 = t\frac{\alpha}{2} + d_9 \\ X^1 = xd_8 - yd_{20} - zd_{18} + d_{21} \\ X^2 = yd_8 + xd_{20} + zd_{17} + d_{23} \\ X^3 = zd_8 + xd_{18} - yd_{17} + d_{24} \end{cases} \},$$

where $d_8, d_9, d_{17}, d_{18}, d_{20}, d_{21}, d_{23}, d_{24} \in R$. The above space-time (2.26) admits seven independent homothetic vector fields in which six are Killing vector fields and one is proper homothetic vector field which is

$$(2.28) Z^5 = (t, 0, 0, 0).$$

Case (6) Eleven independent homothetic vector fields:

In this case the above space-time (2.2) becomes Minkowski space-time

(2.29)
$$ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$$

homothetic vector fields in this case are

$$(2.30) \qquad \begin{array}{c} X^{0} = t\frac{\alpha}{2} + xd_{11} + yd_{10} + zd_{12} + d_{13} \\ X^{1} = x\frac{\alpha}{2} + td_{11} + yd_{17} + zd_{15} + d_{18} \\ X^{2} = y\frac{\alpha}{2} + td_{10} + zd_{14} - xd_{17} + d_{19} \\ X^{3} = z\frac{\alpha}{2} + td_{12} - xd_{15} - yd_{14} + d_{16} \end{array} \right\},$$

where $d_{10}, d_{11}, d_{12}, d_{13}, d_{14}, d_{15}, d_{16}, d_{17}, d_{18}, d_{19} \in R$. The above space-time (2.29) admits eleven independent homothetic vector fields in which ten are Killing vector fields and one is proper homothetic vector field which is

(2.31)
$$Z^6 = (t, x, y, z)$$

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Summary

In this paper a study of Bianchi type-I space-times according to their proper homothetic vector fields is given by using the direct integration technique. From the above study we obtain the following:

(i) The space-times which admit four independent homothetic vector fields are given in equations (2.14) and (2.17) (see for details cases (1) and (2)).

(ii) The space-times which admit five independent homothetic vector fields are given in equations (2.20) and (2.23) (see for details cases (3) and (4)).

(iii) The space-time which admits seven independent homothetic vector fields are given in equation (2.26) (see for details case (5)).

(iv) The space-time which admits eleven independent homothetic vector fields are given in equation (2.29) (see for details case (6)).

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